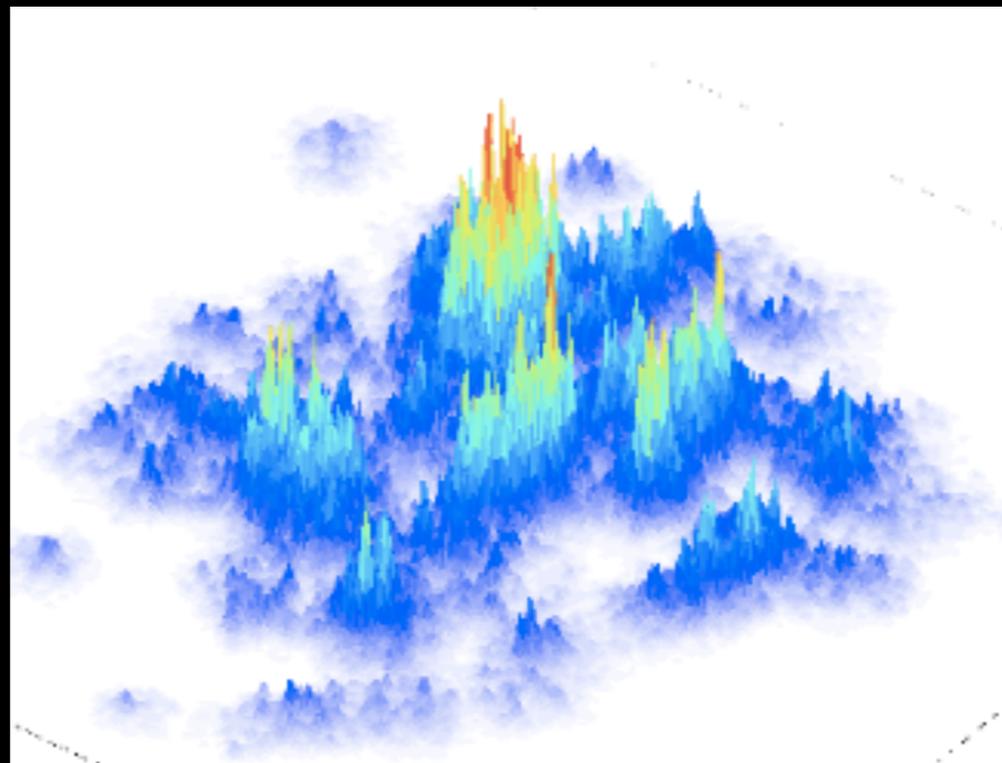


# Initial State Fluctuations and Final State Correlations



Gunther Roland



- Two lectures
  - Geometry, geometry fluctuations and hydrodynamic flow in nucleus-nucleus collisions
  - Correlations in pp, pA and dA collisions

As this is a workshop in honor of Wit Busza, I decided to provide a *historical* overview of the developments in our understanding of flow and correlations over the last 15 years



This will naturally highlight contributions made under Wit's leadership in PHOBOS and the MIT heavy-ion group

I will mostly focus on experimental/conceptual developments

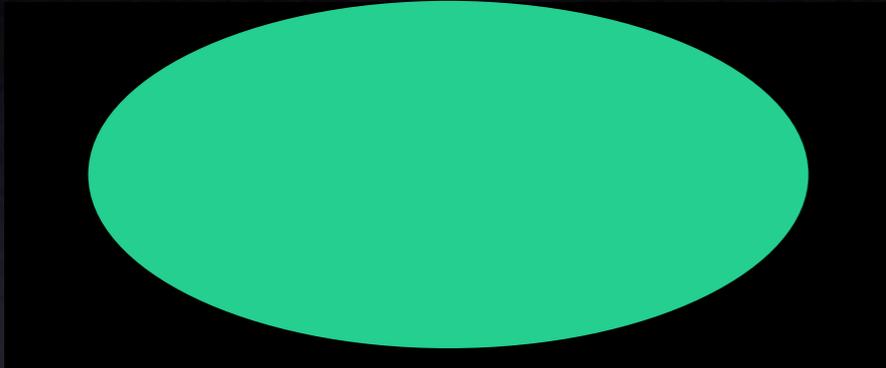
In NA49 I worked on some of the earliest studies of event-by-event fluctuations in transverse momentum and particle ratios, or more generally, particle correlations

Whenever the subject of correlations and in particular correlation functions came up in our meetings at MIT, Wit would ask:

Is there anything that we have learned from correlation measurements?

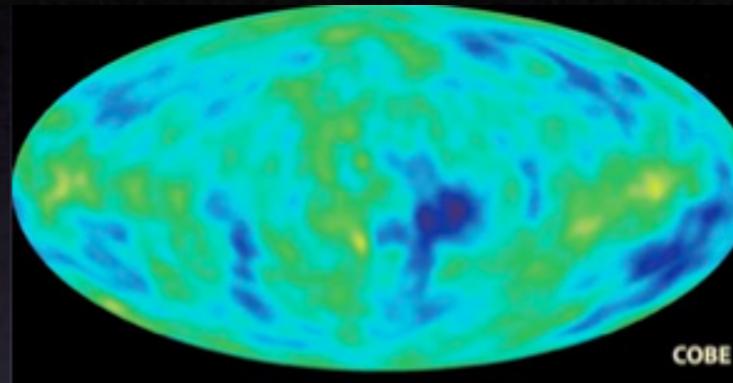
# Analogy: Cosmic Microwave Background

## Cosmic microwave background



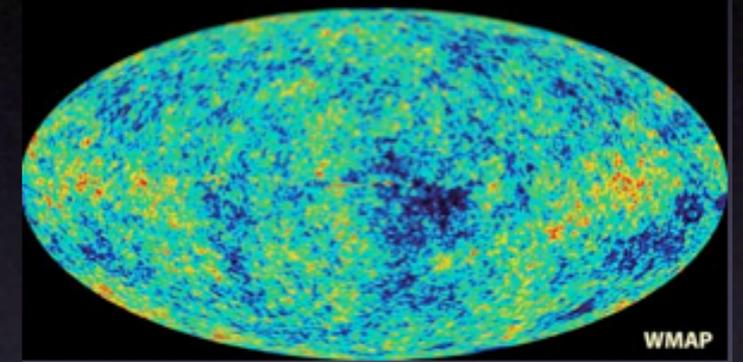
Penzias, Wilson  
1964

$$\langle T \rangle = 3\text{K}$$

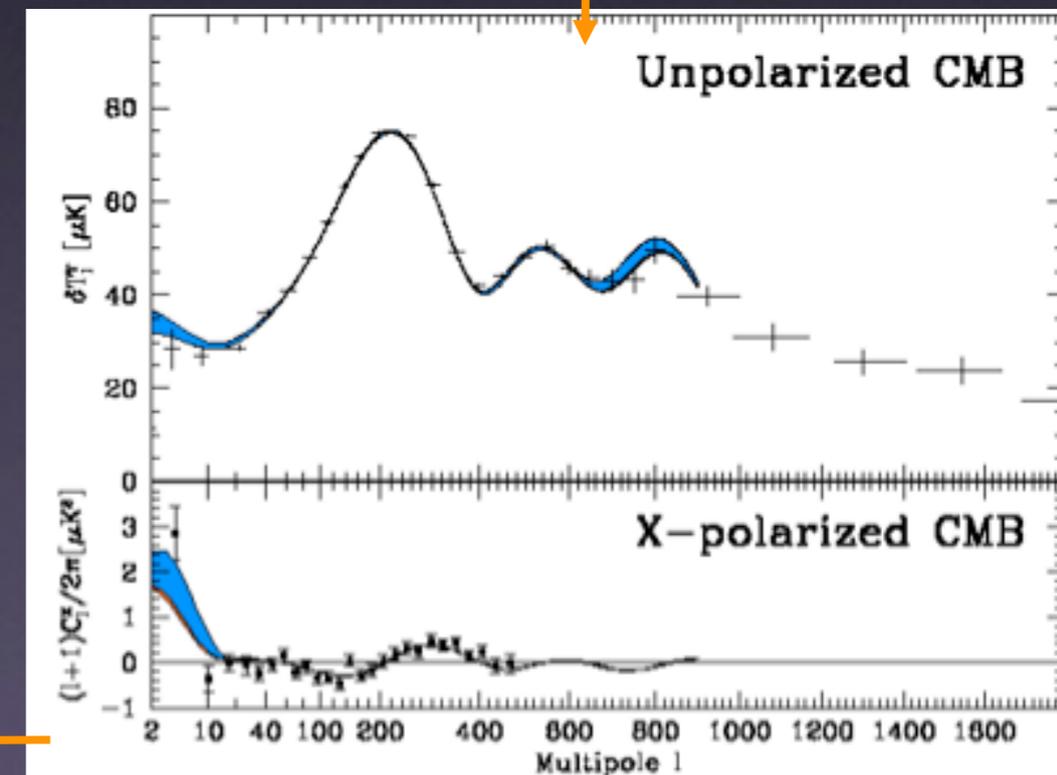


COBE 1992

$$\Delta T/T \sim 10^{-5}$$



WMAP et al.  
2003



"Best" Cosmological Parameters:  
Table 1 from Wilkinson Microwave Anisotropy Probe (WMAP) Observations:  
Preliminary Maps and Basic Results,  
C. L. Bennett et al. (2003), accepted by the *Astrophysical Journal*,  
available at <http://lambda.gsfc.nasa.gov/>

| Description  | Symbol                     | Value                 | + uncertainty         | - uncertainty         |
|--|----------------------------|-----------------------|-----------------------|-----------------------|
| Total density  | $\Omega_{tot}$             | 1.02                  | 0.02                  | 0.02                  |
| Equation of state of quintessence                                | $w$                        | $< -0.78$             | 95% CL                | —                     |
| Dark energy density  | $\Omega_{\Lambda}$         | 0.73                  | 0.04                  | 0.04                  |
| Baryon density   | $\Omega_b h^2$             | 0.0224                | 0.0009                | 0.0009                |
| Baryon density   | $\Omega_b$                 | 0.044                 | 0.004                 | 0.004                 |
| Baryon density ( $\text{cm}^{-3}$ )                              | $n_b$                      | $2.5 \times 10^{-7}$  | $0.1 \times 10^{-7}$  | $0.1 \times 10^{-7}$  |
| Matter density   | $\Omega_m h^2$             | 0.135                 | 0.008                 | 0.009                 |
| Matter density   | $\Omega_m$                 | 0.27                  | 0.04                  | 0.04                  |
| Light neutrino density   | $\Omega_{\nu} h^2$         | $< 0.0078$            | 95% CL                | —                     |
| CMB temperature (K)  | $T_{mb}$                   | 2.725                 | 0.002                 | 0.002                 |
| CMB photon density ( $\text{cm}^{-3}$ )                          | $n_{\gamma}$               | 410.4                 | 0.9                   | 0.9                   |
| Baryon-to-photon ratio   | $\eta$                     | $6.1 \times 10^{-10}$ | $0.2 \times 10^{-10}$ | $0.2 \times 10^{-10}$ |
| Baryon-to-matter ratio   | $\Omega_b \Omega_m^{-1}$   | 0.17                  | 0.01                  | 0.01                  |
| Fluctuation amplitude in $8h^{-1}$ Mpc spheres                   | $\sigma_8$                 | 0.84                  | 0.04                  | 0.04                  |
| Low- $z$ cluster abundance scaling                               | $\sigma_8 \Omega_m^{-0.5}$ | 0.44                  | 0.04                  | 0.05                  |
| Power spectrum normalization (at $k_0 = 0.05 \text{ Mpc}^{-1}$ ) | $A$                        | 0.833                 | 0.086                 | 0.083                 |
| Scalar spectral index (at $k_0 = 0.05 \text{ Mpc}^{-1}$ )        | $n_s$                      | 0.96                  | 0.03                  | 0.03                  |
| Running index slope (at $k_0 = 0.05 \text{ Mpc}^{-1}$ )          | $dn_s/d \ln k$             | -0.051                | 0.016                 | 0.016                 |
| Tensor-to-scalar ratio (at $k_0 = 0.002 \text{ Mpc}^{-1}$ )      | $r$                        | $< 0.90$              | 95% CL                | —                     |
| Redshift of decoupling   | $z_{dec}$                  | 1089                  | 1                     | 1                     |
| Thickness of decoupling (FWHM)                                   | $\Delta z_{dec}$           | 196                   | 2                     | 2                     |
| Hubble constant  | $h$                        | 0.71                  | 0.04                  | 0.03                  |
| Age of universe (Gyr)  | $t_0$                      | 13.7                  | 0.2                   | 0.2                   |
| Age at decoupling (kyr)  | $t_{dec}$                  | 379                   | 8                     | 7                     |
| Age at reionization (Myr, 95% CL)                                | $t_r$                      | 180                   | 20                    | 80                    |
| Decoupling time interval (kyr)                                   | $\Delta t_{dec}$           | 118                   | 3                     | 2                     |
| Redshift of matter-energy equality                               | $z_{eq}$                   | 3233                  | 194                   | 210                   |
| Reionization optical depth                                       | $\tau$                     | 0.17                  | 0.04                  | 0.04                  |
| Redshift of reionization (95% CL)                                | $z_r$                      | 20                    | 10                    | 9                     |
| Sound horizon at decoupling ( $^{\circ}$ )                       | $\theta_A$                 | 0.598                 | 0.002                 | 0.002                 |
| Angular size distance to decoupling (Gpc)                        | $d_A$                      | 14.0                  | 0.2                   | 0.3                   |
| Acoustic scale   | $\ell_A$                   | 301                   | 1                     | 1                     |
| Sound horizon at decoupling (Mpc)                                | $r_s$                      | 147                   | 2                     | 2                     |



# Correlations and Fluctuations

- Correlations/fluctuations induced by interaction of perturbations with medium
- Essential for understanding the nature of medium and of perturbations
- Need to interface experiment/theory and experiment/experiment
- Find a representation of correlation content, i.e. our 'Power spectrum'

Workshop on Correlations and Fluctuations in  
Relativistic Heavy Ion Collisions  
MIT, 4/21 to 4/23 2005  
Organizers T.Trainor, G.Roland





# A brief history of correlations in HIC



Is there anything that we have learned from correlation measurements?

2002

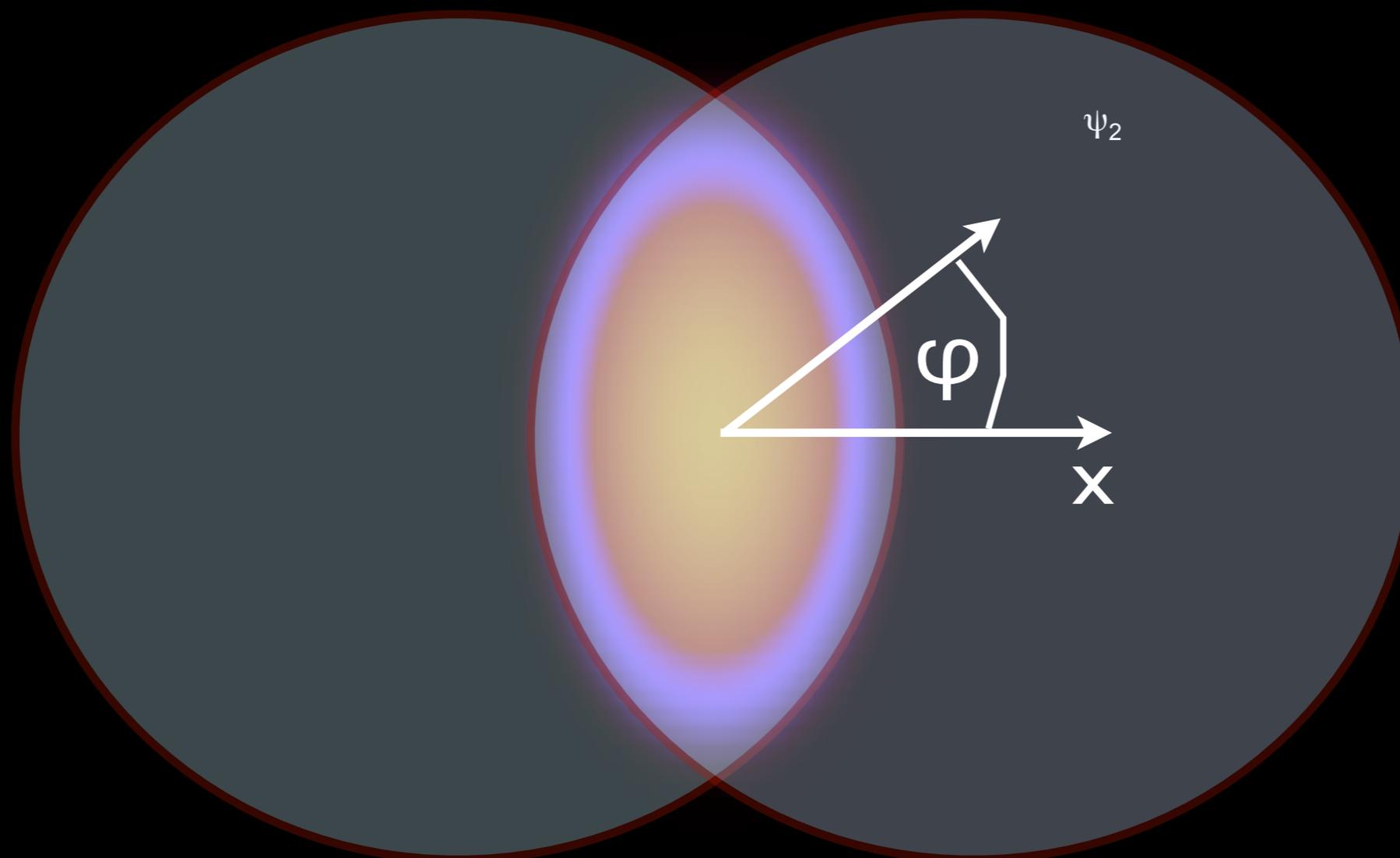
2006

2010

2014

# Anisotropic initial state

Non-central collision (Transverse plane)

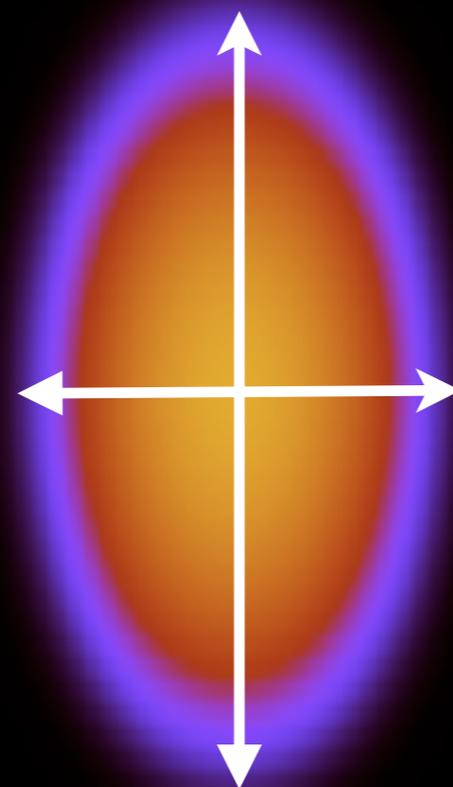


Nucleus II  
(into plane)

Nucleus I  
(out-of-plane)

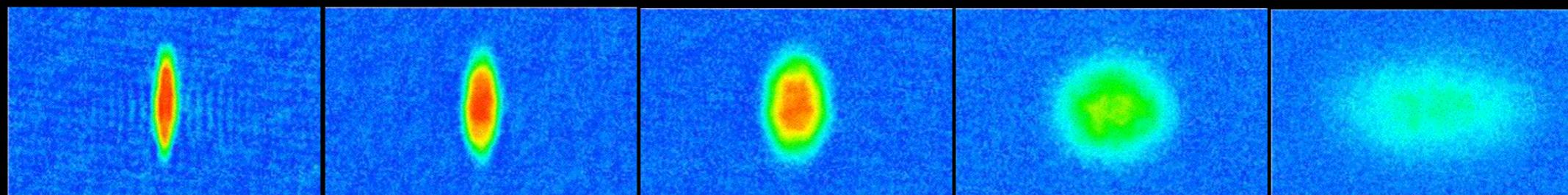
Initial overlap in transverse plane  
is asymmetric in azimuth  $\varphi$

Small pressure gradient



Large pressure gradient

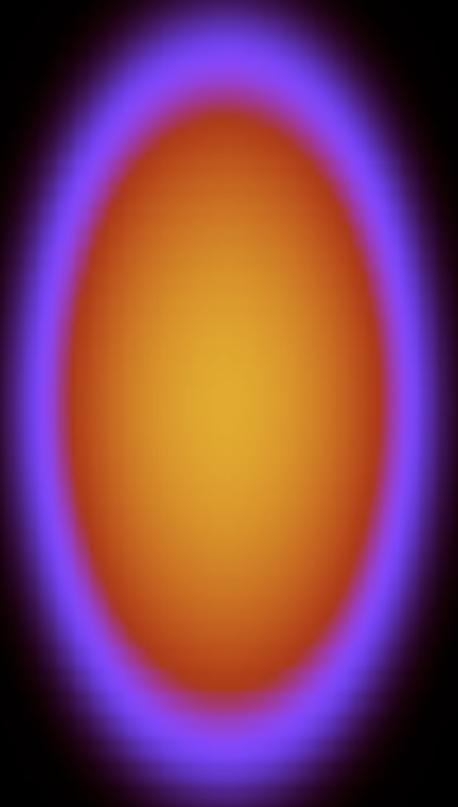
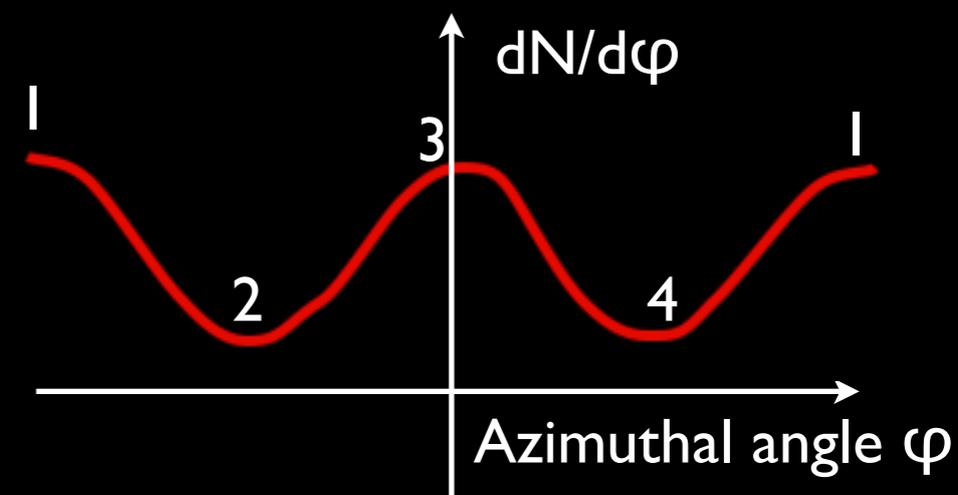
Time →



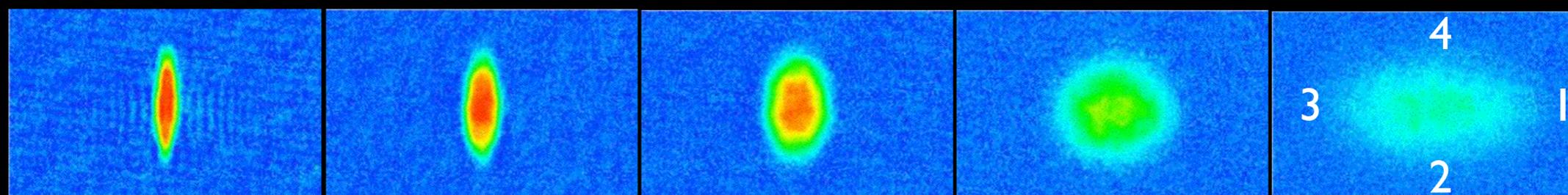
n.b. picture shows expansion of ultracold atoms released from trap

# “Elliptic Flow”

Shape information transformed into momentum space  
 $\cos(2\varphi)$  modulation of azimuthal distribution



Time →

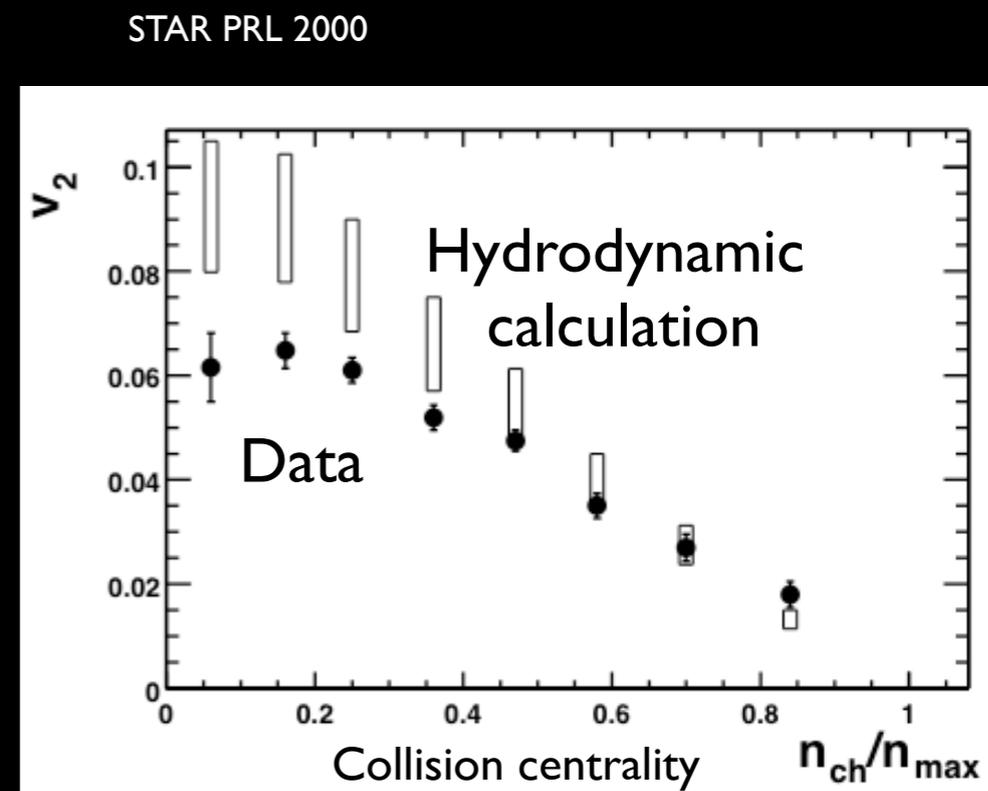
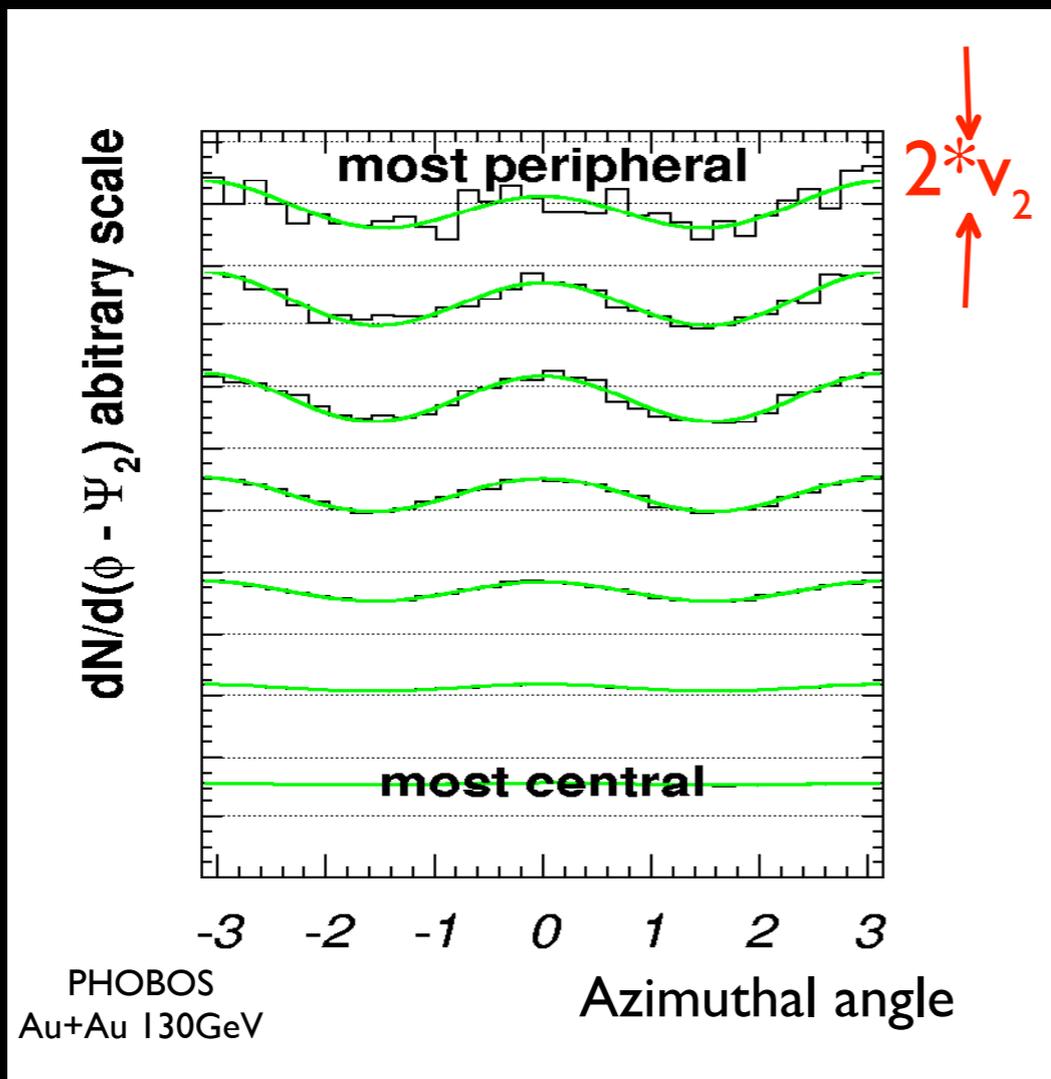


Pressure driven *hydrodynamic* expansion

Azimuthal distribution  

$$dN/d\varphi = 1 + 2 v_2 \cos(2(\varphi - \varphi_0))$$

## “Elliptic Flow”



Peripheral collisions

central collisions

“Elliptic Flow”  
is clearly seen

Close to “ideal  
hydrodynamics”

**Hydrodynamics:** conservation laws for long wavelength modes

$$\partial_\mu T^{\mu\nu} = 0$$

Generally:

$$T^{\mu\nu} = (\epsilon + P)u^\mu u^\nu - P g^{\mu\nu} + \pi^{\mu\nu}.$$

First order Navier Stokes theory:

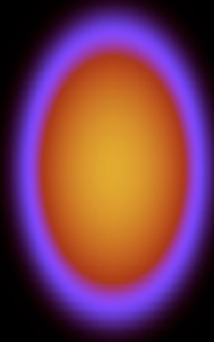
$$\pi^{\mu\nu} = \pi_{(1)}^{\mu\nu} = \eta \left( \nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} \nabla_\alpha u^\alpha \right).$$

$$\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$$

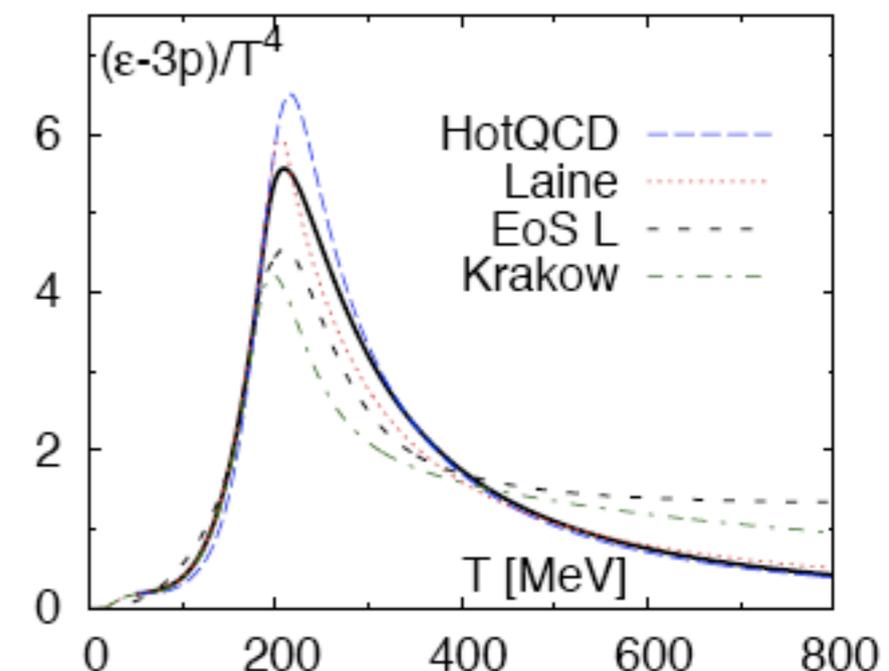
$\eta$ : Shear viscosity

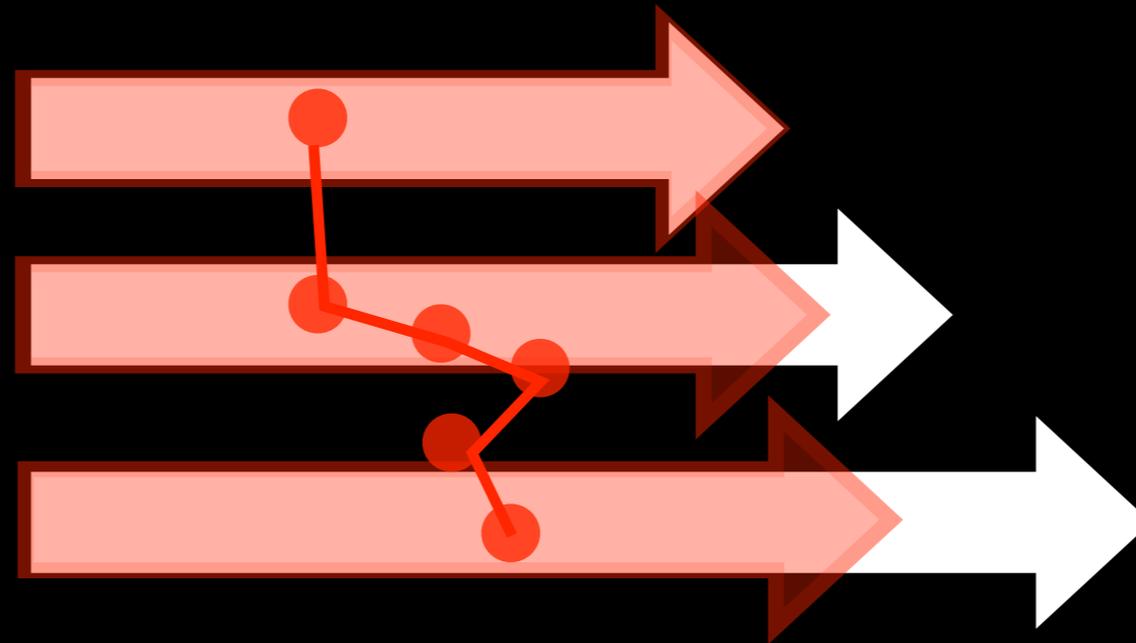
Large shear viscosity  $\rightarrow$  transport of momentum across fluid layers

+ initial conditions



+ Equation of State

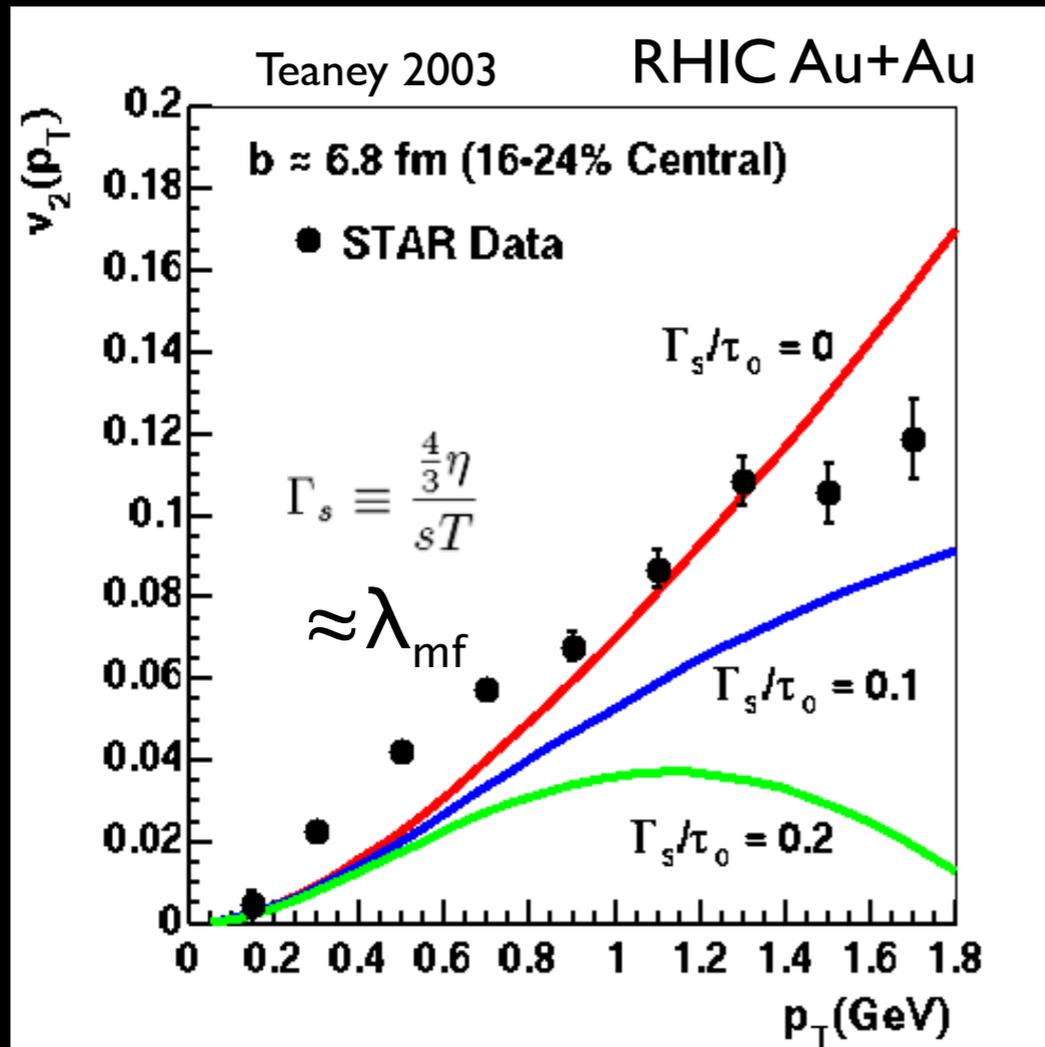




Shear viscosity: Momentum transport across fluid reduces gradients  $\rightarrow$  less elliptic flow

To compare systems: Divide by entropy density  $\rightarrow \eta/s$   
 (“ $\eta$ ” for QGP is very large, but “ $s$ ” is even larger..)

Weakly interacting gas: Large  $\eta/s$



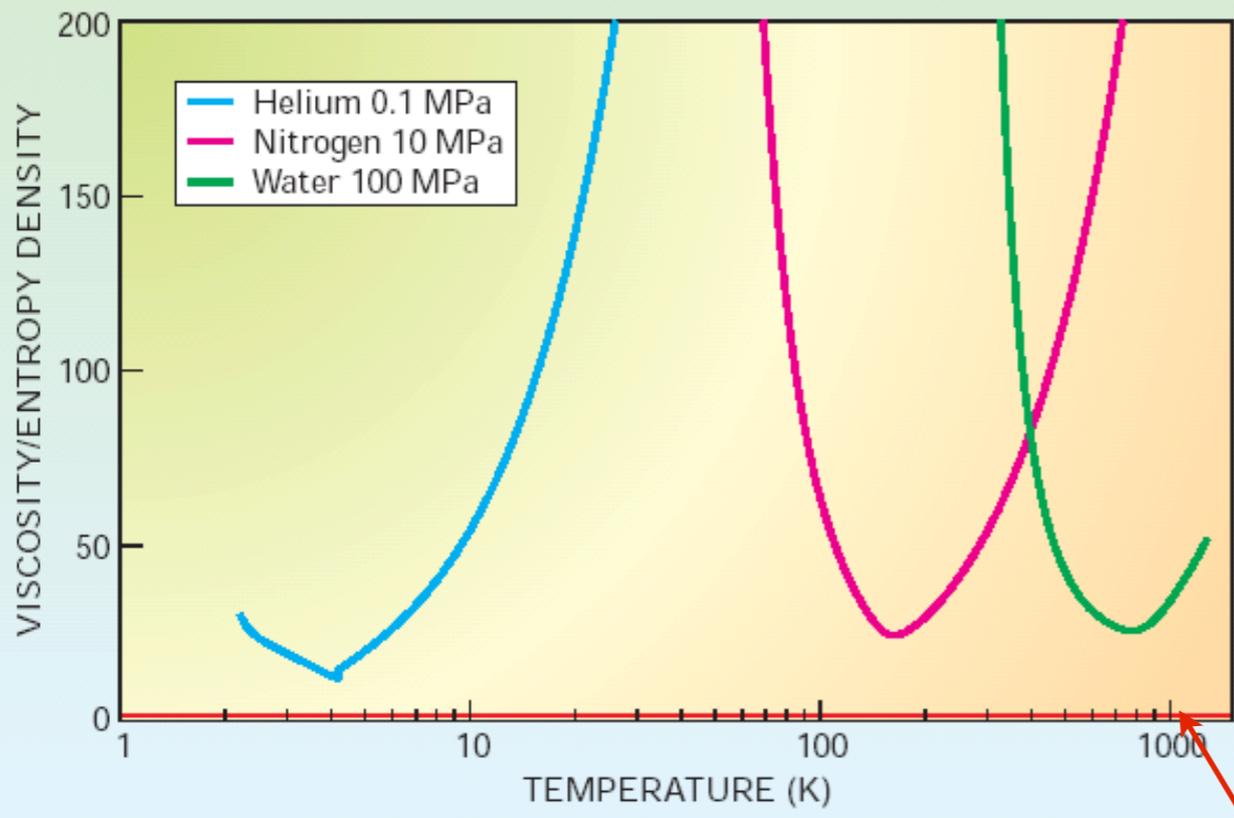
Phys.Rev. C68 (2003) 034913

First estimate of viscous corrections: Magnitude of  $v_2$  changes rapidly with  $\eta/s$

The observed elliptic flow **places** a constraint on the shear viscosity. Indeed, unless  $\Gamma_s/\tau_0$  is less than 0.1,  $v_2$  as a function of  $p_T$  falls well below the ideal curve by  $p_T \approx 1.0$  GeV. For the blast wave model, the viscous corrections to elliptic observables become large *before* the corresponding corrections to the transverse momentum spectra.



# $\eta/s \ll 1$ makes QGP unique



find eprint hep-th and topcite 1000+ and date after 2005

find | "Phys.Rev.Lett.,105" :: more

Sort by: latest first | desc. | - or rank by - | Display results: 25 results | single list

No exact match found for hep-th, using hep th instead...

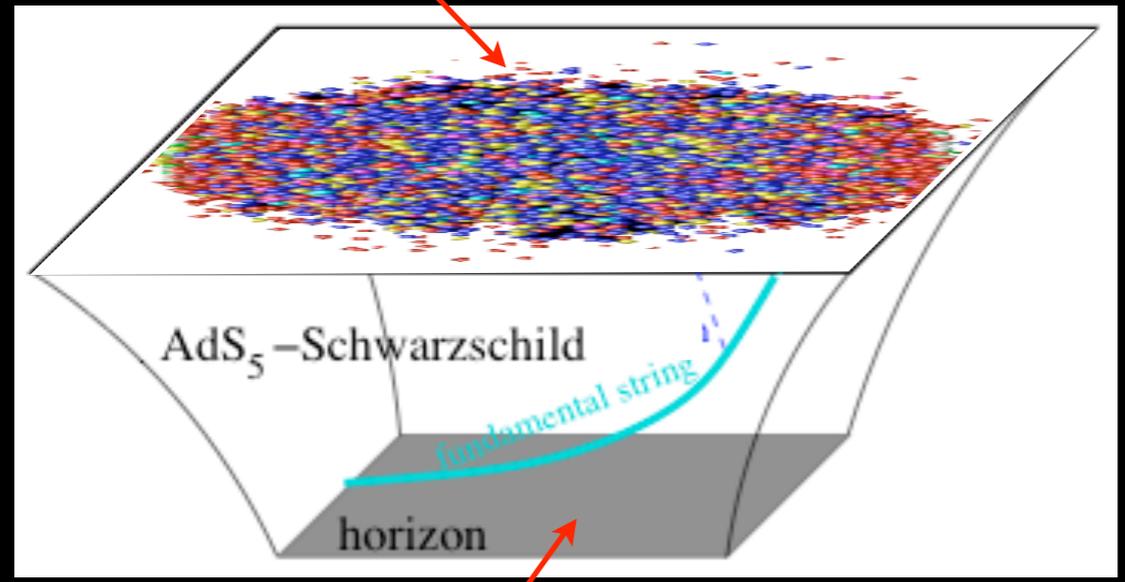
**HEP** 2 records found Search took 2.24 seconds

- Dynamics of dark energy.**  
Edmund J. Copeland (Nottingham U.), M. Sami (Jamia Millia Islamia), Shinji Tsujikawa (Gunma Coll. Tech.). Mar 2006. 84 pp.  
Published in Int.J.Mod.Phys. D15 (2006) 1753-1936  
DOI: 10.1142/S021827180600942X  
e-Print: [hep-th/0603057](#) | PDF  
[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)  
[ADS Abstract Service](#); [Int.J.Mod.Phys.D Server](#)  
[Detailed record](#) - Cited by 1812 records
- Viscosity in strongly interacting quantum field theories from black hole physics.**  
P. Kovtun, D.T. Son, A.O. Starinets (Washington U., Seattle). Mar 2004. 8 pp.  
Published in Phys.Rev.Lett. 94 (2005) 111601  
INT-PUB-04-09, UW-PT-04-04  
DOI: 10.1103/PhysRevLett.94.111601  
e-Print: [hep-th/0405231](#) | PDF  
[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)  
[ADS Abstract Service](#); [Phys. Rev. Lett. Server](#)  
[Detailed record](#) - Cited by 1007 records

Thermal Field Theory (distant cousin of QCD...)

$$\frac{\eta}{s} = \frac{\hbar}{4\pi}$$

Shear viscosity in  $N=4$  supersymmetric Yang-Mills in strong coupling, large  $N$  limit related to absorption cross-section of graviton on black three-brane in 10D classical gravity, using AdS/CFT



Viscosity in CFT

Graviton absorption by black hole



Kovtun, Son, Starinets hep-th/0405231  
Policastro, Son, Starinets hep-th/0104066

Black hole horizon

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**Number 757 #1, December 7, 2005 by Phil Schewe and Ben Stein**

### The Top Physics Stories for 2005

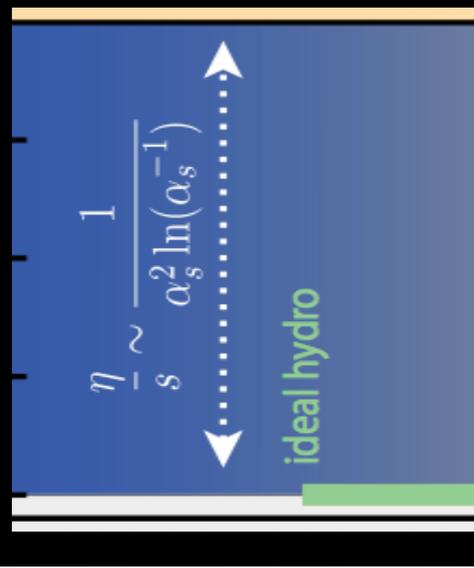
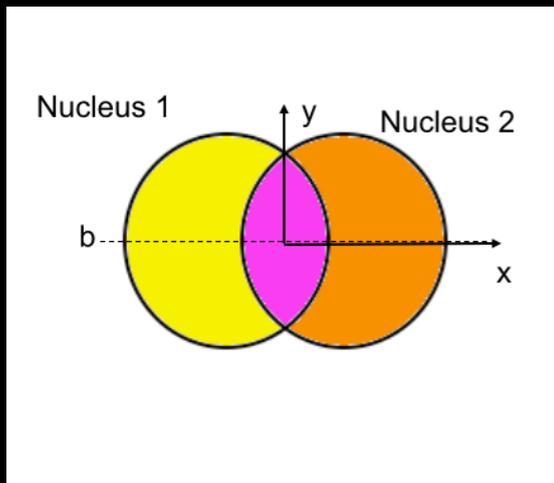
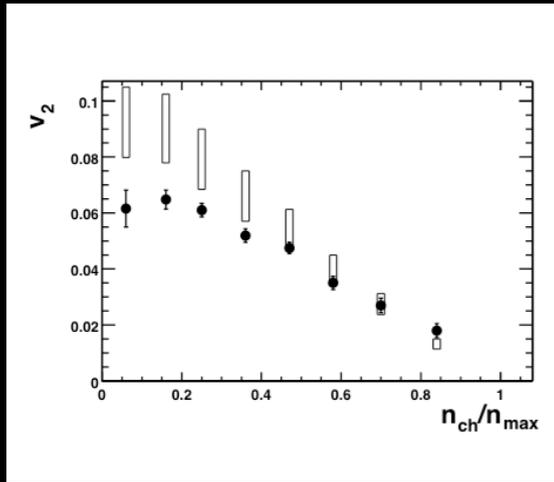
At the Relativistic Heavy Ion Collider (RHIC) on Long Island, the four large detector groups agreed, for the first time, on a consensus interpretation of several year's worth of high-energy ion collisions: the fireball made in these collisions -- a sort of stand-in for the primordial universe only a few microseconds after the big bang -- was not a gas of weakly interacting quarks and gluons as earlier expected, but something more like a liquid of strongly interacting quarks and gluons ([PNU 728](#)).

“...the fireball made in these [heavy-ion] collisions...was not a gas of weakly interacting quarks and gluons as earlier expected, but something more like a liquid..”

based on Whitepapers by BRAHMS, PHENIX, PHOBOS and STAR collaborations at RHIC



# A brief history of correlations in HIC

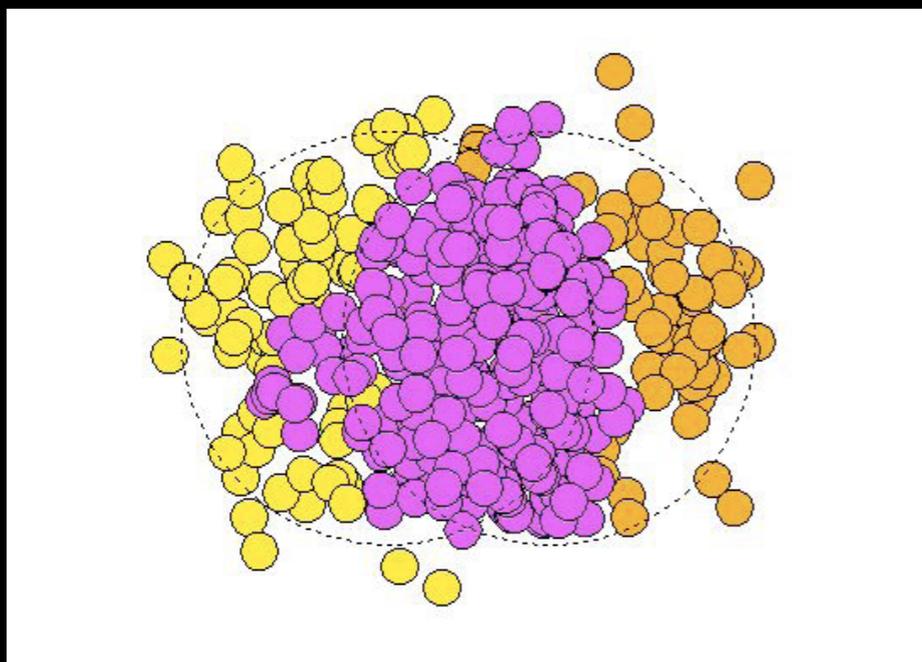
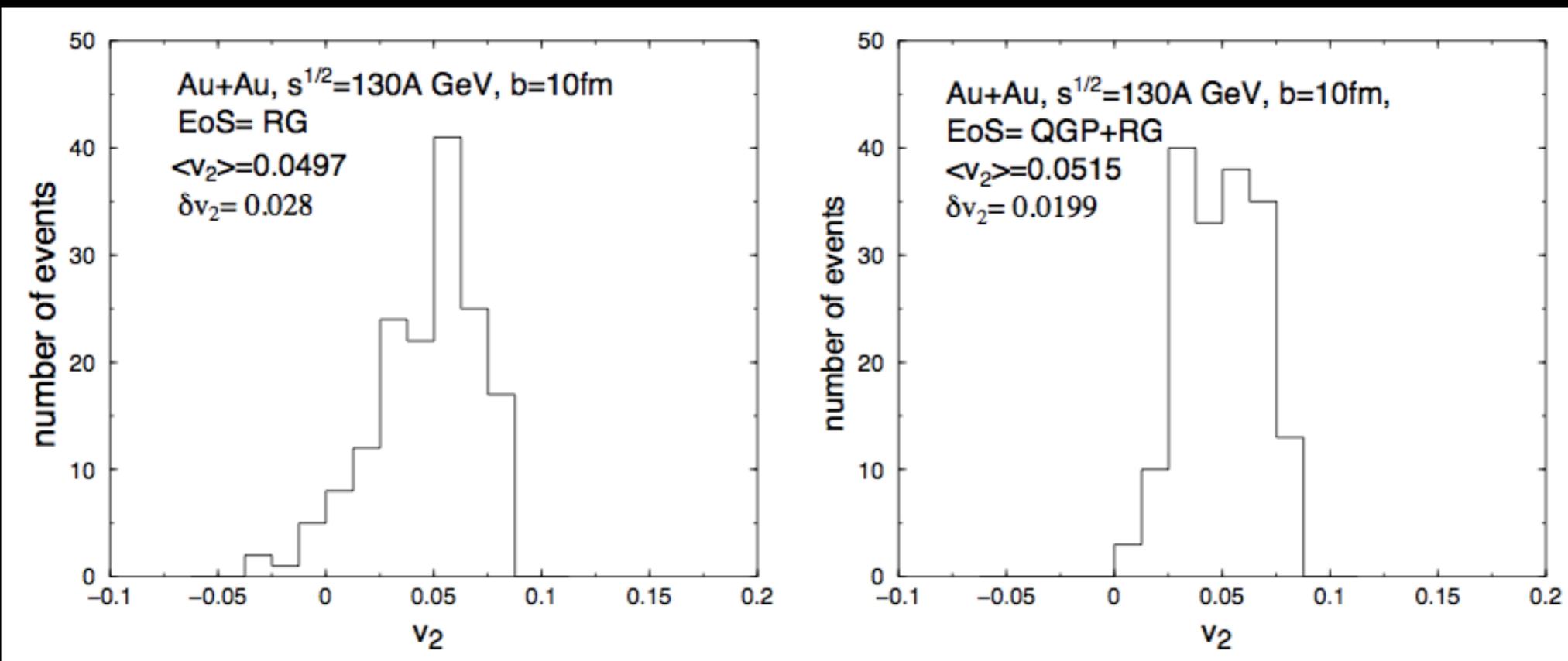


2002

2006

2010

2014



Aguiar, Hama et al  
 Nucl.Phys.A698 (2002) 639-642

Using NeXus (K.Werner) MC  
 Glauber initial conditions

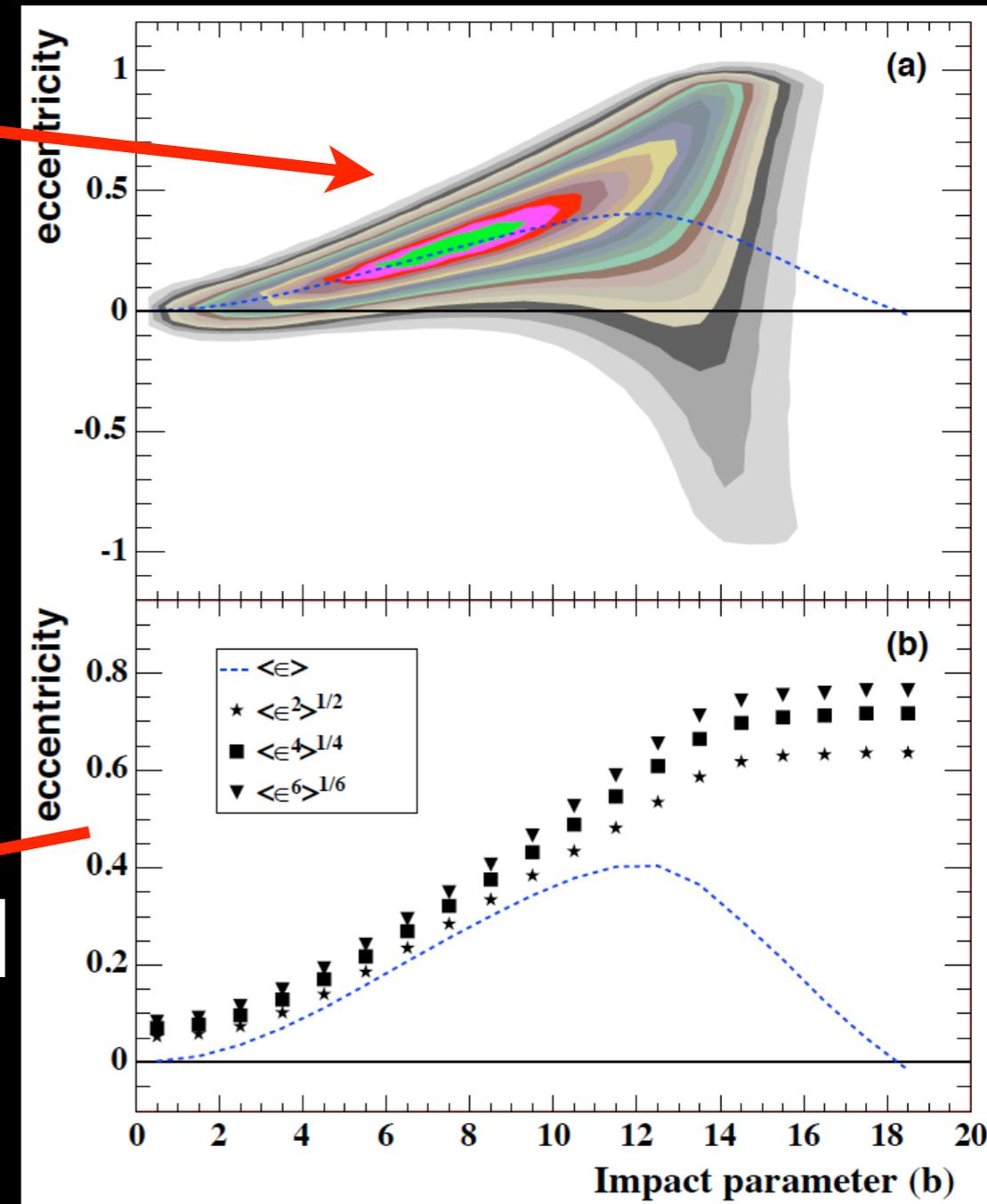


# Importance of Geometry Fluctuations



## MC Glauber eccentricity fluctuations relative to the reaction plane

Miller, Snellings  
nucl-ex/03 | 2008



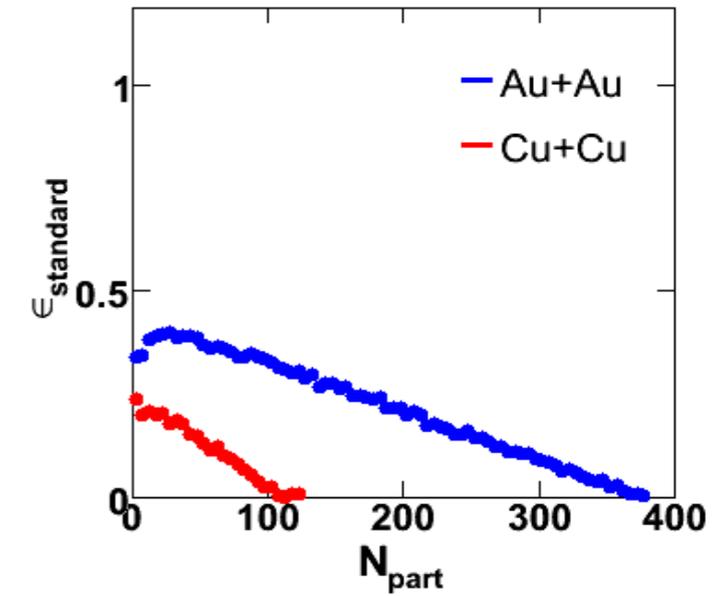
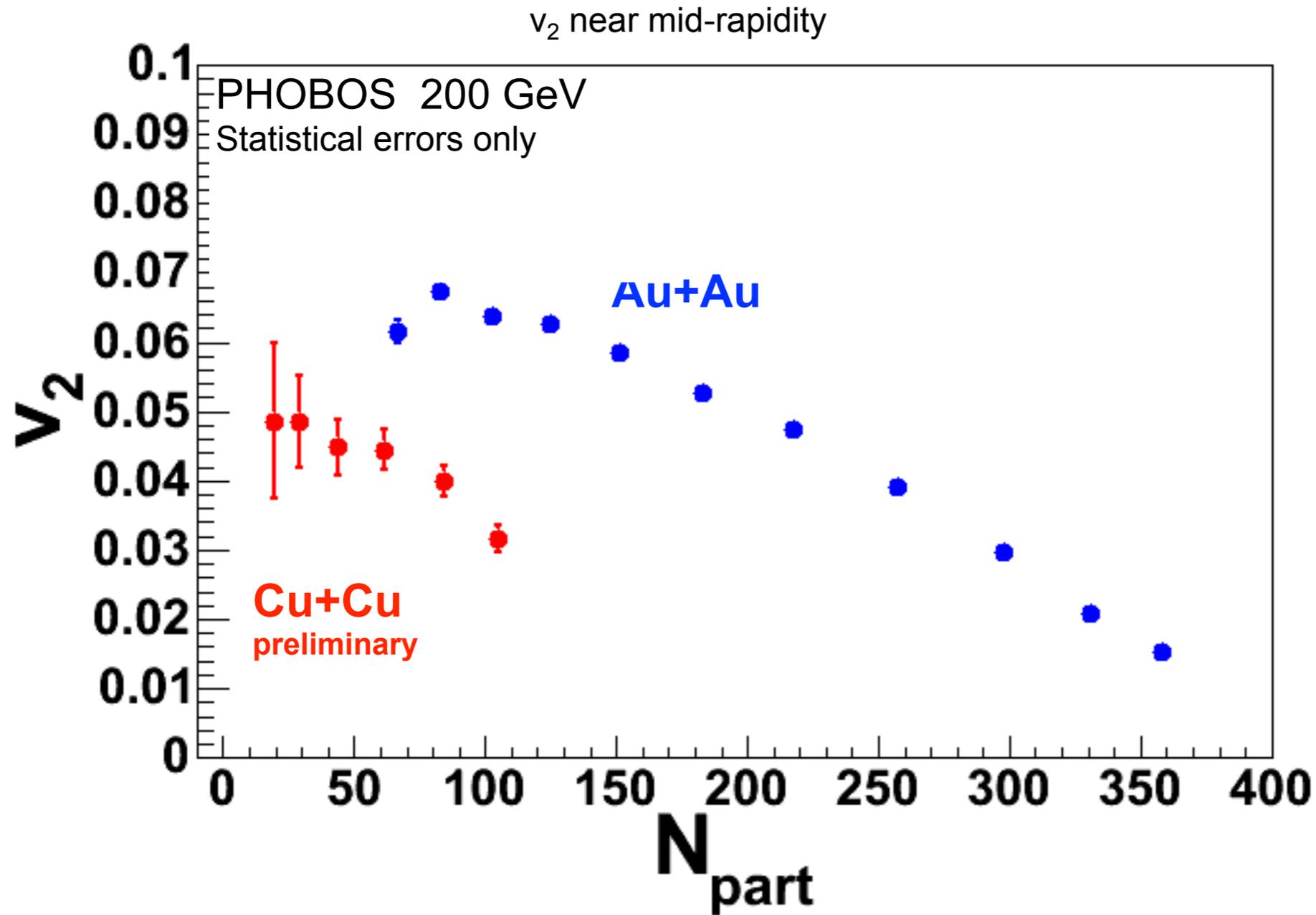
In summary, experimental measurements of elliptic flow ( $v_2$ ) might be affected by fluctuations.

Geometry fluctuations still treated as a perturbation of reaction plane eccentricity

## Elliptic Flow, Geometry & Density

**Q: How does elliptic flow scale with geometry and density?**

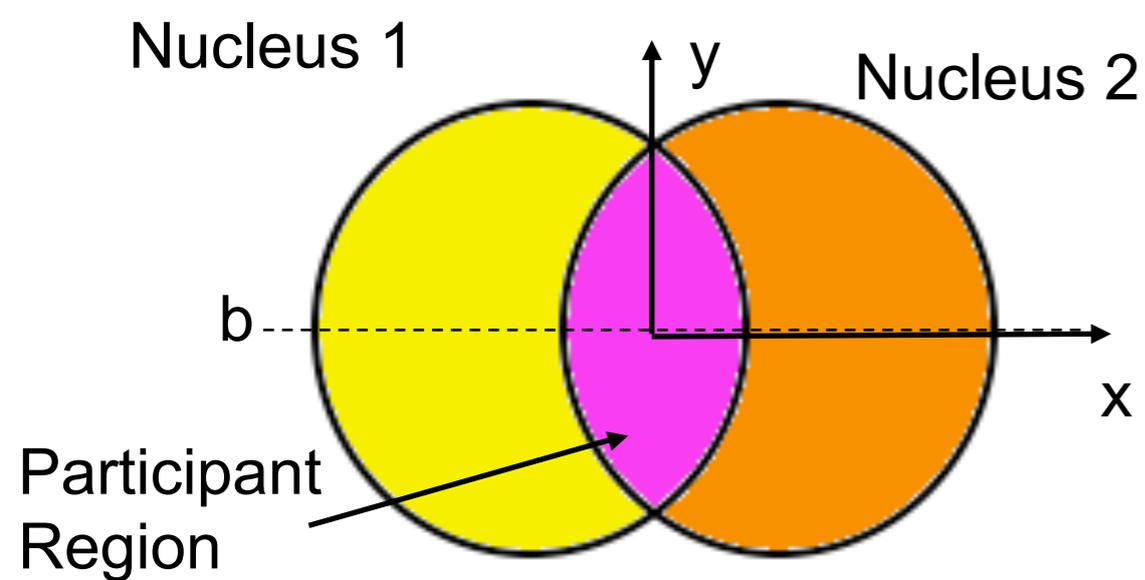
# Elliptic Flow vs $N_{part}$



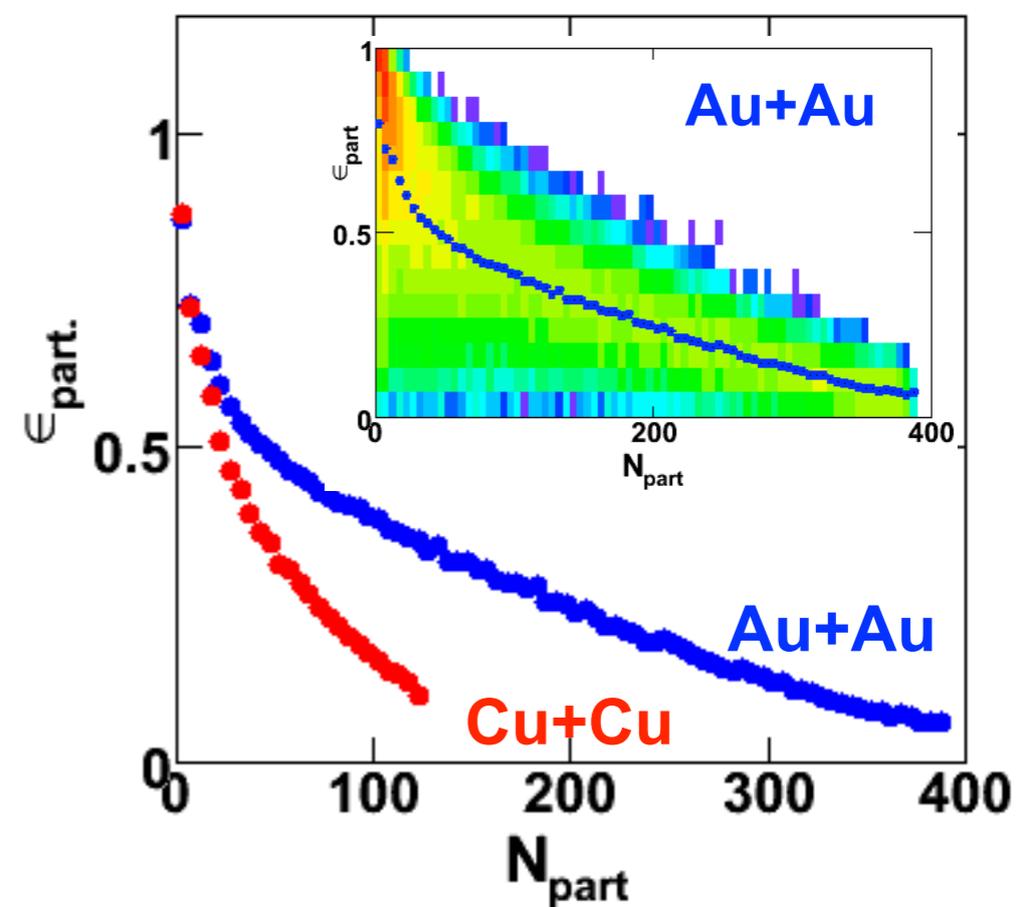
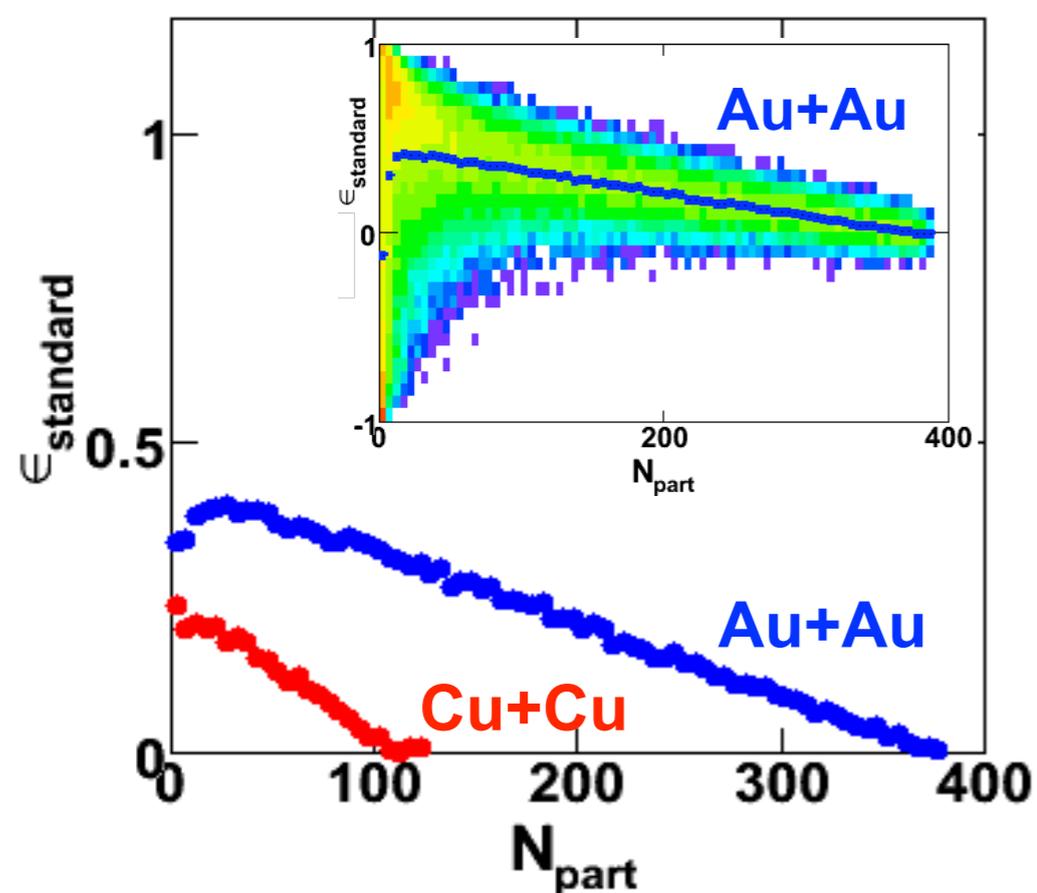
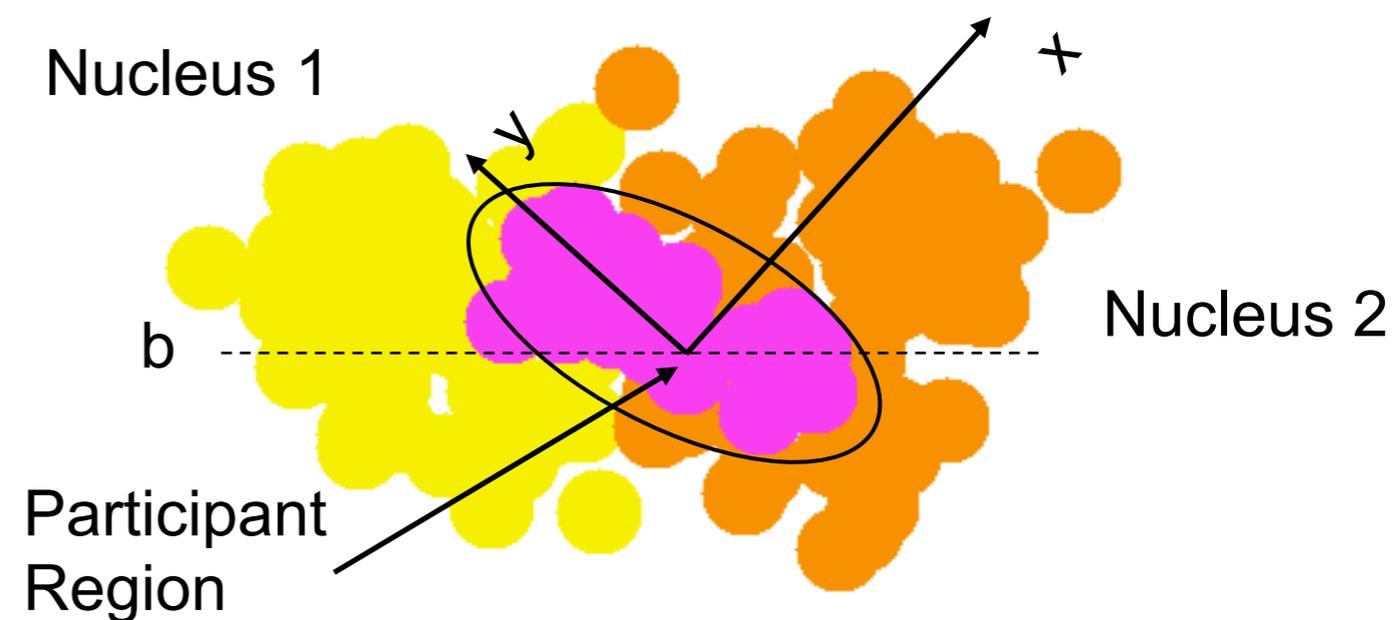
**Substantial  $v_2$  even for most central bin in Cu+Cu**

# Eccentricity Calculation

## Standard Eccentricity

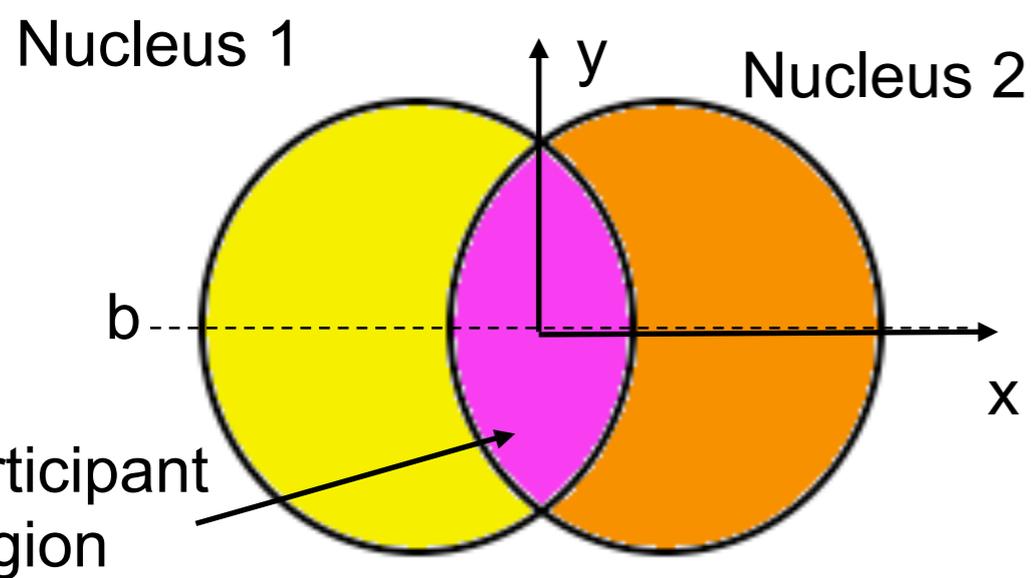


## Participant Eccentricity

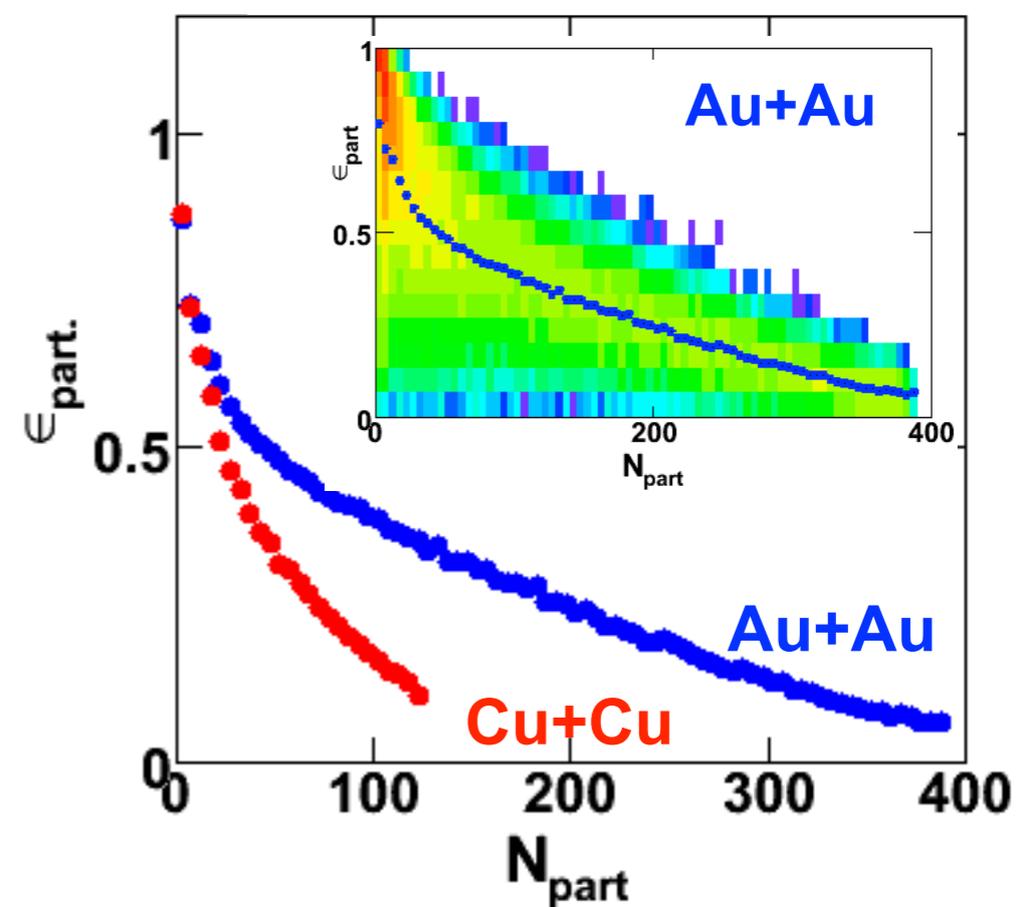
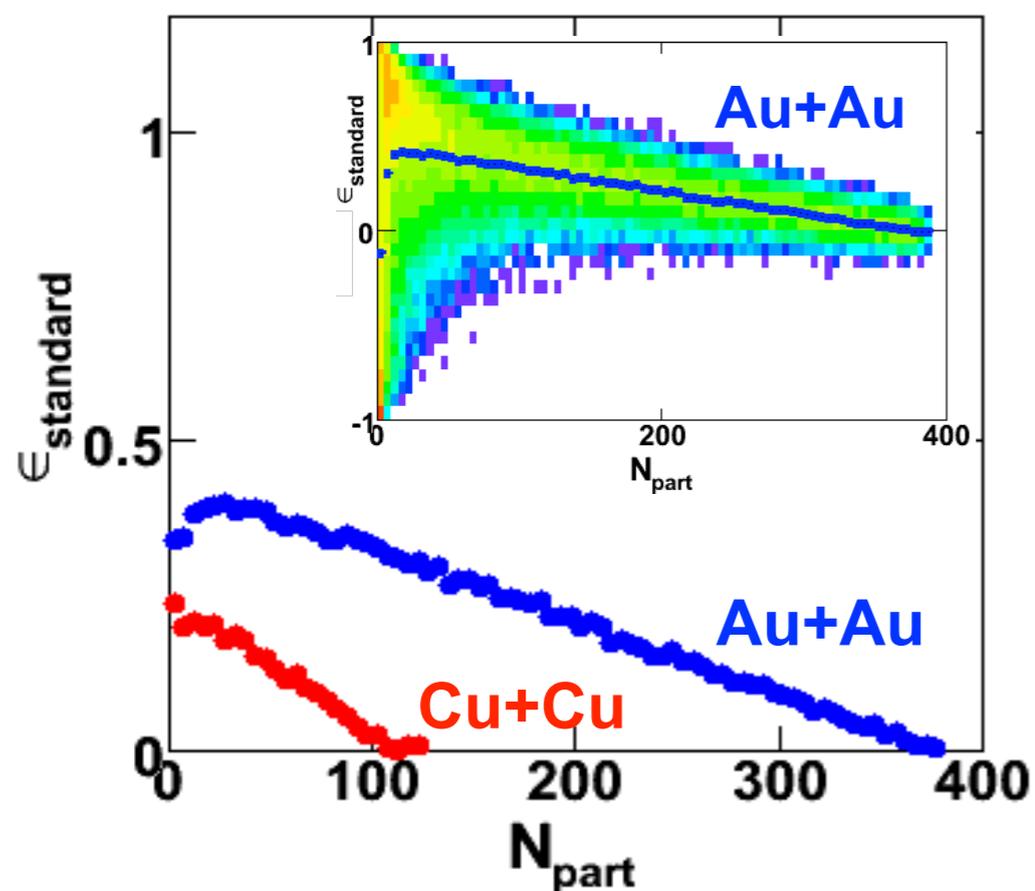
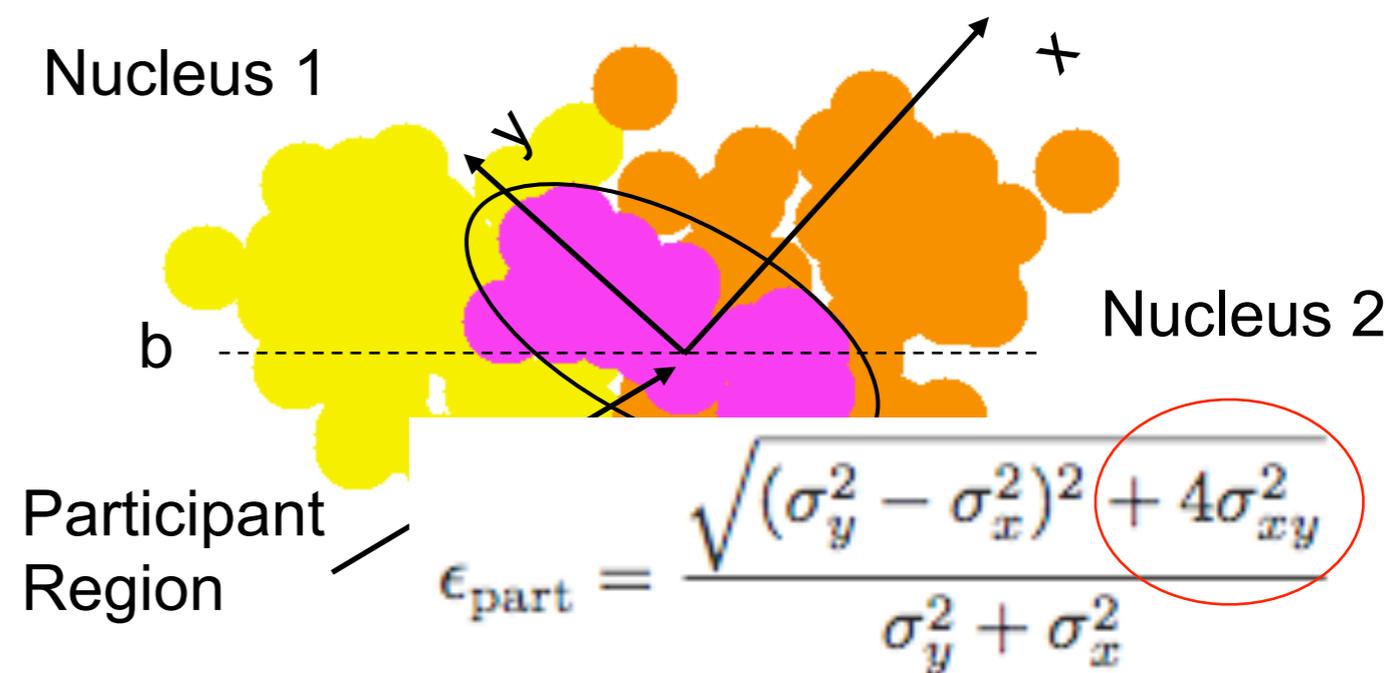


# Eccentricity Calculation

## Standard Eccentricity

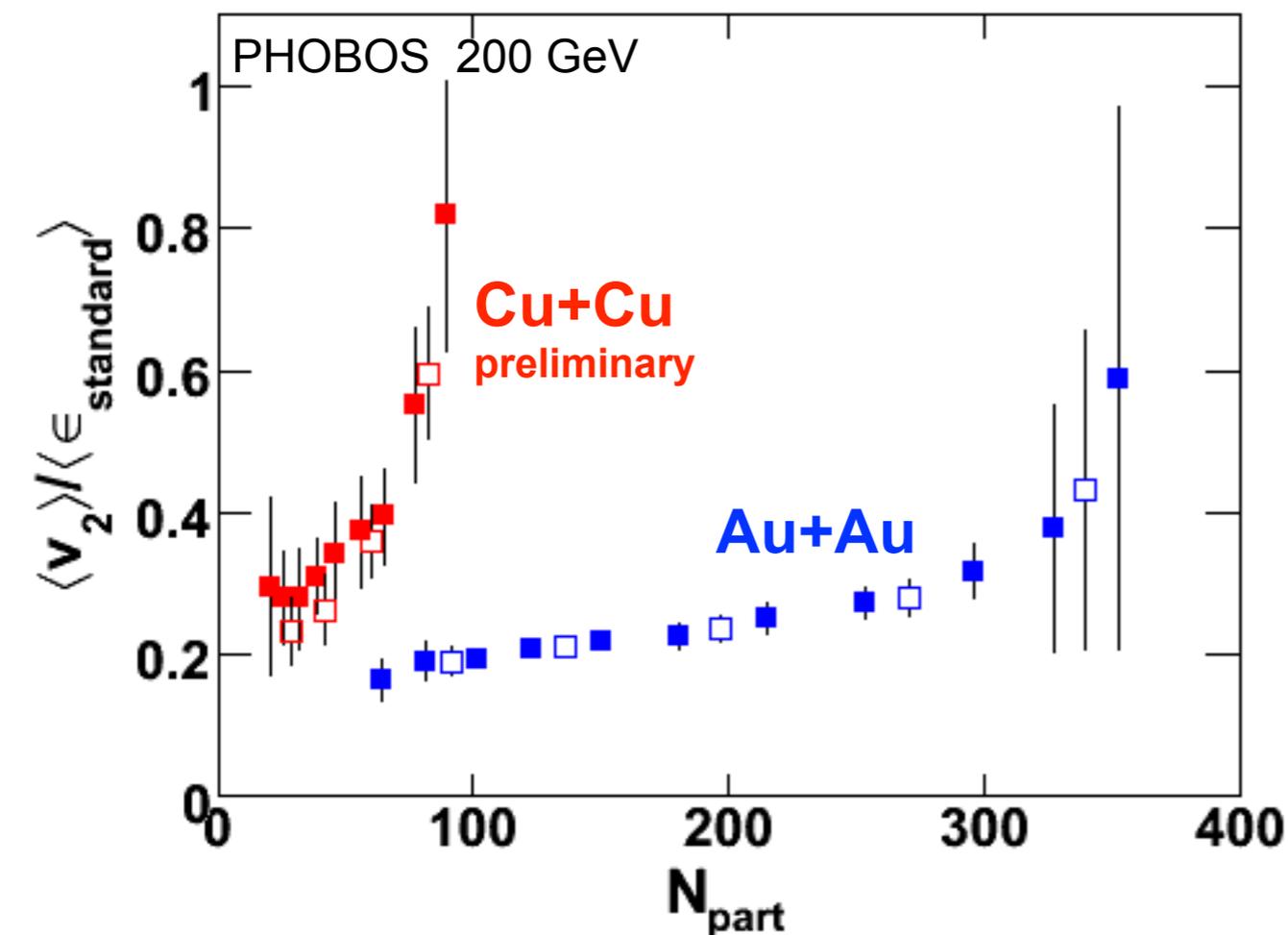


## Participant Eccentricity

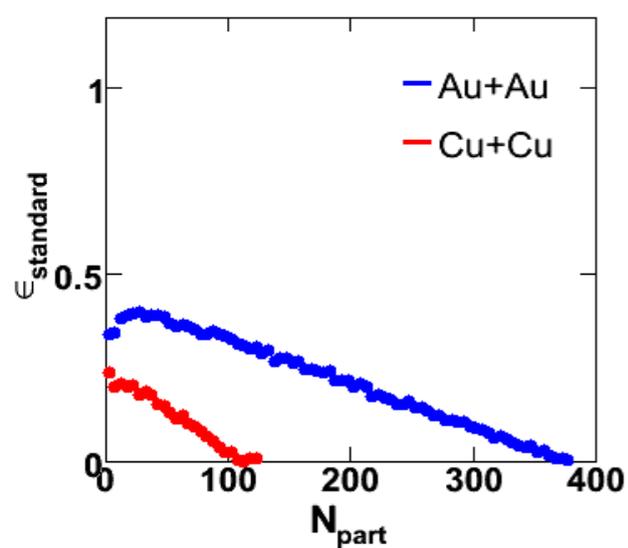
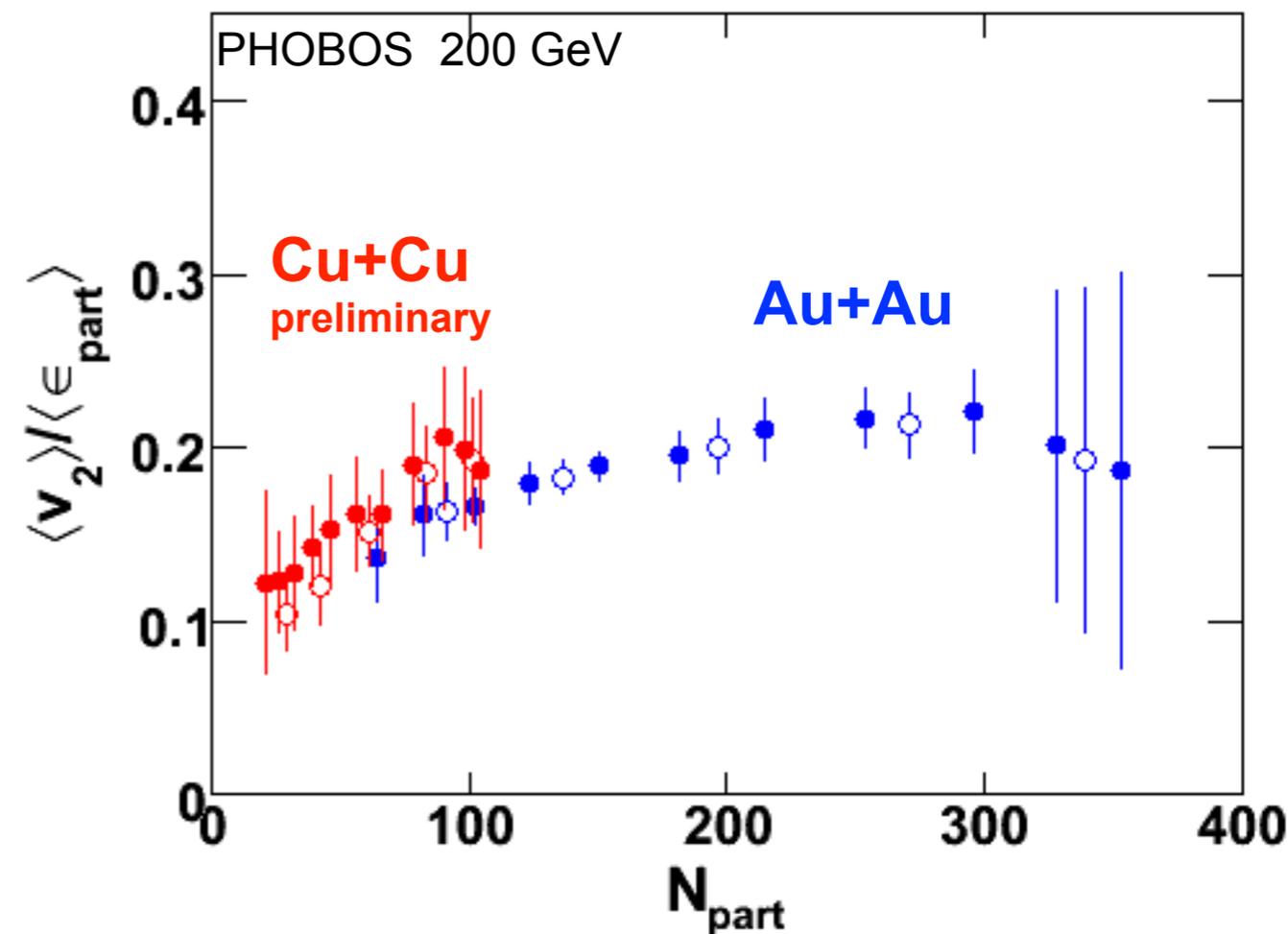


# Elliptic Flow vs $N_{part}$ , II

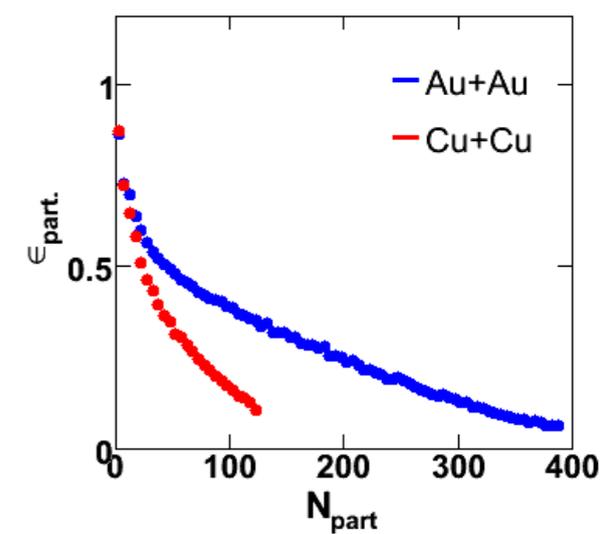
## Standard Eccentricity



## Participant Eccentricity

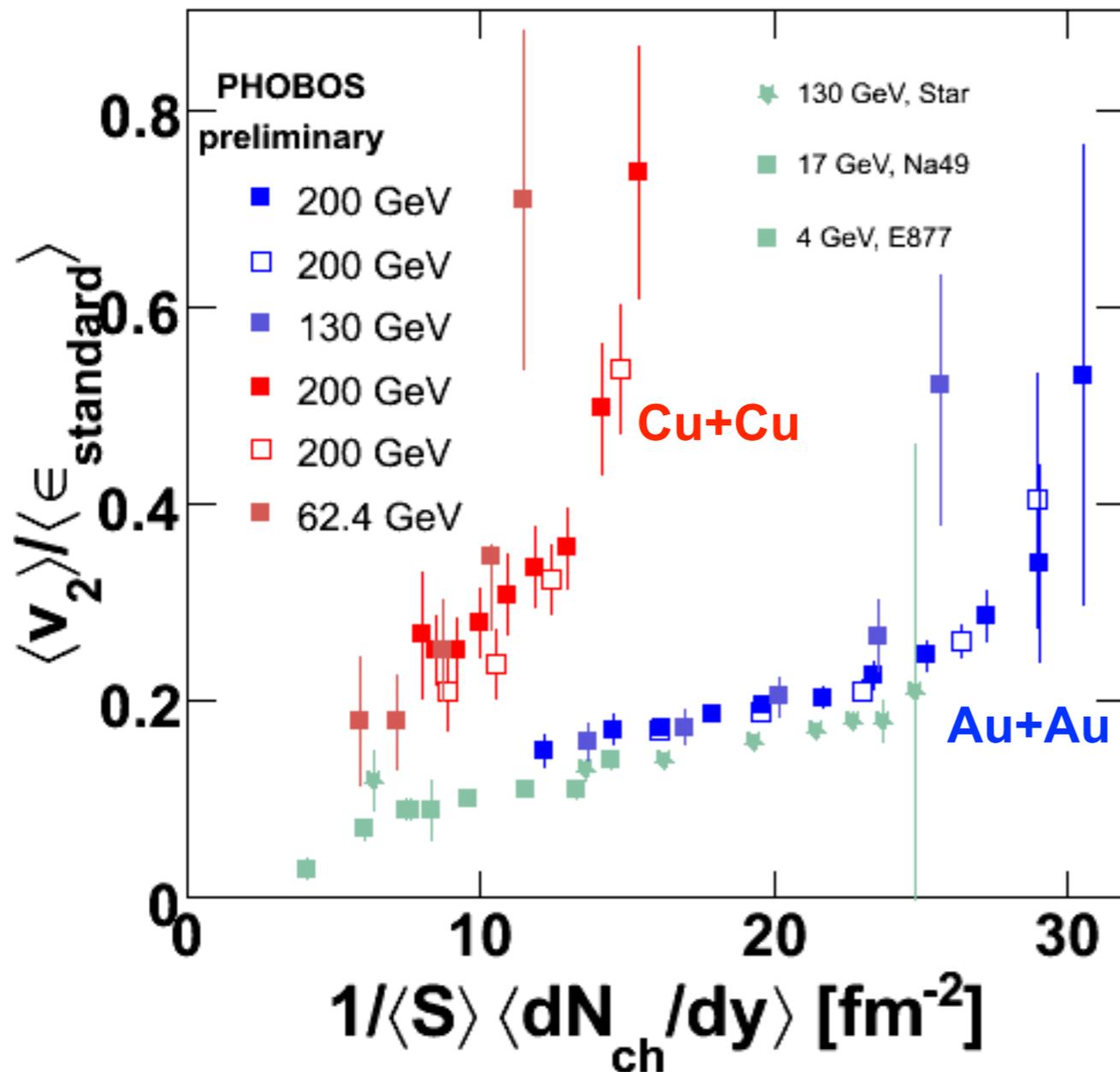


“Participant Eccentricity”  
allows  $v_2$  scaling from  
Cu+Cu to Au+Au

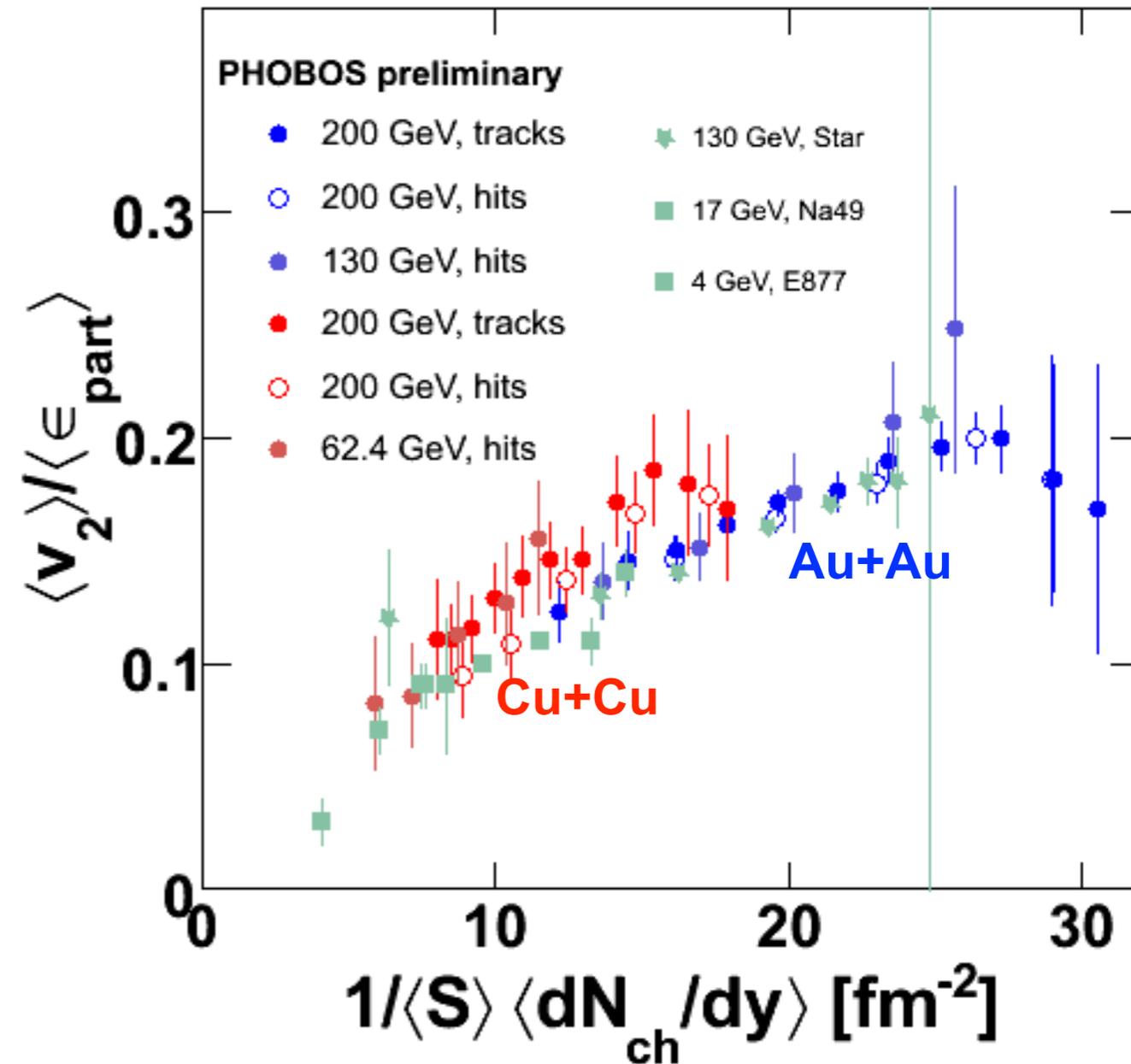


# 'Low Density Limit'-Scaling

## Standard Eccentricity



## Participant Eccentricity



Low Density Limit:

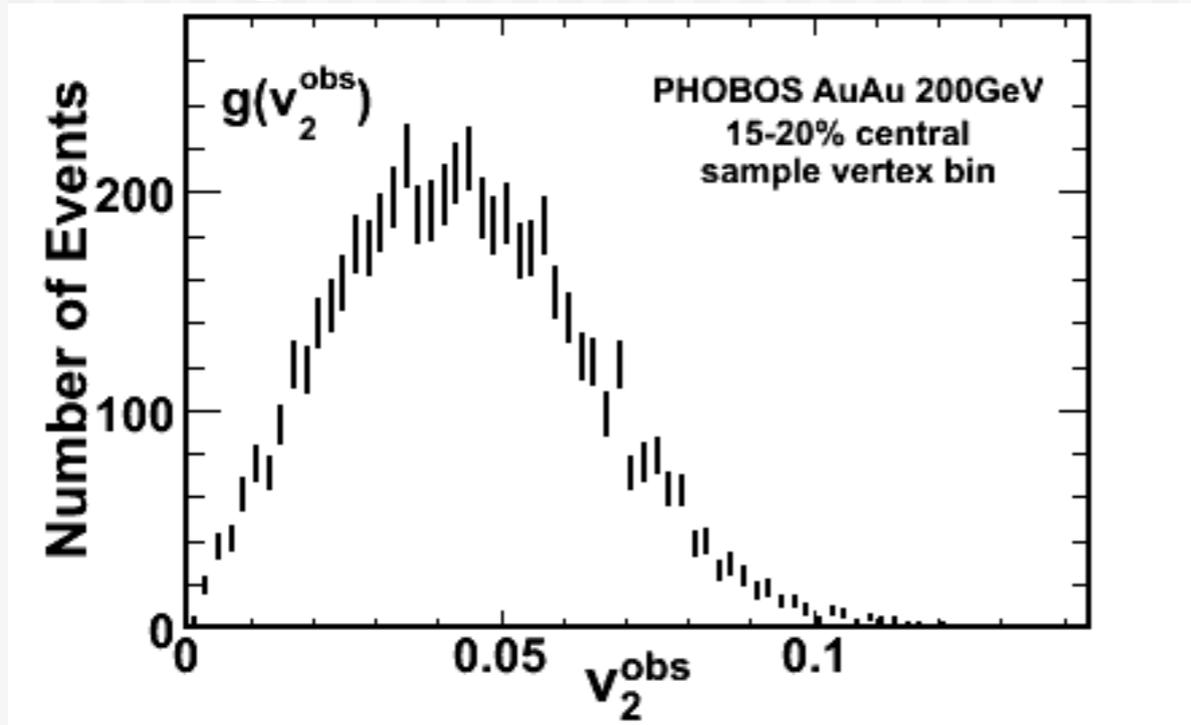
STAR, PRC 66 034904 (2002)

Voloshin, Poskanzer, PLB 474 27 (2000)

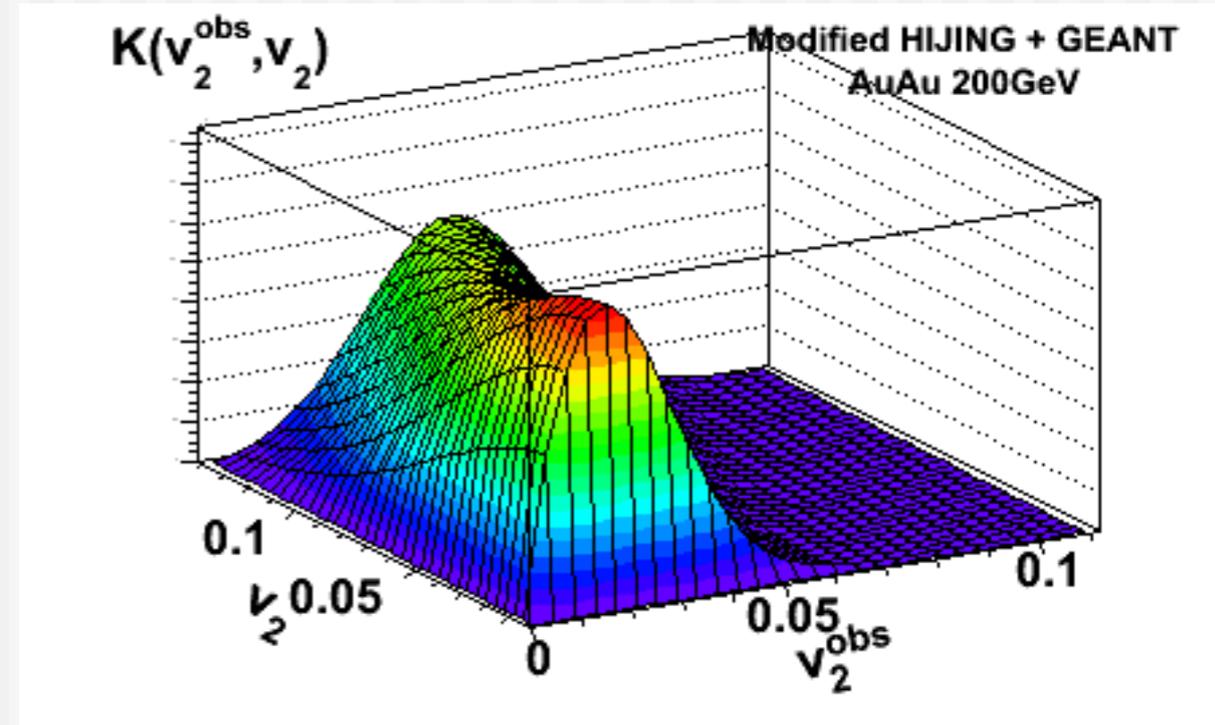
Heiselberg, Levy, PRC 59 2716, (1999)

# Quark Matter 2006

$v_2^{\text{obs}}$  distribution in "data"



Kernel – Response Function

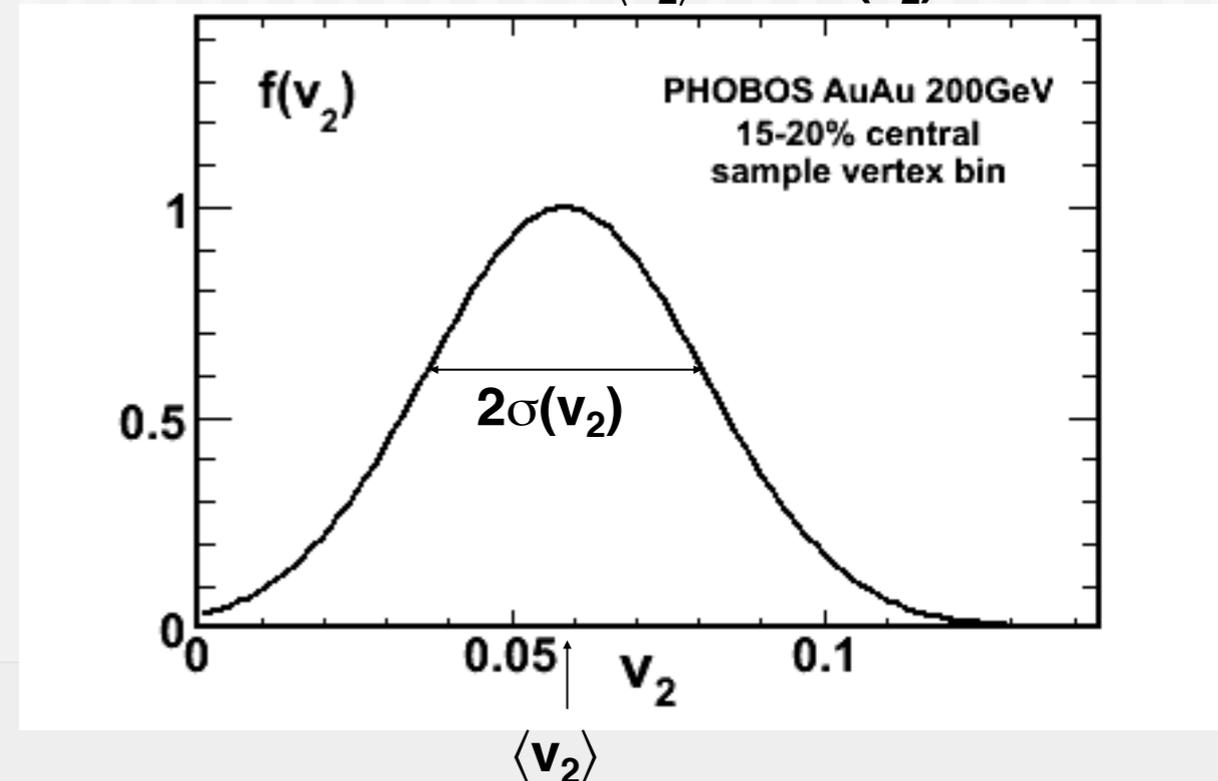


- Event-by-event measurement
- Determination of response in MC
- Extraction of true  $\langle v_2 \rangle$  and  $\sigma(v_2)$

$$g(v_2^{\text{obs}}) = \int K(v_2^{\text{obs}}, v_2) f(v_2) dv_2$$

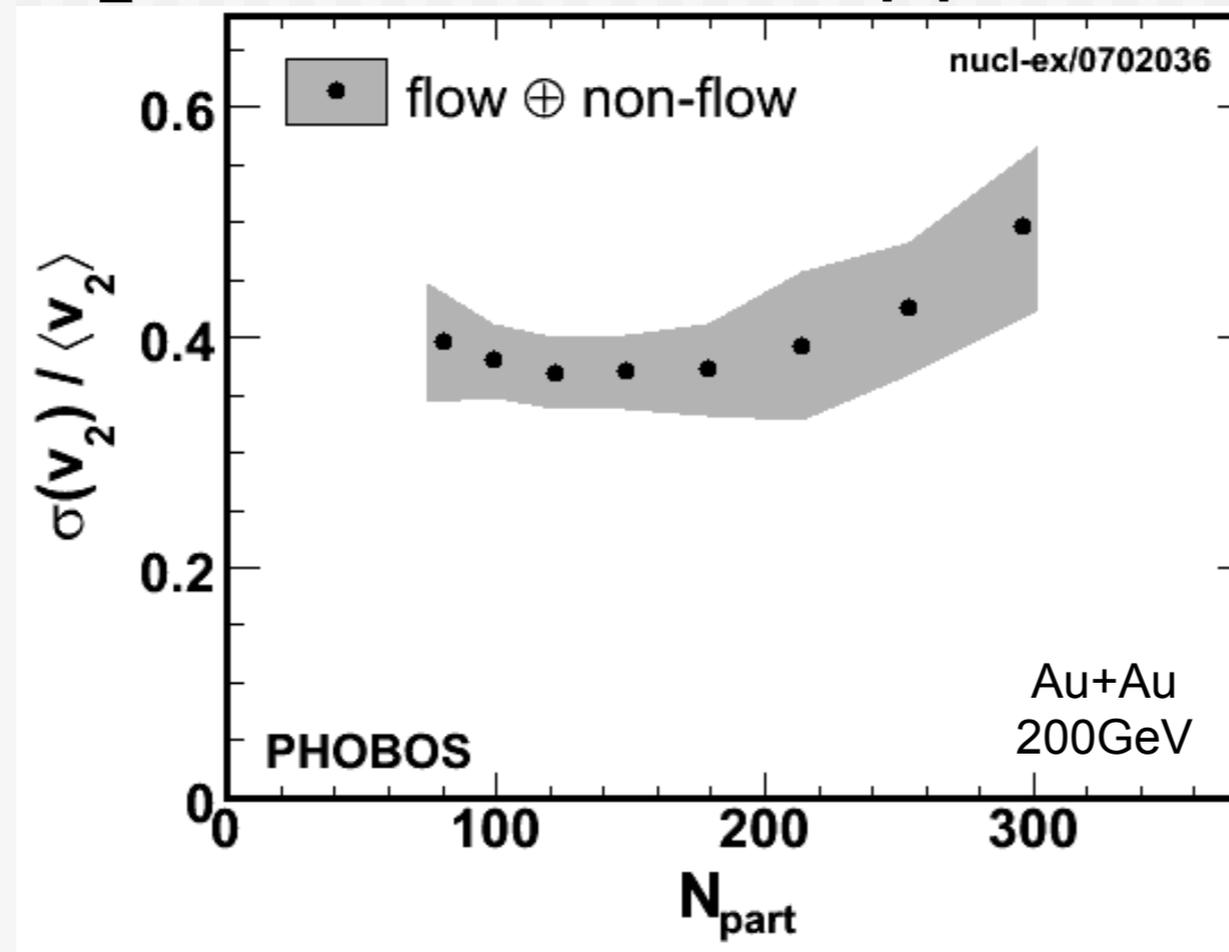
arXiv:nucl-ex/0702036

Extracted true  $\langle v_2 \rangle$  and  $\sigma(v_2)$



# Quark Matter 2006

- Relative  $v_2$  fluctuations of approximately 40%



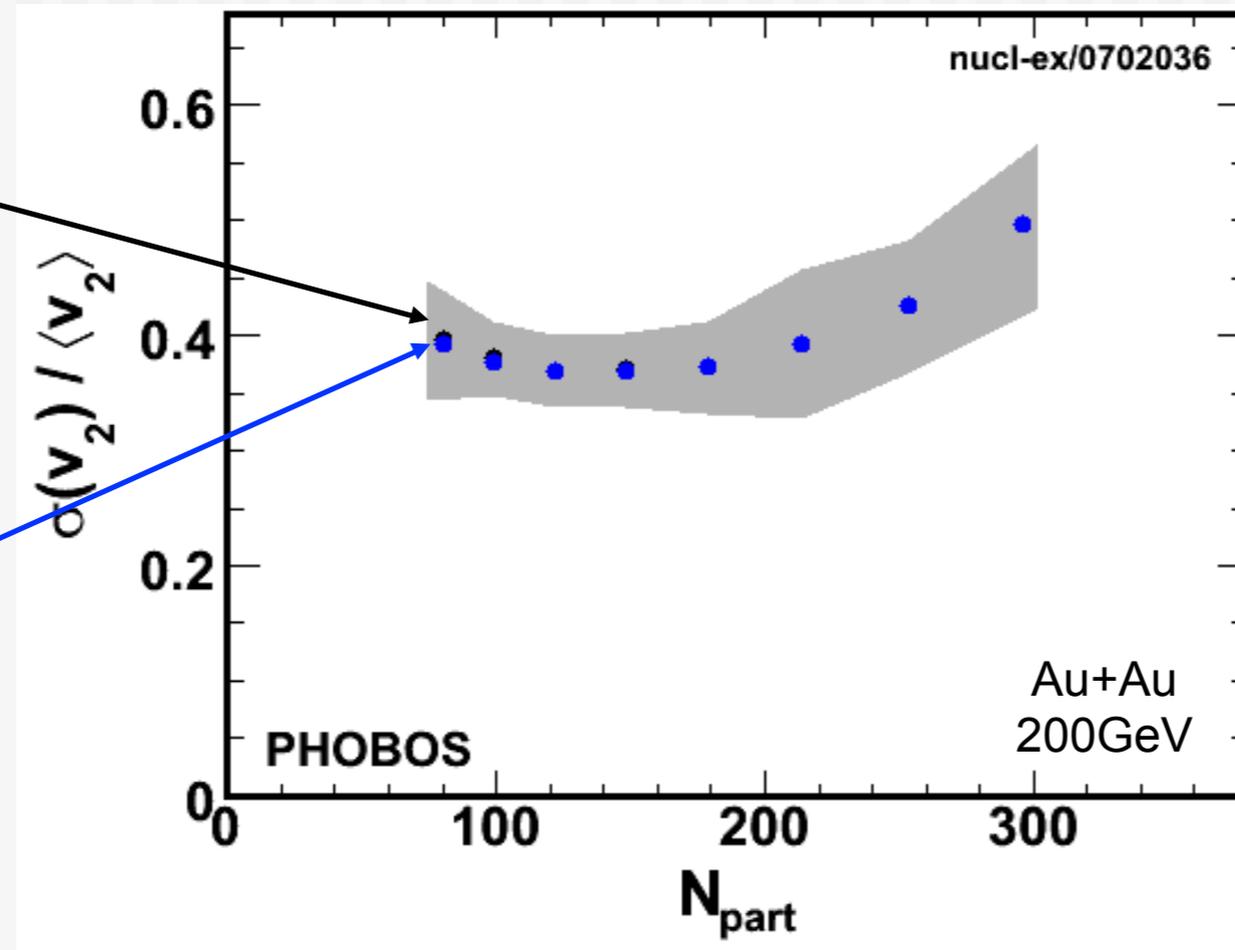
- Correlated particle production (non-flow correlations) can broaden the  $v_2^{\text{obs}}$  distribution and affect the fluctuation measurement.

# Quark Matter 2006

- Estimate non-flow contribution with HIJING

Measured  
fluctuations

HIJING  
Non-flow  
Correction  
(very small)

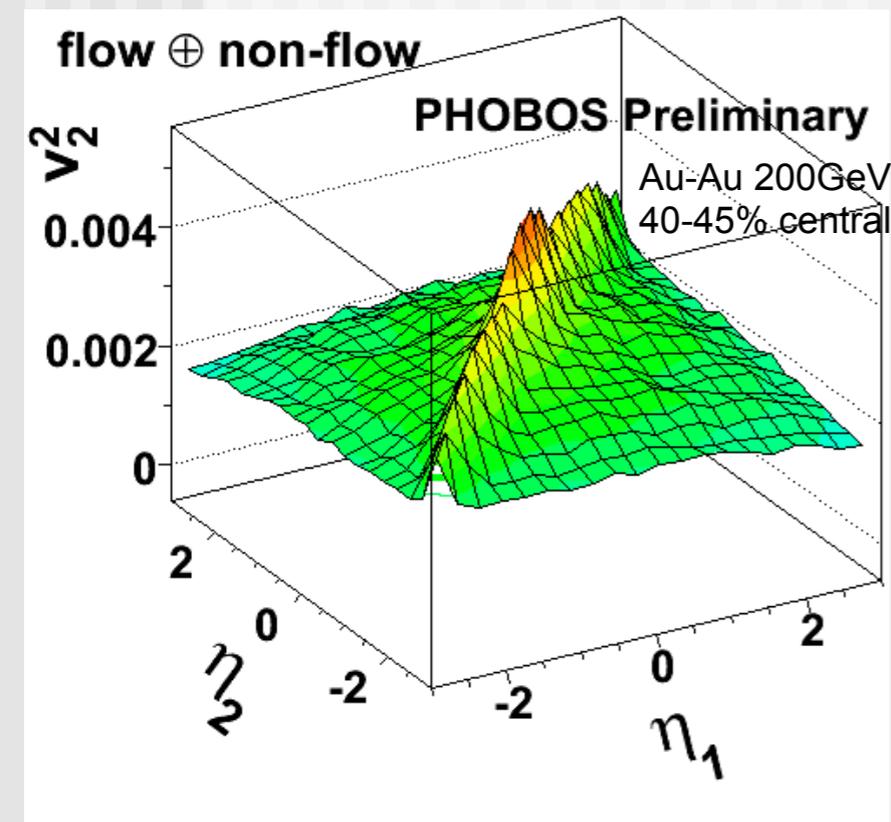


- We used response function calculated from HIJING with correlations preserved to estimate non-flow effect.

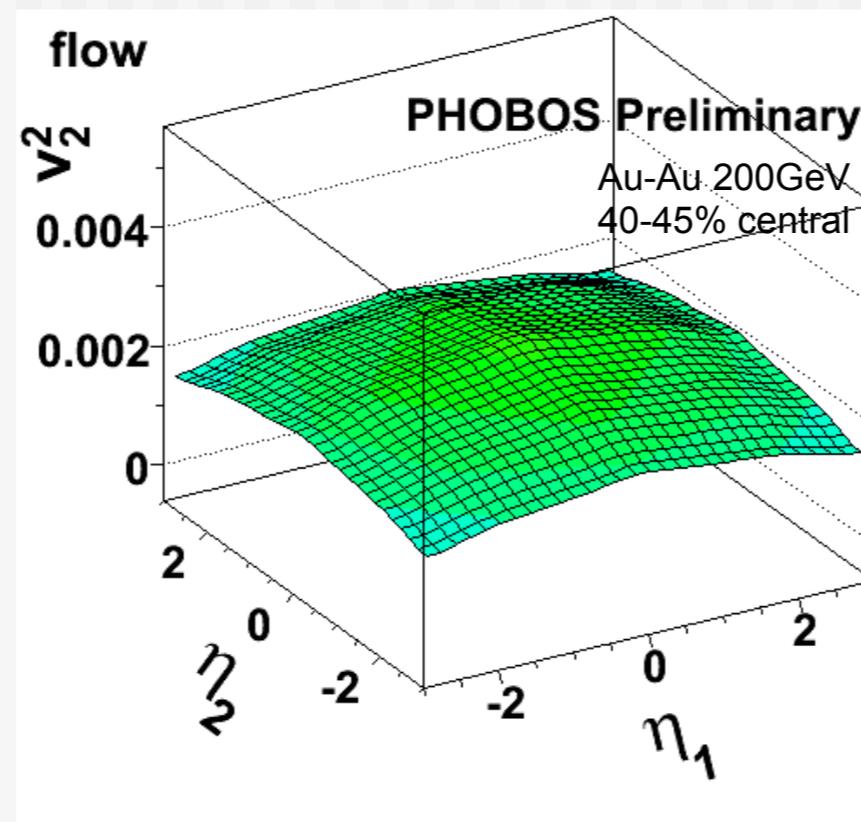
# Separating flow and non-flow

- Subtract to find  $\delta(\eta_1, \eta_2)$  at all ranges:

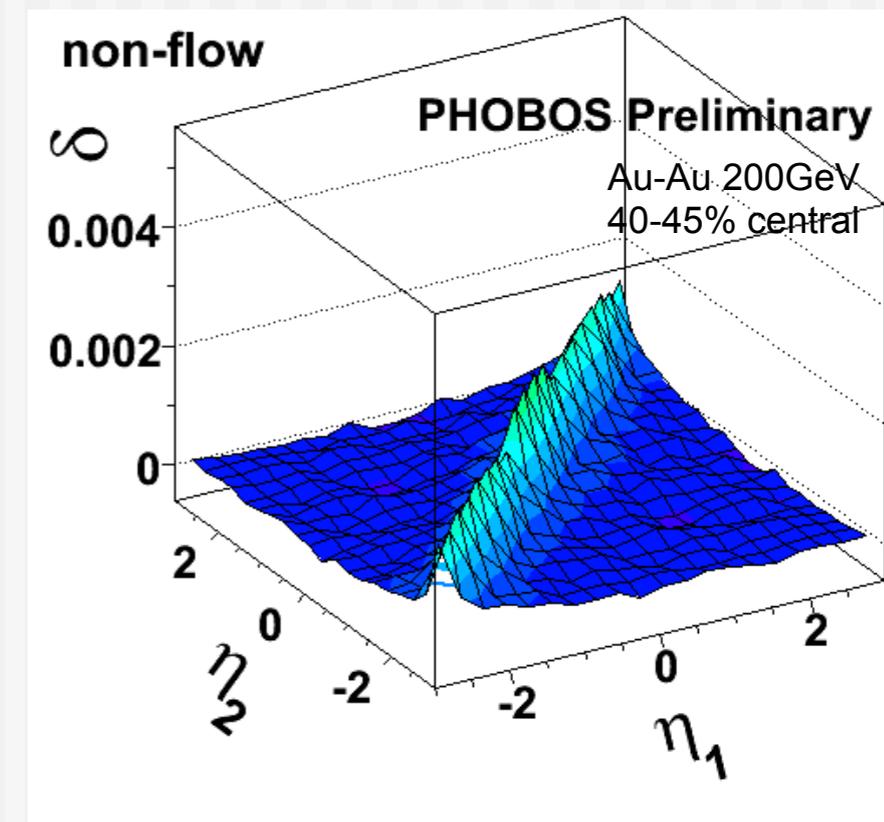
$$\delta(\eta_1, \eta_2) = v_2^2(\eta_1, \eta_2) - v_2(\eta_1) \times v_2(\eta_2)$$



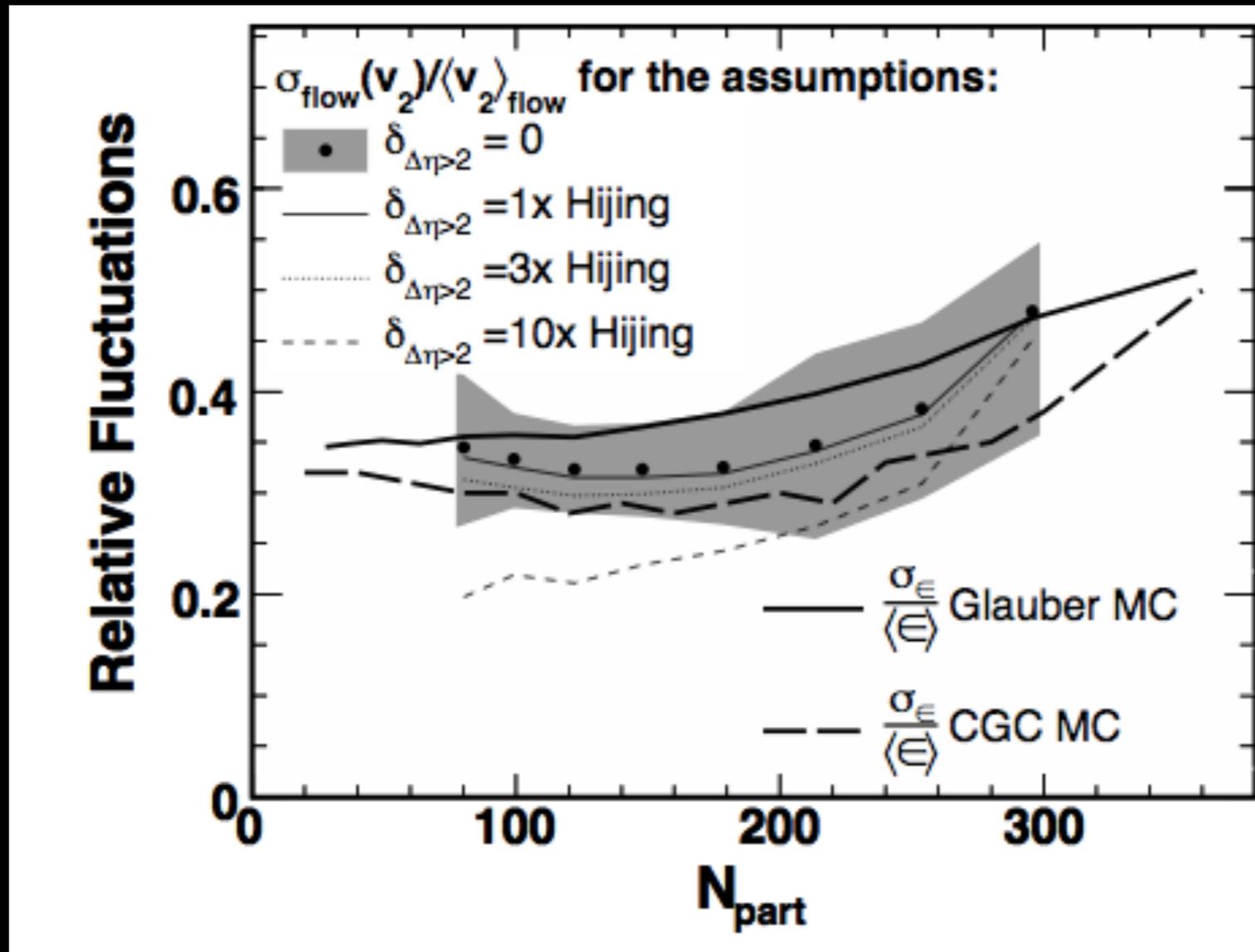
flow $\oplus$ non-flow



flow



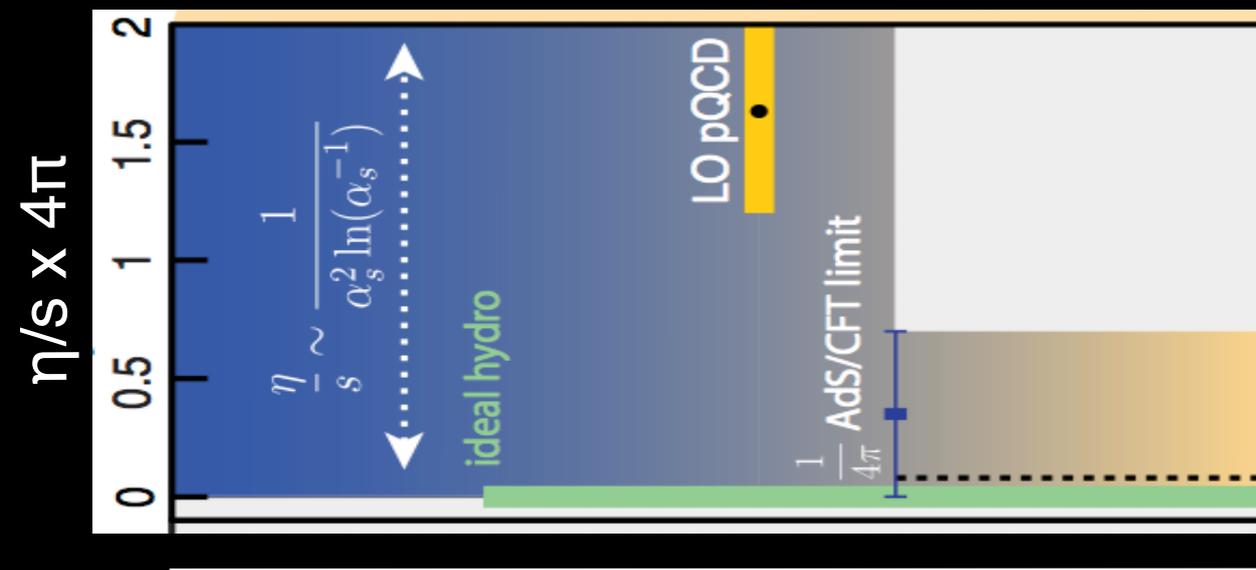
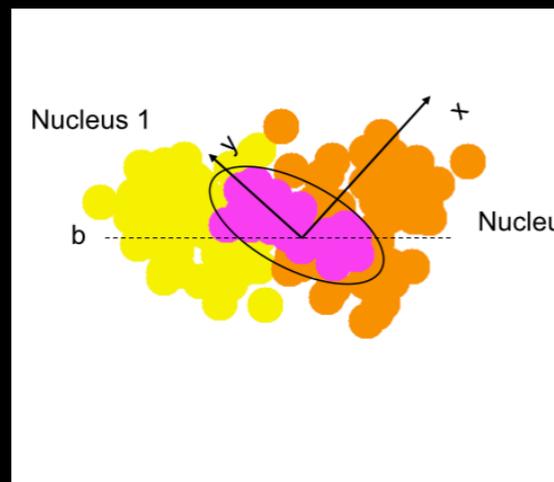
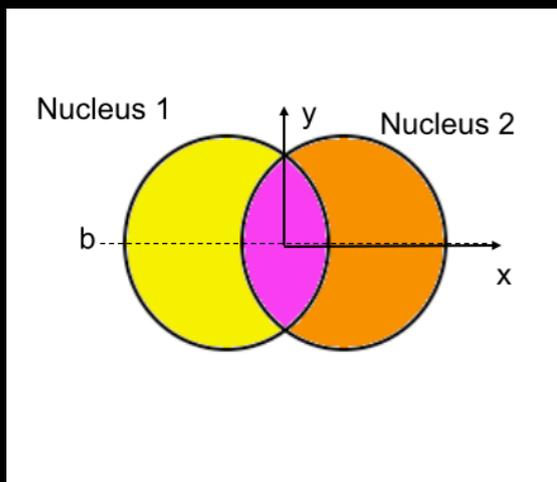
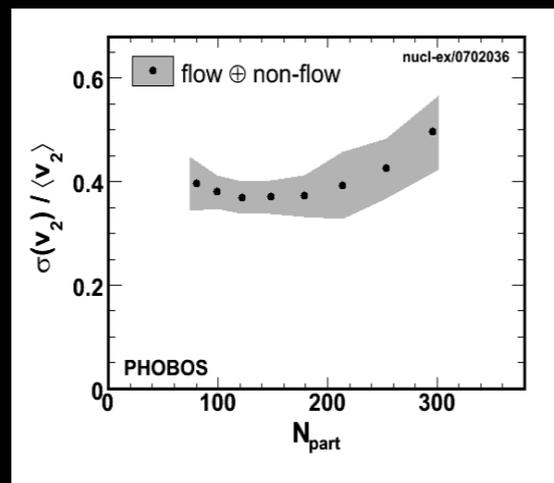
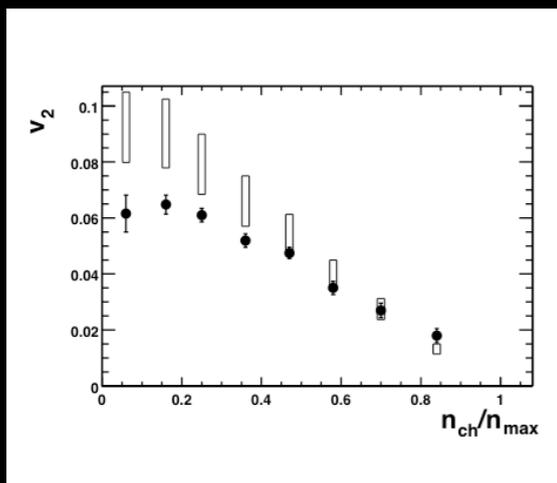
non-flow



Elliptic flow fluctuations for different non-flow assumptions  
(constrained by long-range pseudo-rapidity factorization)



# A brief history of correlations in HIC



2002

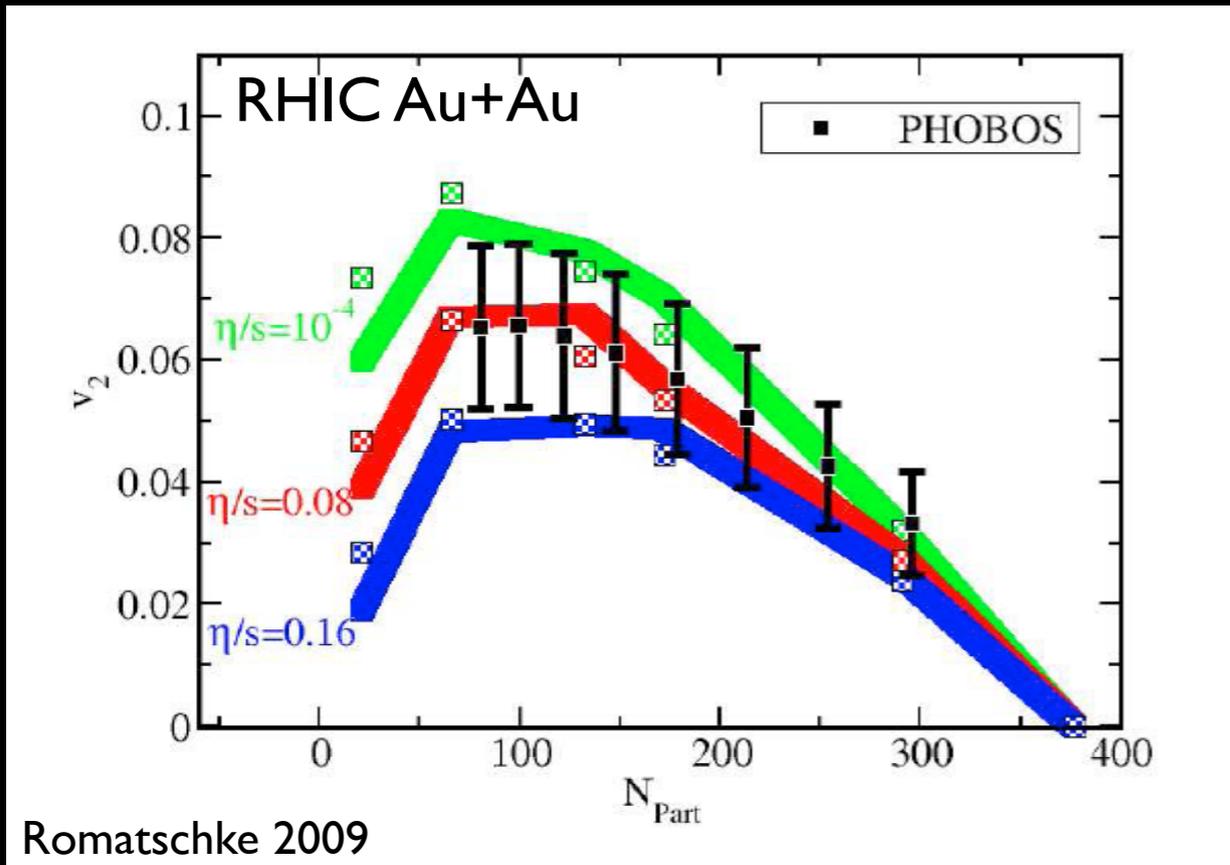
2006

2010

2014

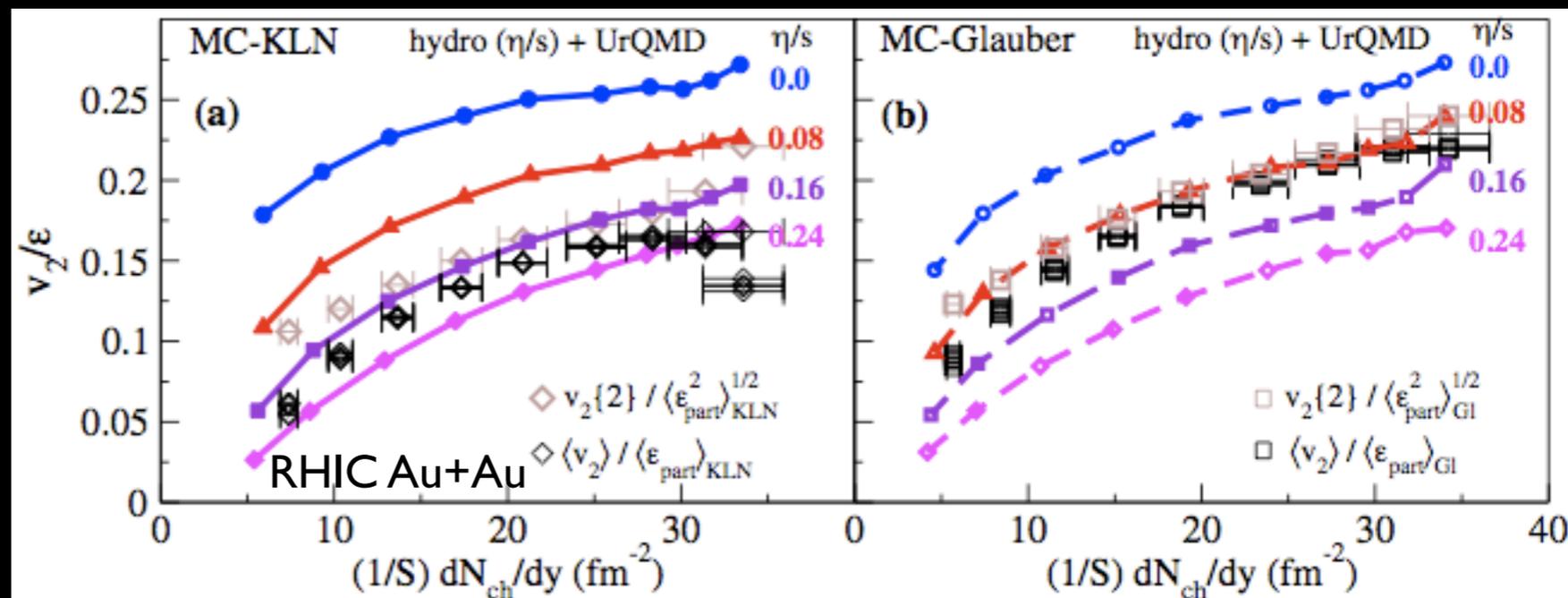


# Development of viscous hydrodynamics

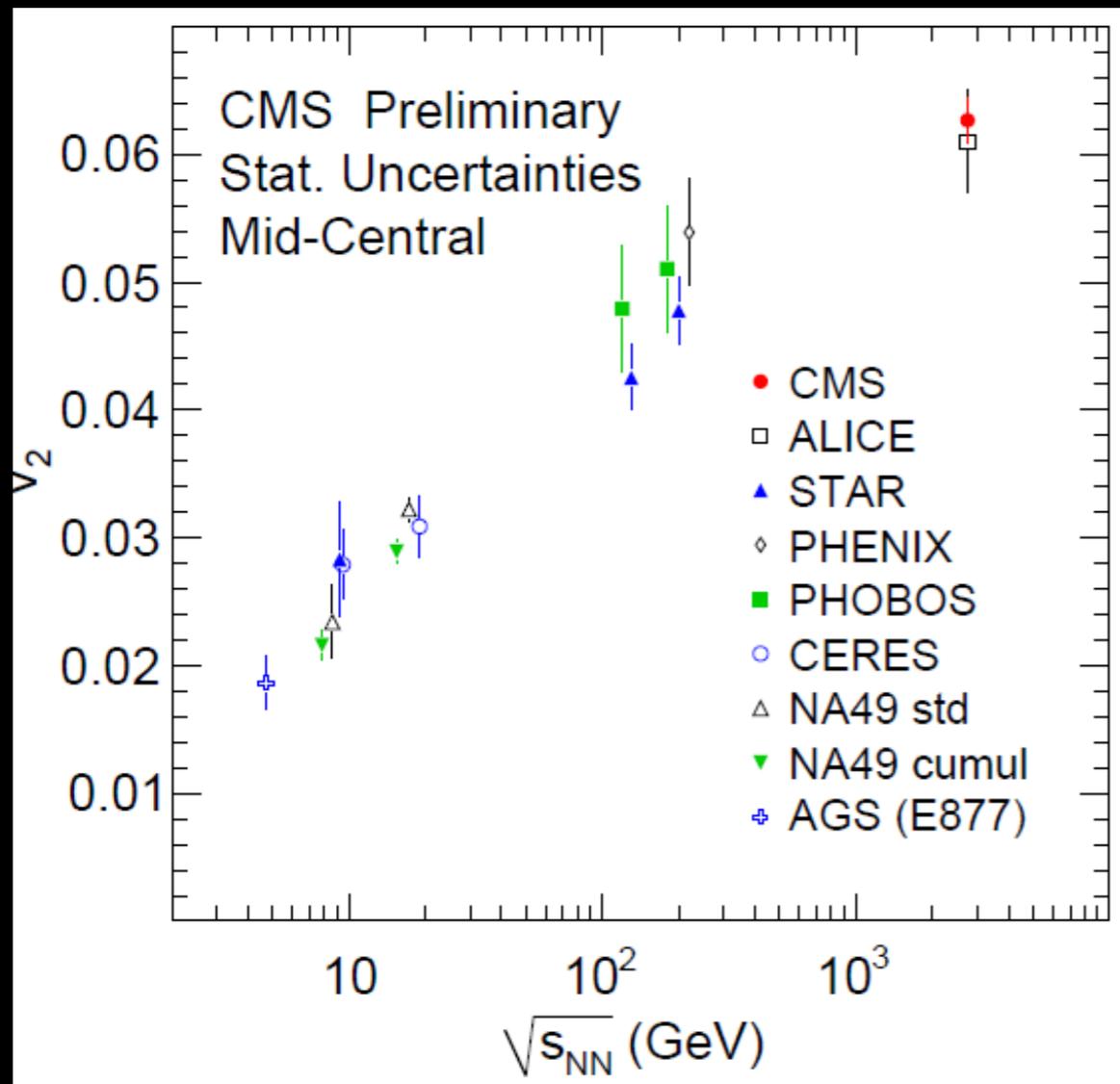


$$0.1 < \eta/s < 0.3$$

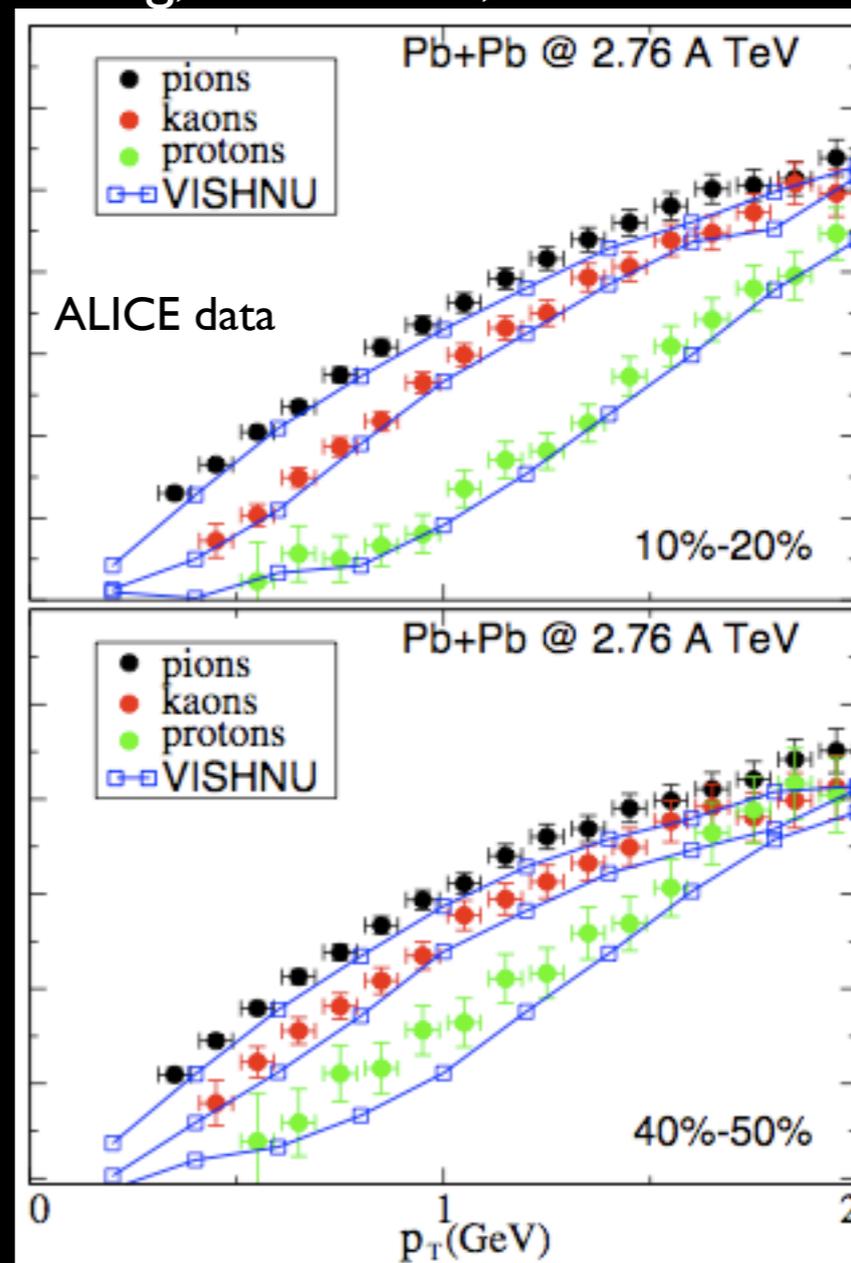
Large contribution to uncertainty from initial geometry



Song, Heinz et al, PANIC 2011



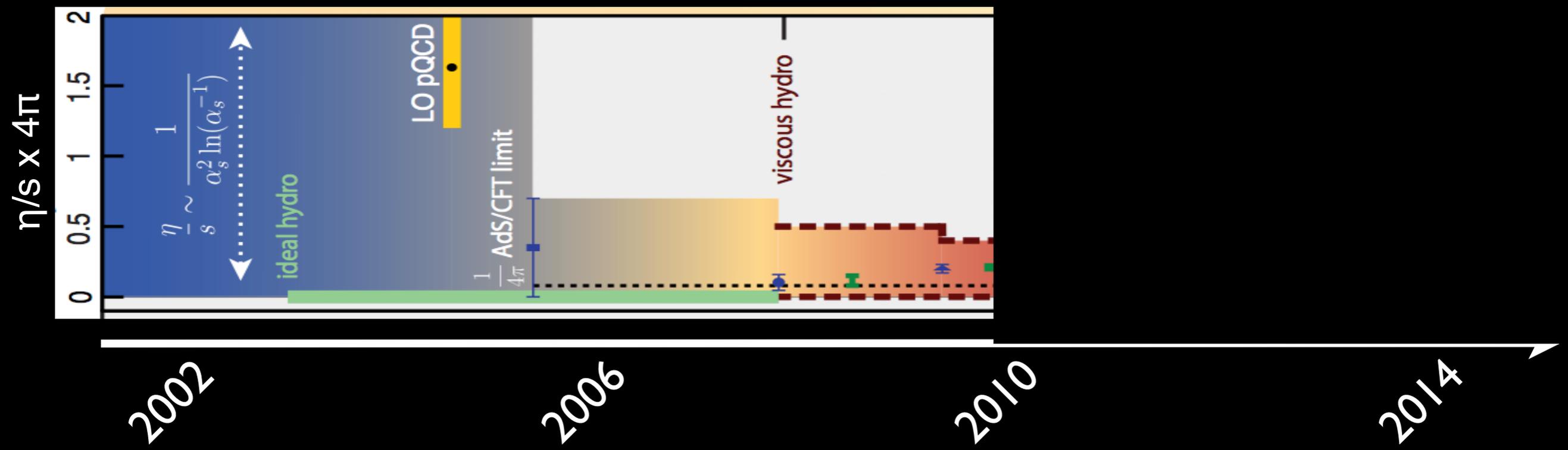
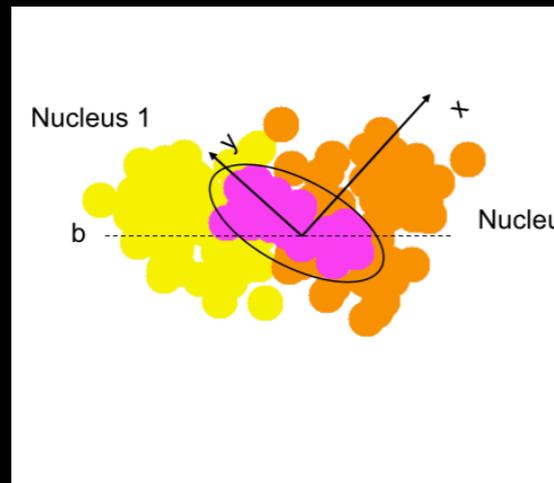
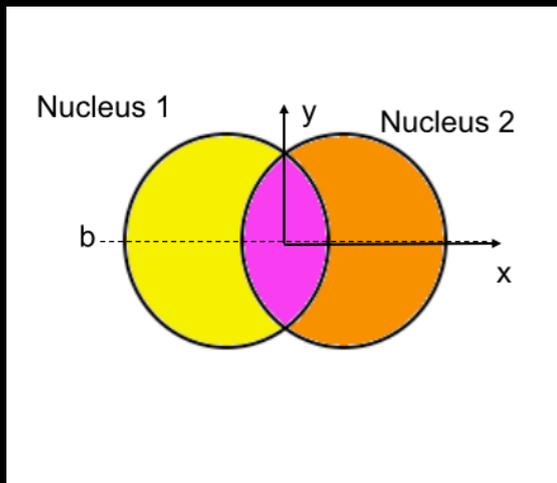
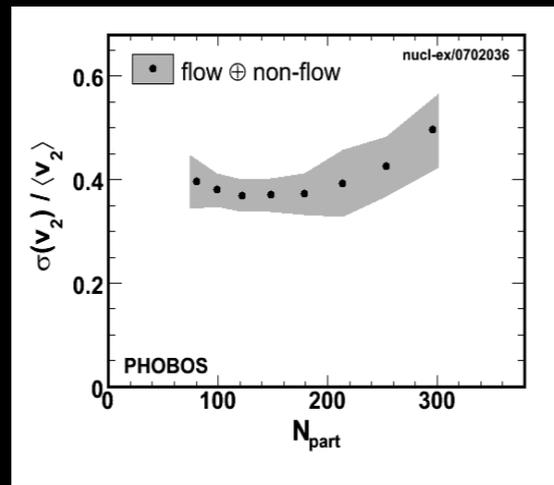
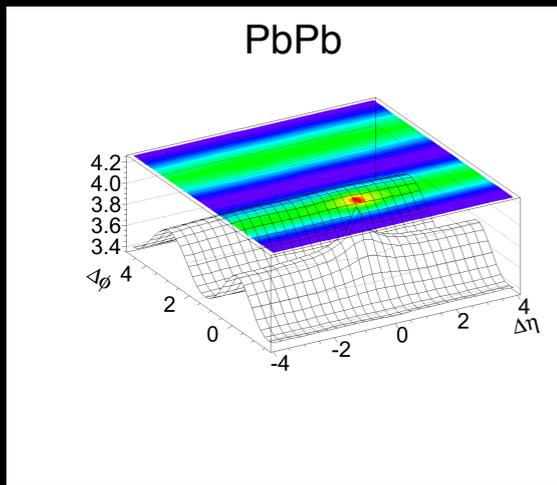
Elliptic flow 10% larger at LHC:  
Stronger initial push due to  
higher density

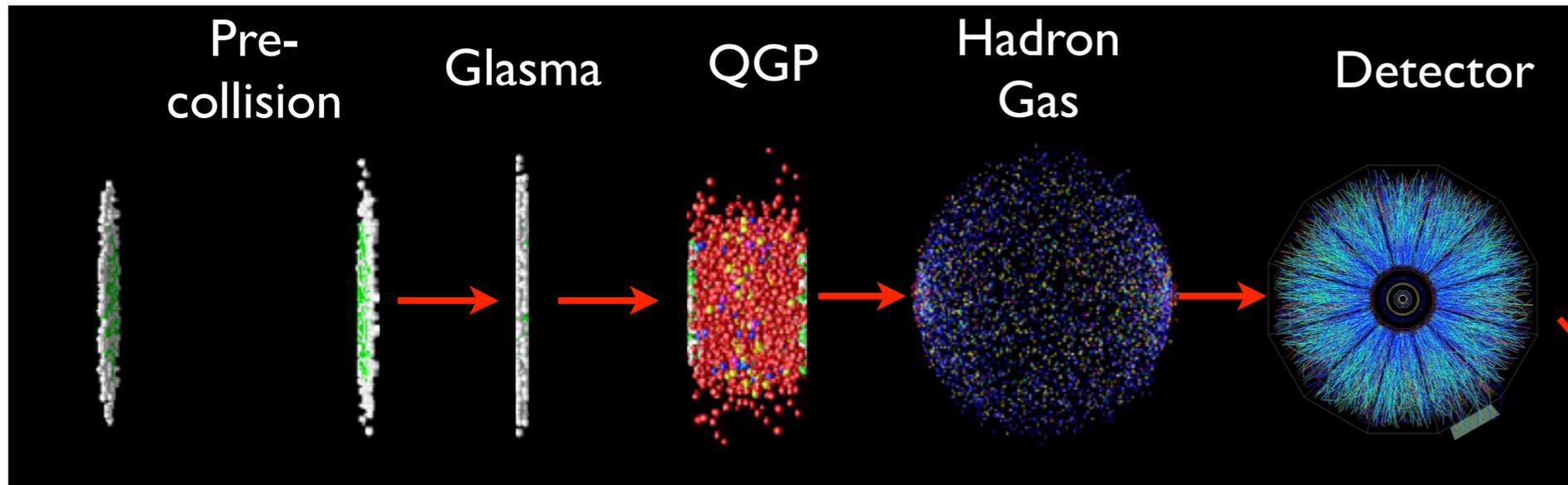


Comparison to state-of-the-art  
hydro calculations suggests:  
 $\eta/s_{\text{(LHC)}} \sim \eta/s_{\text{(RHIC)}}$

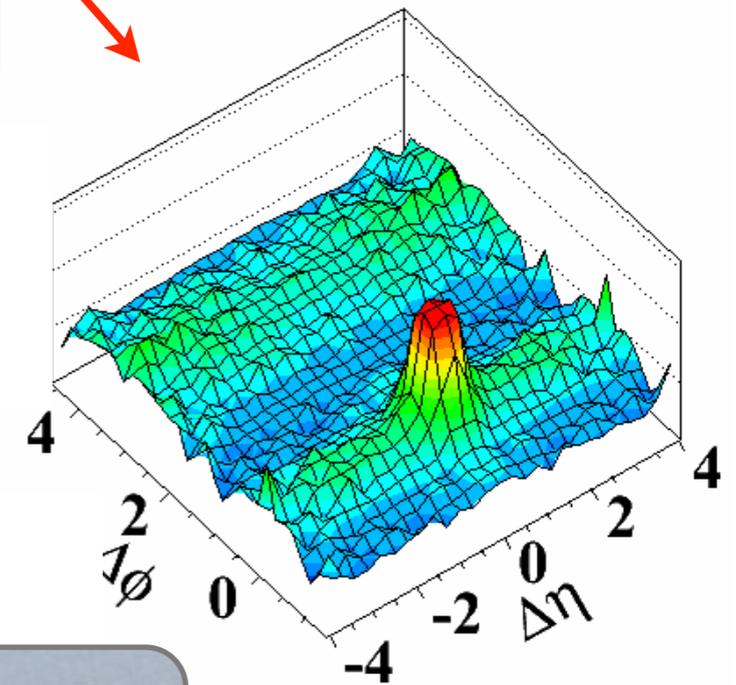


# A brief history of correlations in HIC





Two-hadron  
Correlation Function

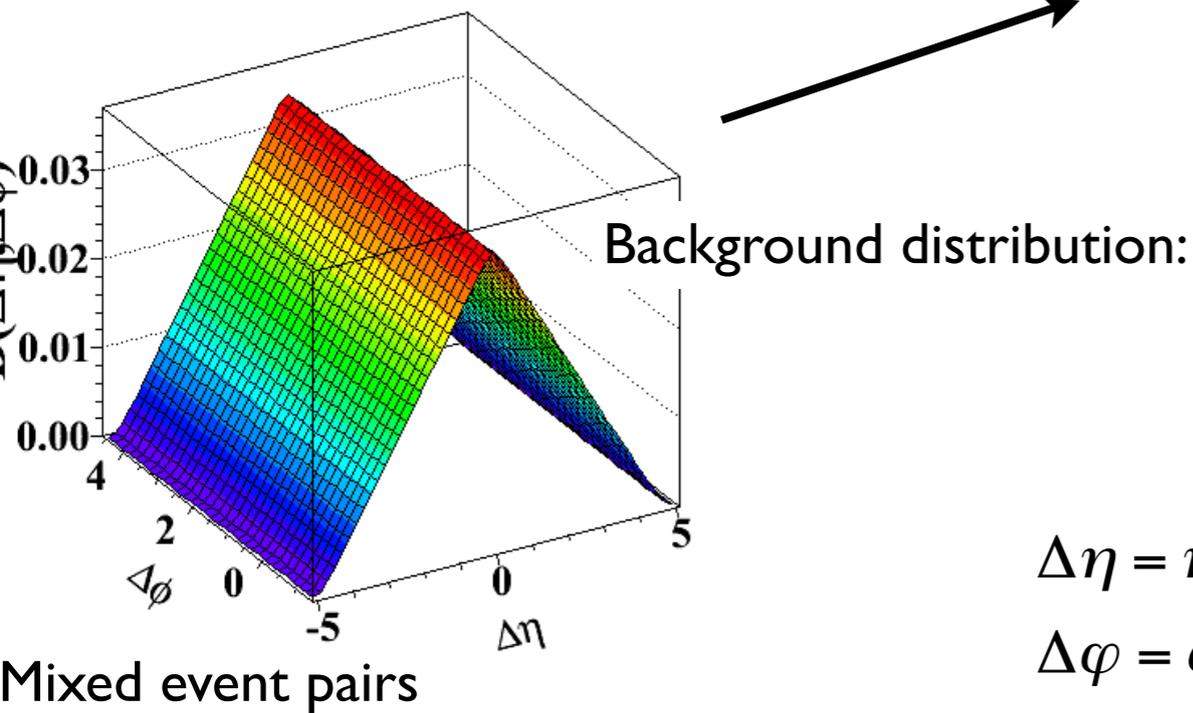
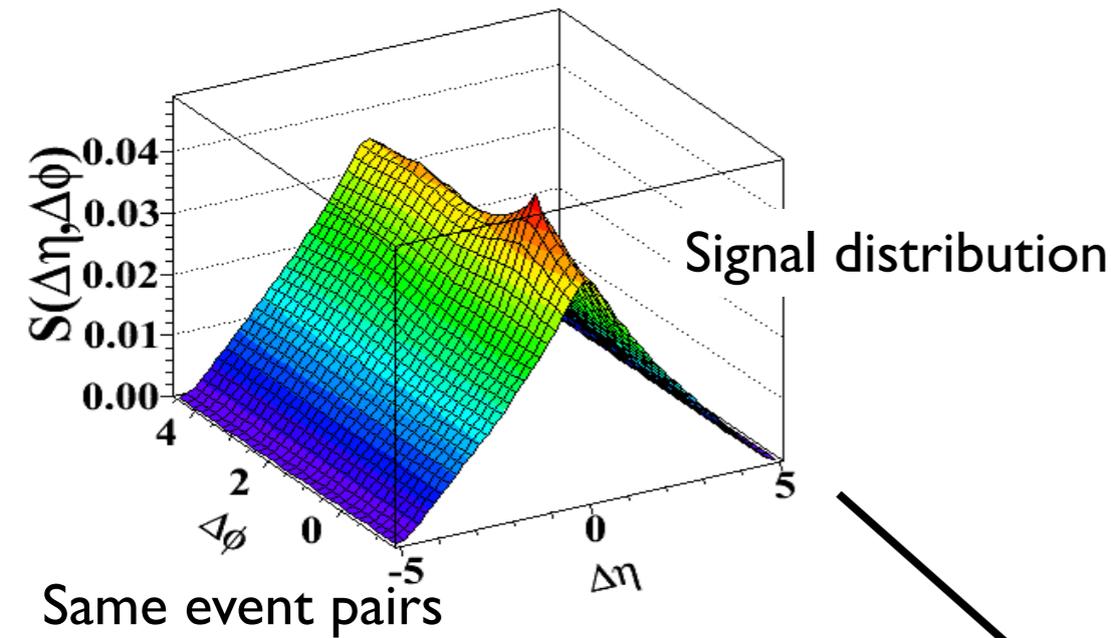


Initial State Correlations  $\otimes$  Hydrodynamic Flow =

Jets  
Glasma Flux Tubes  
Hot Spots  
Glauber Fluctuations

## Key Physics Topics:

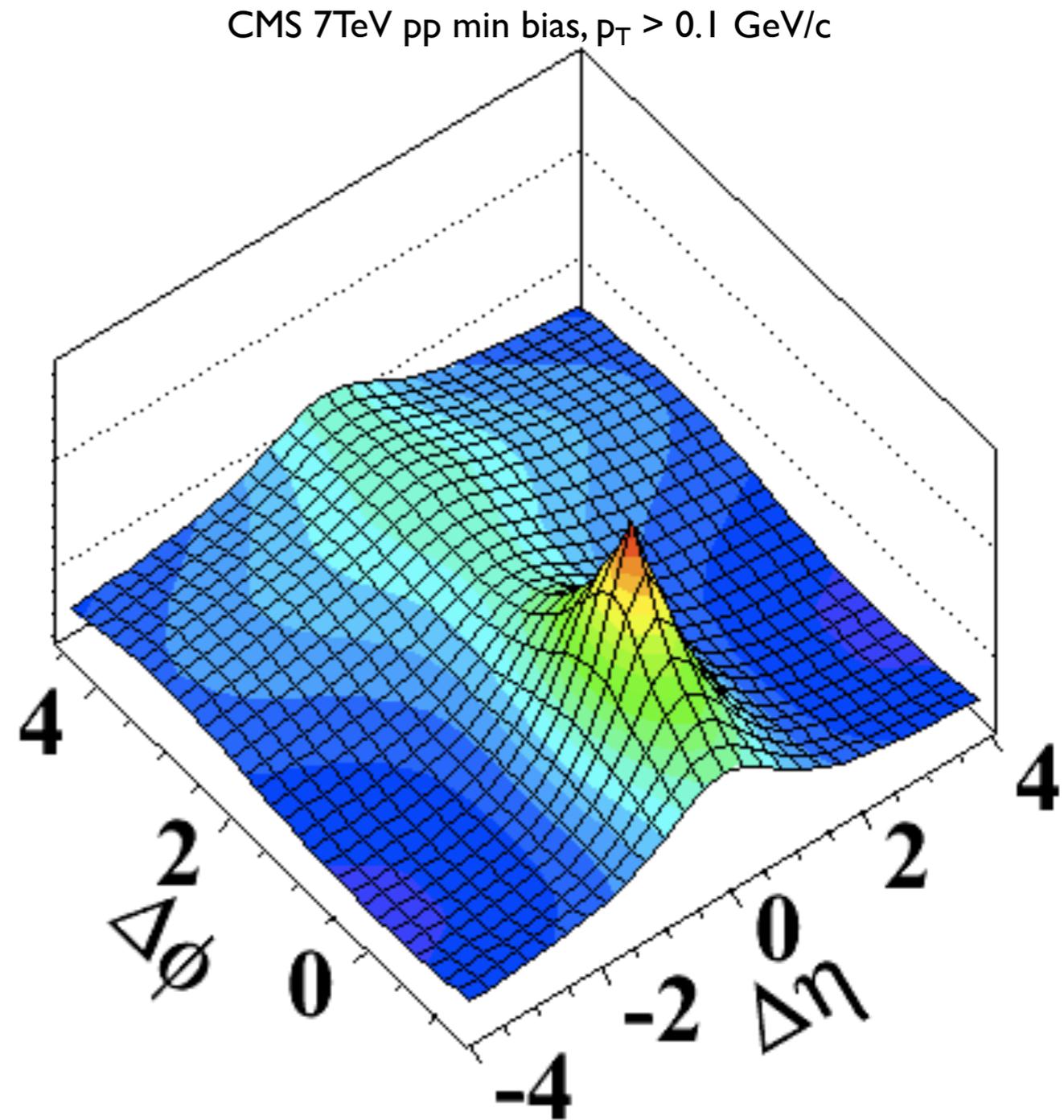
- Jet energy loss
- Jet-medium interactions
- Transport properties of the medium



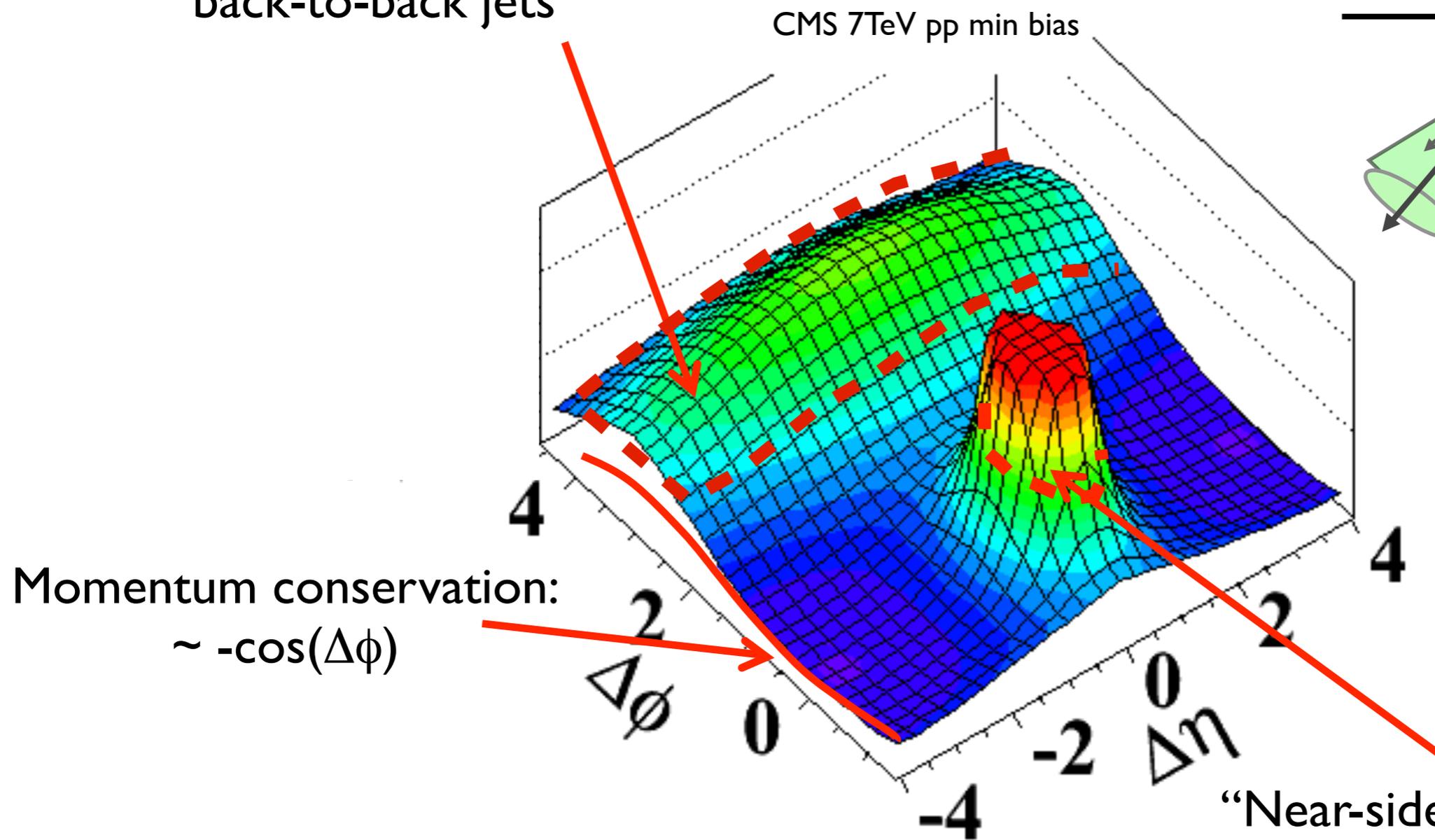
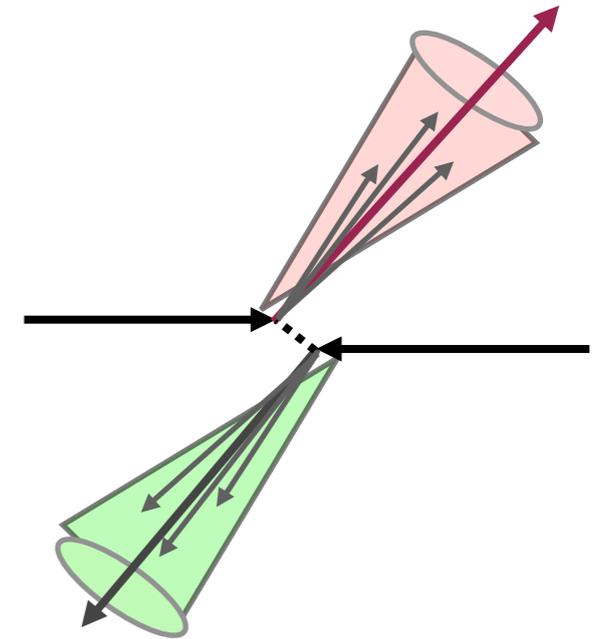
Ratio  
Signal/Background

$$\Delta\eta = \eta_1 - \eta_2$$

$$\Delta\varphi = \varphi_1 - \varphi_2$$



“Away-side” ( $\Delta\phi \sim \pi$ ) jet correlations:  
 Correlation of particles between  
 back-to-back jets

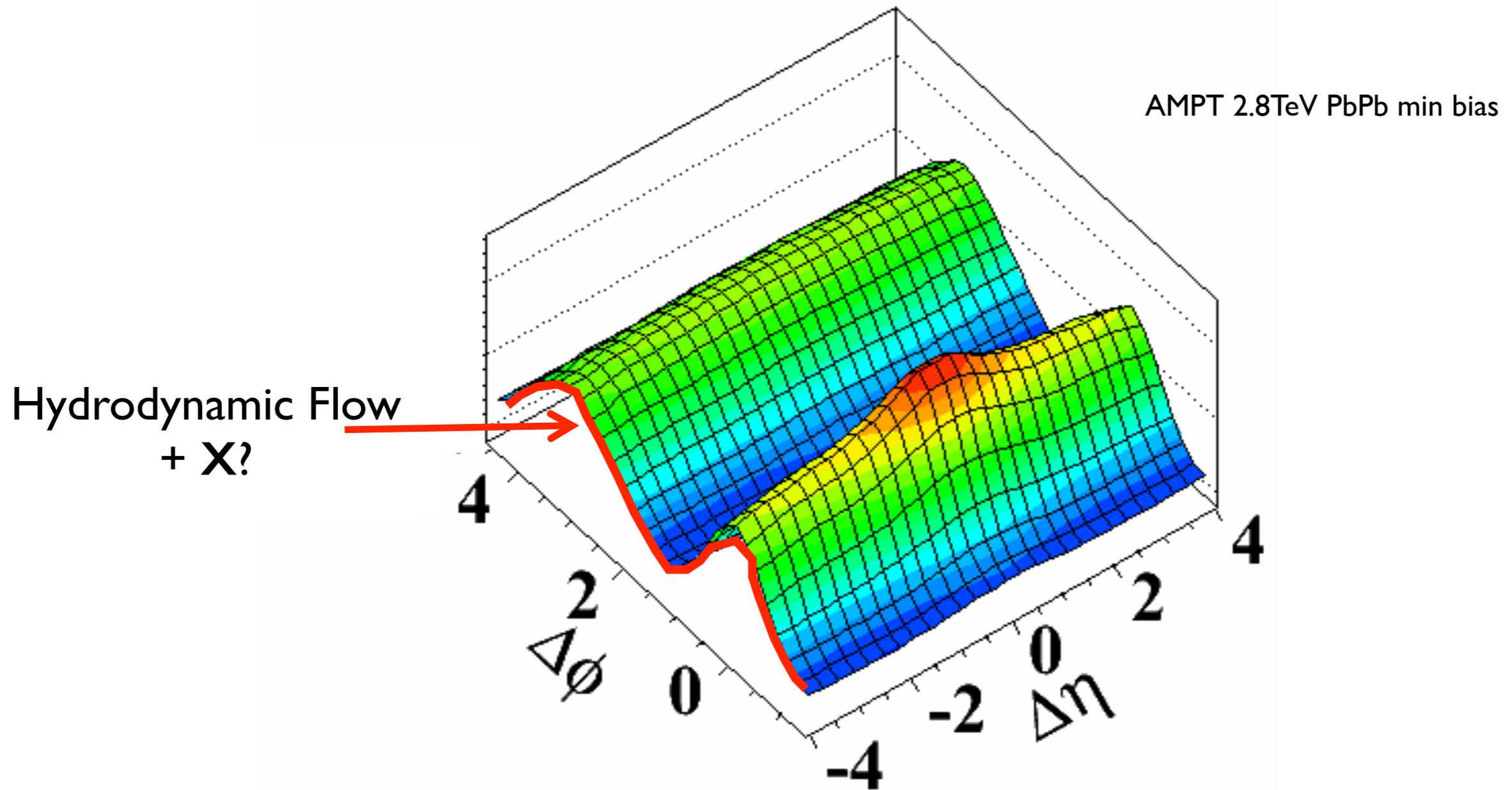


Momentum conservation:  
 $\sim -\cos(\Delta\phi)$

$1 < p_T < 3 \text{ GeV}/c$

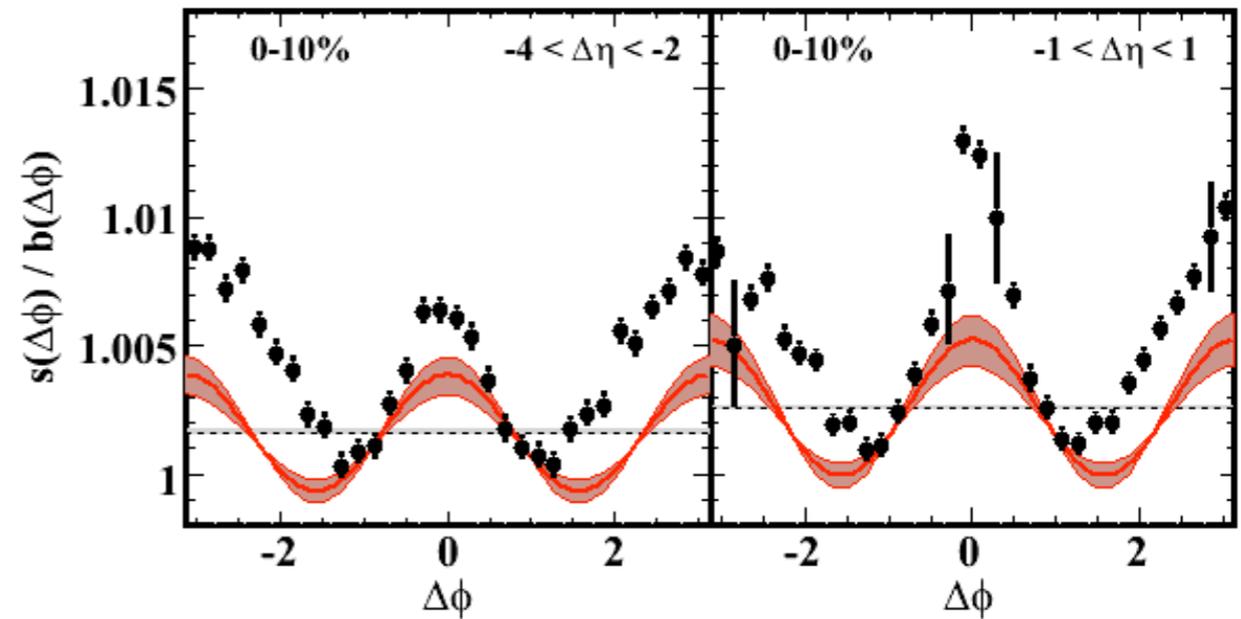
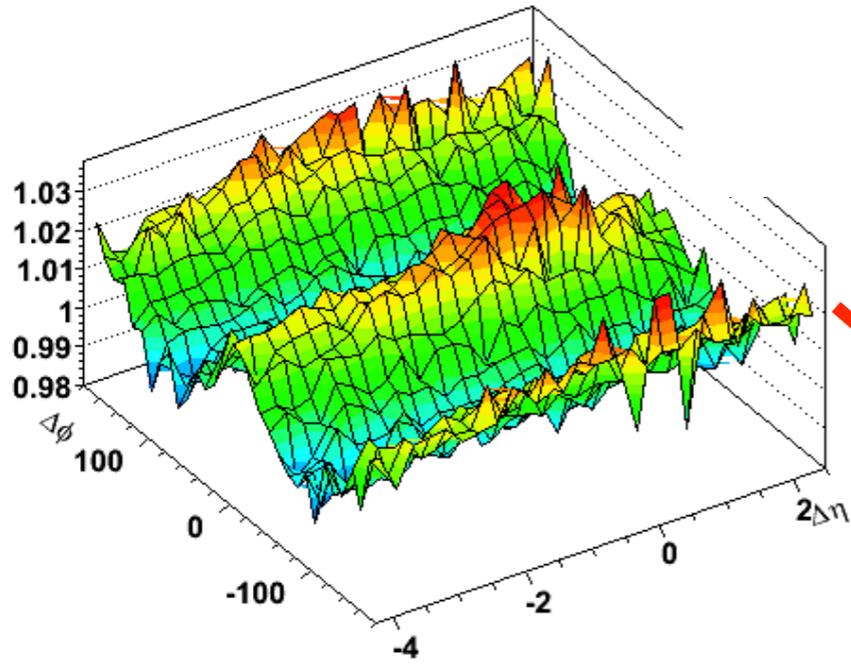
Dominated by jet-like correlations

“Near-side” ( $\Delta\phi \sim 0$ ) jet peak:  
 Correlation of particles  
 within a single jet



PbPb correlations,  $1 < p_T < 3$  GeV/c  
 Large  $v_2$  term from elliptic flow dominates correlations

“Raw” correlation function

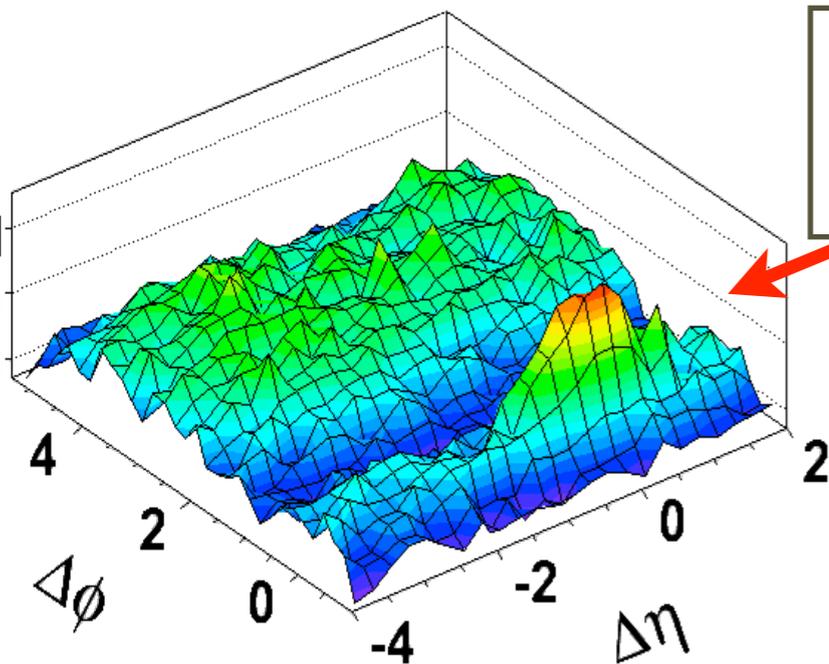


“ZYAM” (zero yield at minimum): assume that one component of the correlations (jets) gives zero contribution at some  $\Delta\phi$ ; match  $v_2$  flow at that  $\Delta\phi$  and subtract

$$\frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{ch}}}{d\Delta\phi d\Delta\eta} = B(\Delta\eta) \left\{ \frac{s(\Delta\phi, \Delta\eta)}{b(\Delta\phi, \Delta\eta)} - a(\Delta\eta) [1 + 2V(\Delta\eta) \cos(2\Delta\phi)] \right\}$$

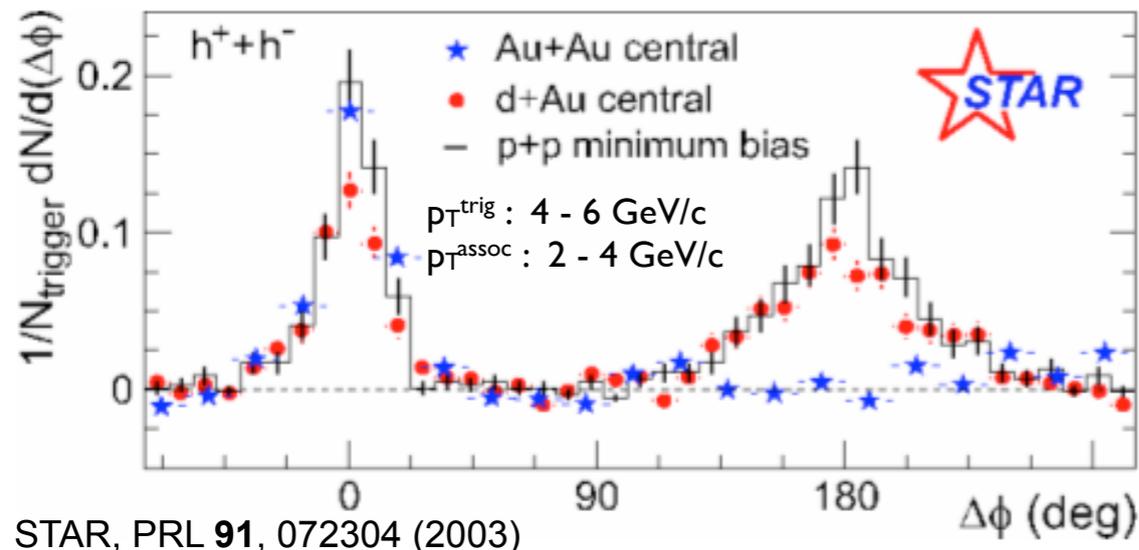
for some measurements: also  $v_4$

Normalization term to go from correlation amplitude to yield per trigger particle



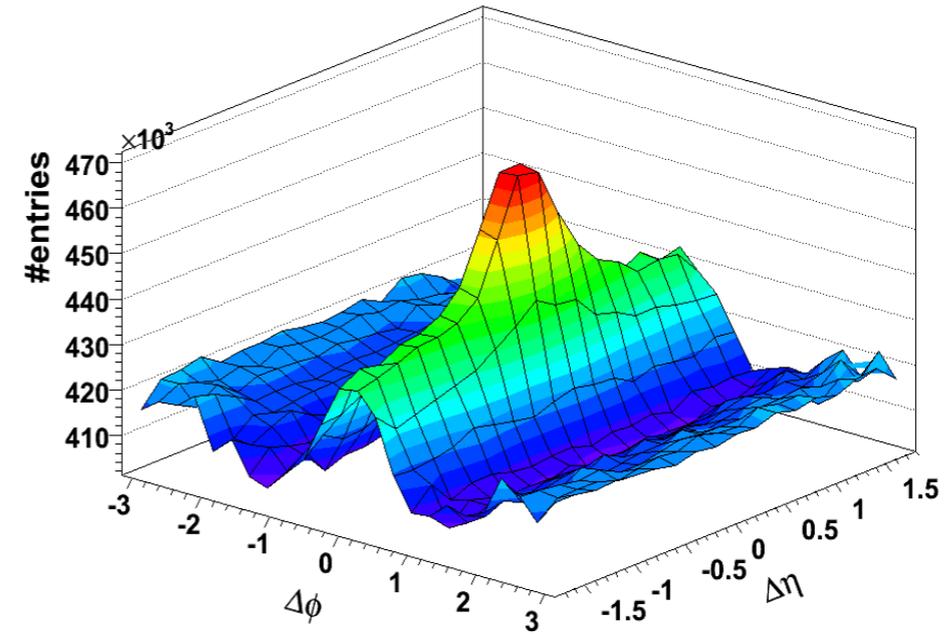
$v_2$  subtracted “associated yield”

## “Disappearance” of away-side jet correlations in central AuAu



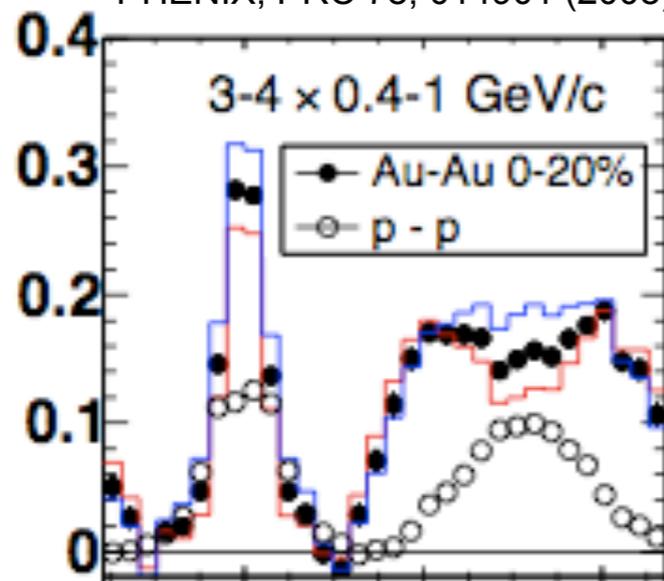
$p_T^{\text{trig}} : 3-4, p_T^{\text{assoc}} : >2 \text{ GeV/c}$

STAR, J.Phys.G34:S679-684,2007



$p_T^{\text{trig}} : 3-4, p_T^{\text{assoc}} : 0.4-1 \text{ GeV/c}$

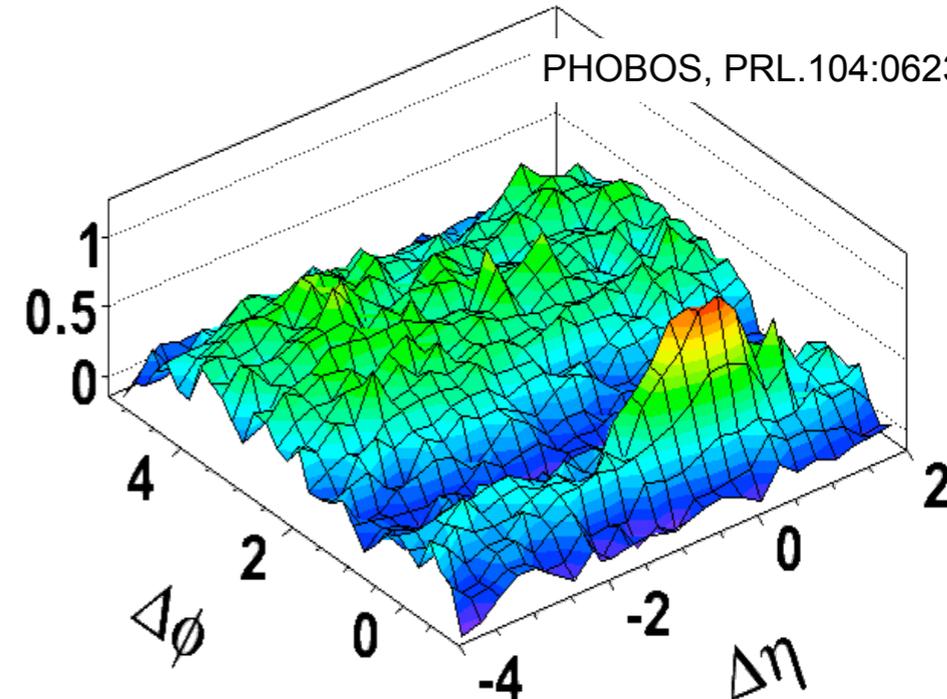
PHENIX, PRC **78**, 014901 (2008)



## Emergence of long-range near-side correlations in AuAu: Ridge

$p_T^{\text{trig}} : >2.5, p_T^{\text{assoc}} : >0 \text{ GeV/c}$

PHOBOS, PRL.104:062301 (2010)

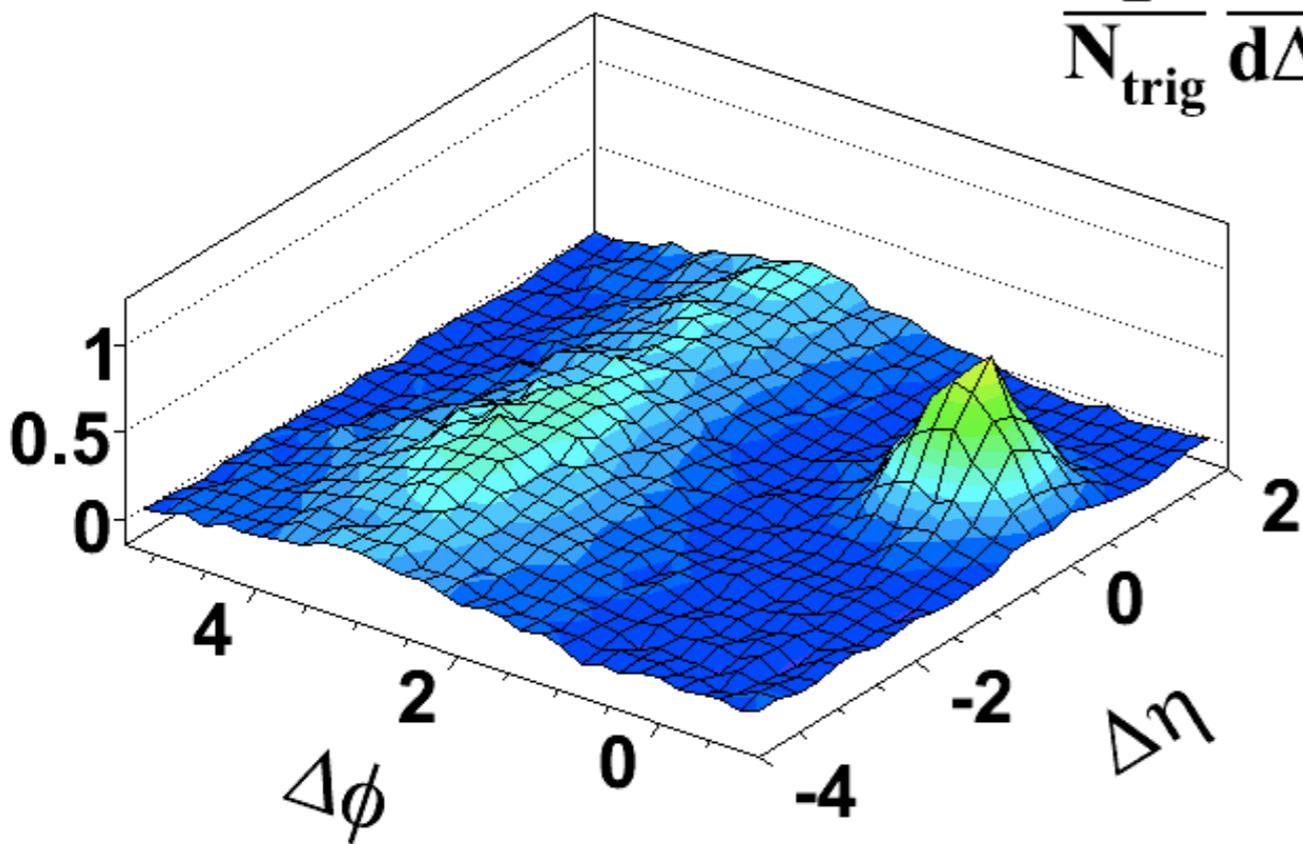


Broadening of away-side jet correlations in central AuAu:  
Conical emission?

# Results

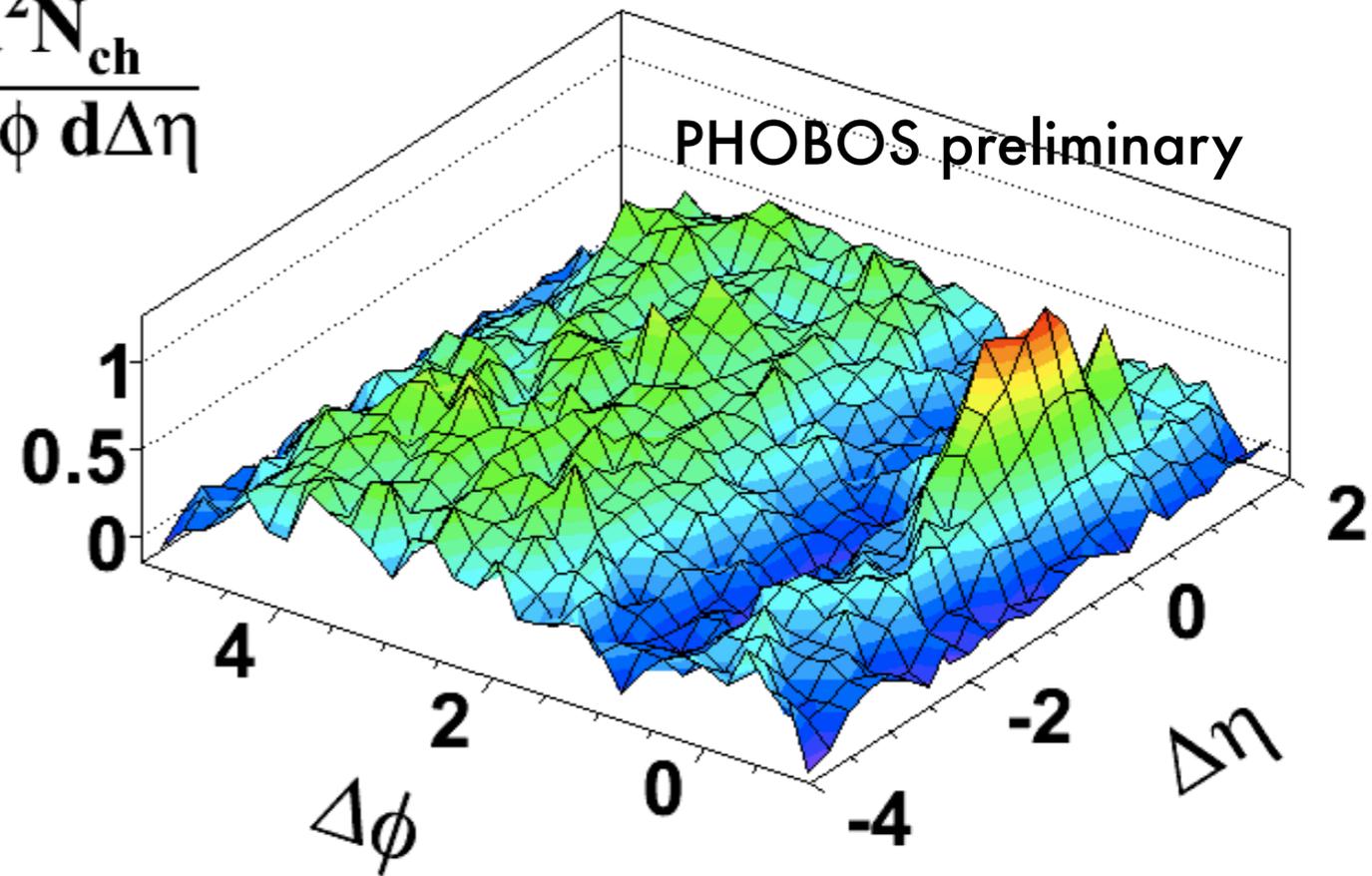
$p_T^{\text{trig}} > 2.5 \text{ GeV}/c$   
 $p_T^{\text{assoc}} \geq 20 \text{ MeV}/c$

p+p PYTHIA v6.325



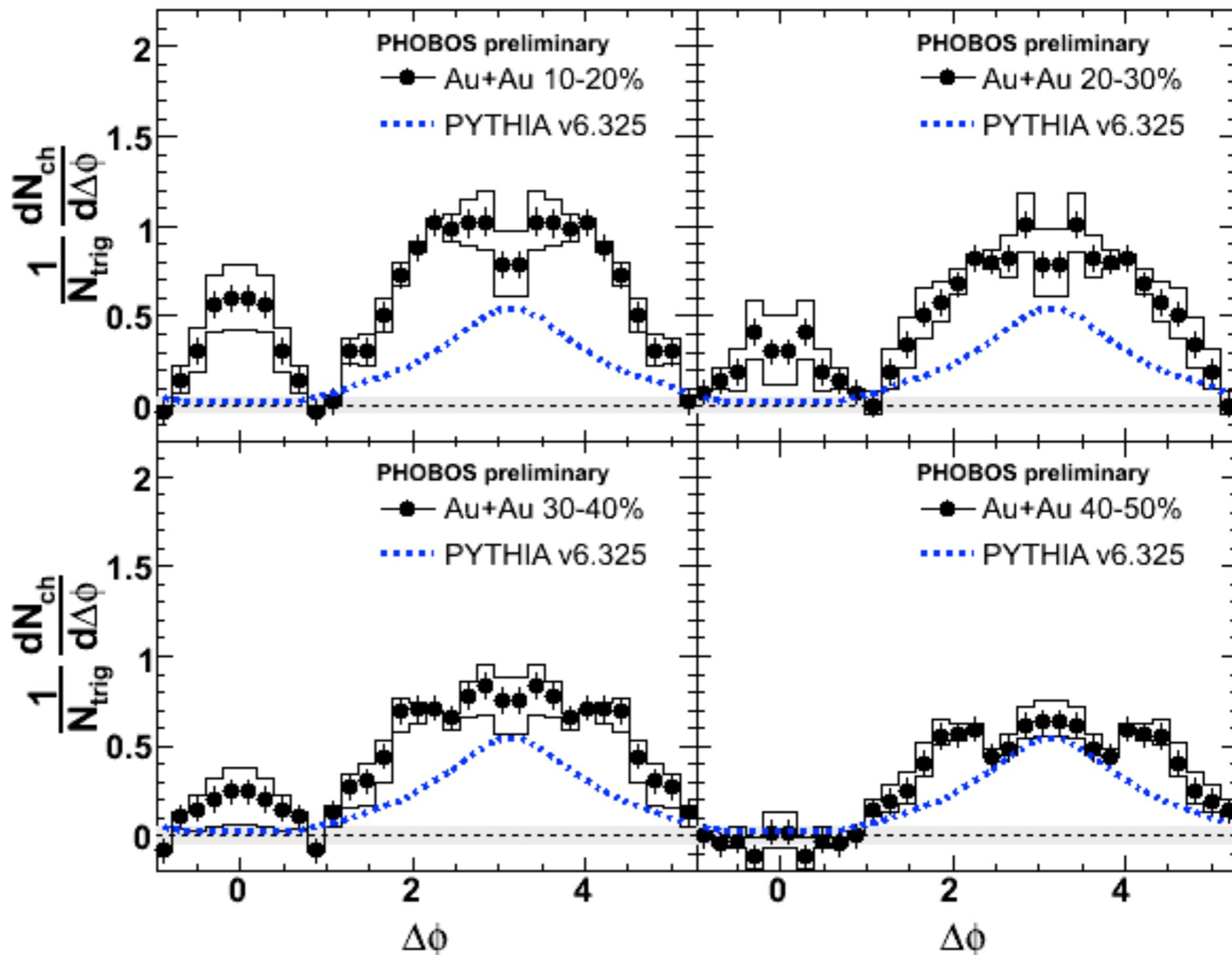
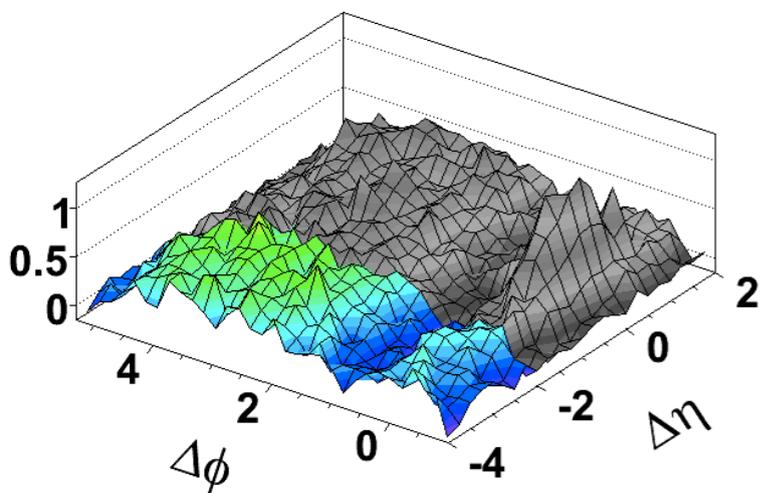
$$\frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{ch}}}{d\Delta\phi d\Delta\eta}$$

Au+Au 0-30% central



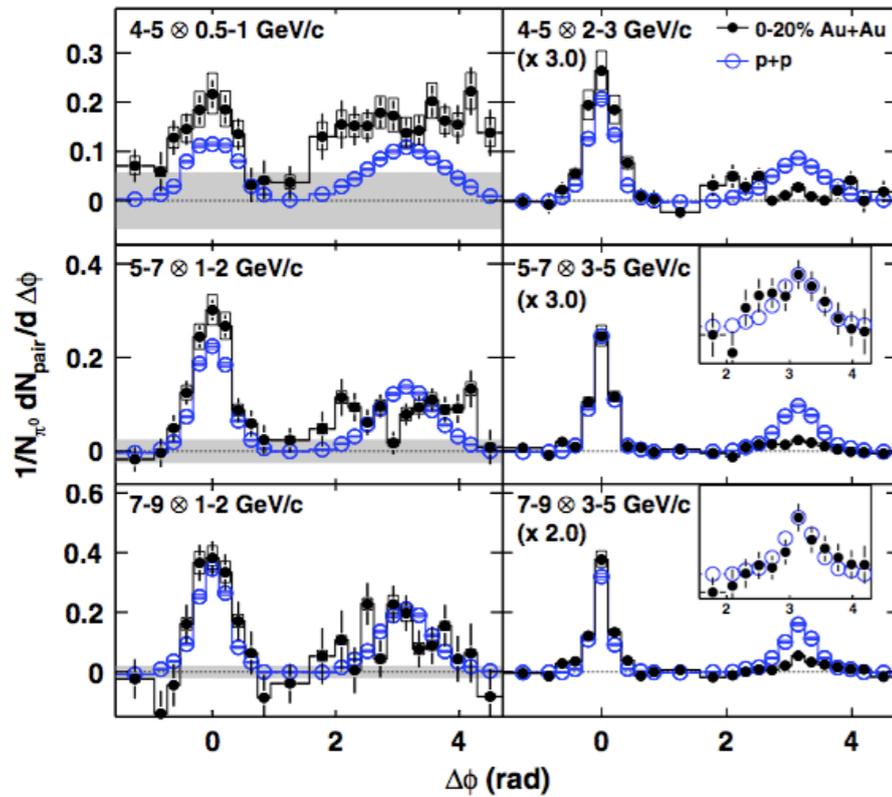
# Long-Range Centrality Dependence

$$-4 < \Delta\eta < -2$$

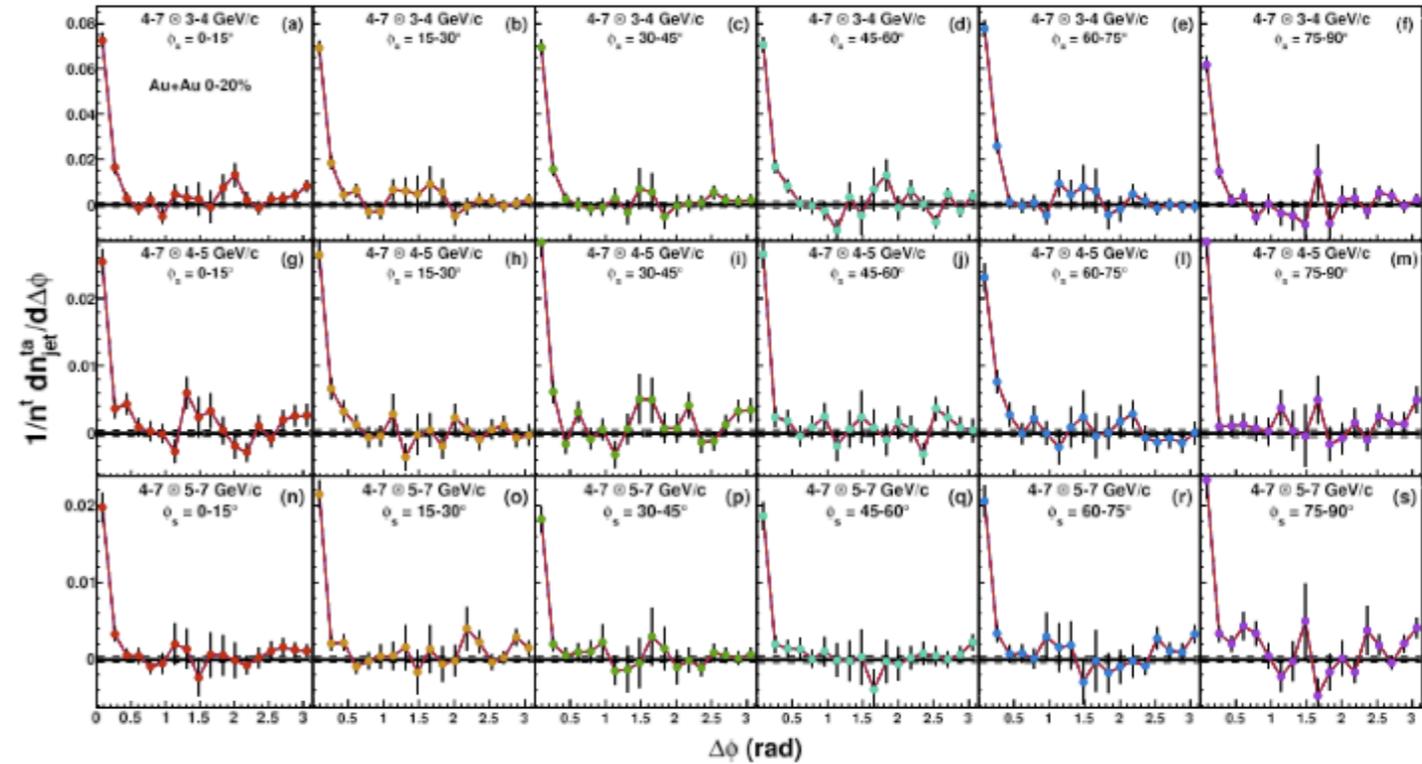




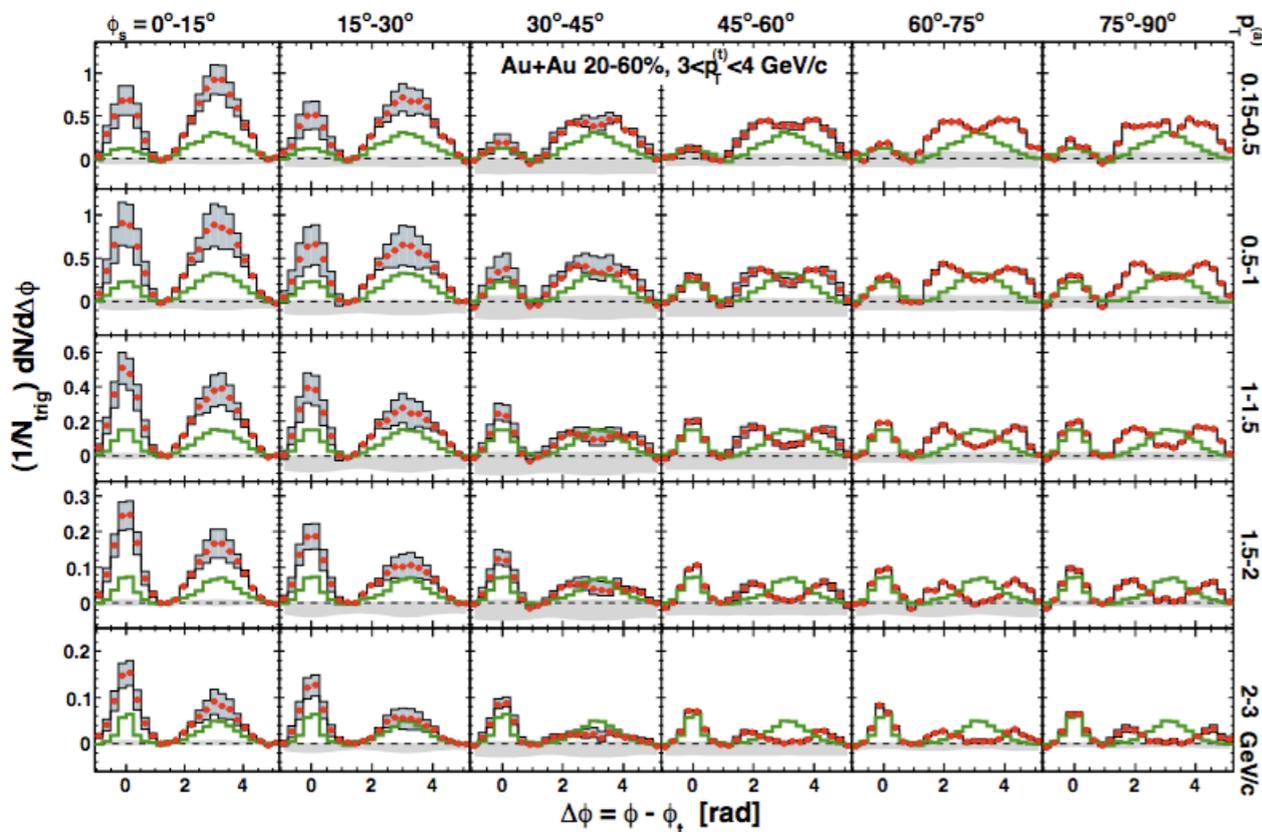
# This Talk: Selected Recent Results



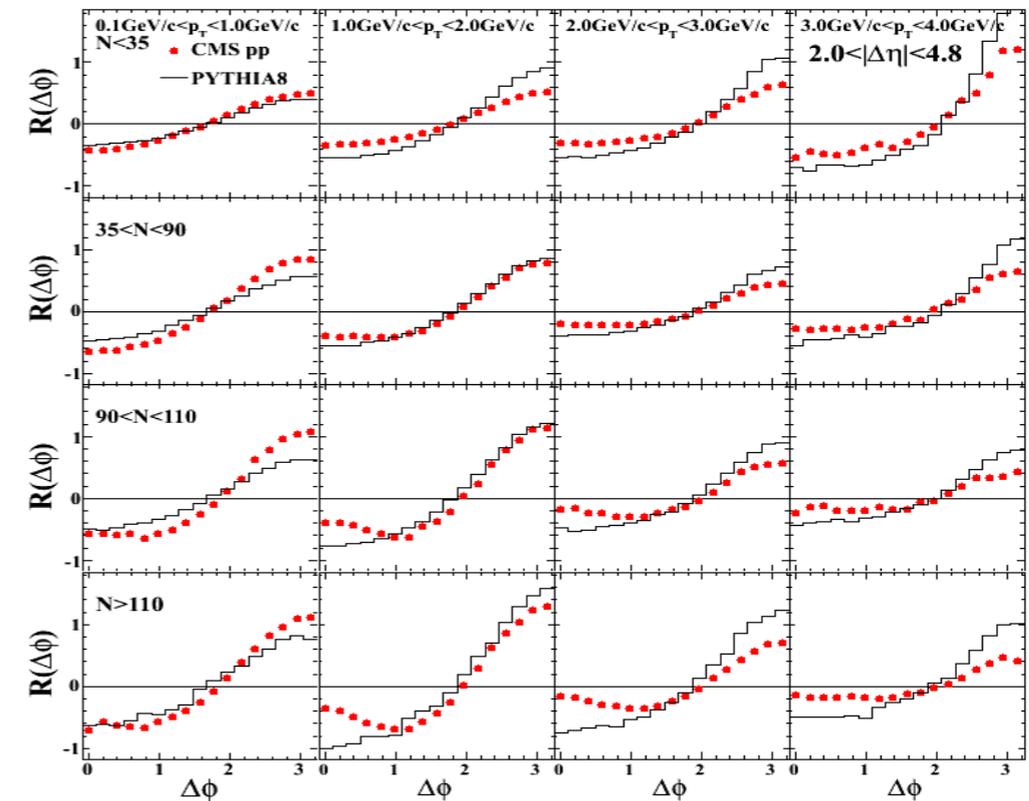
PHENIX arXiv:1002.1077



PHENIX arXiv:1010.1521



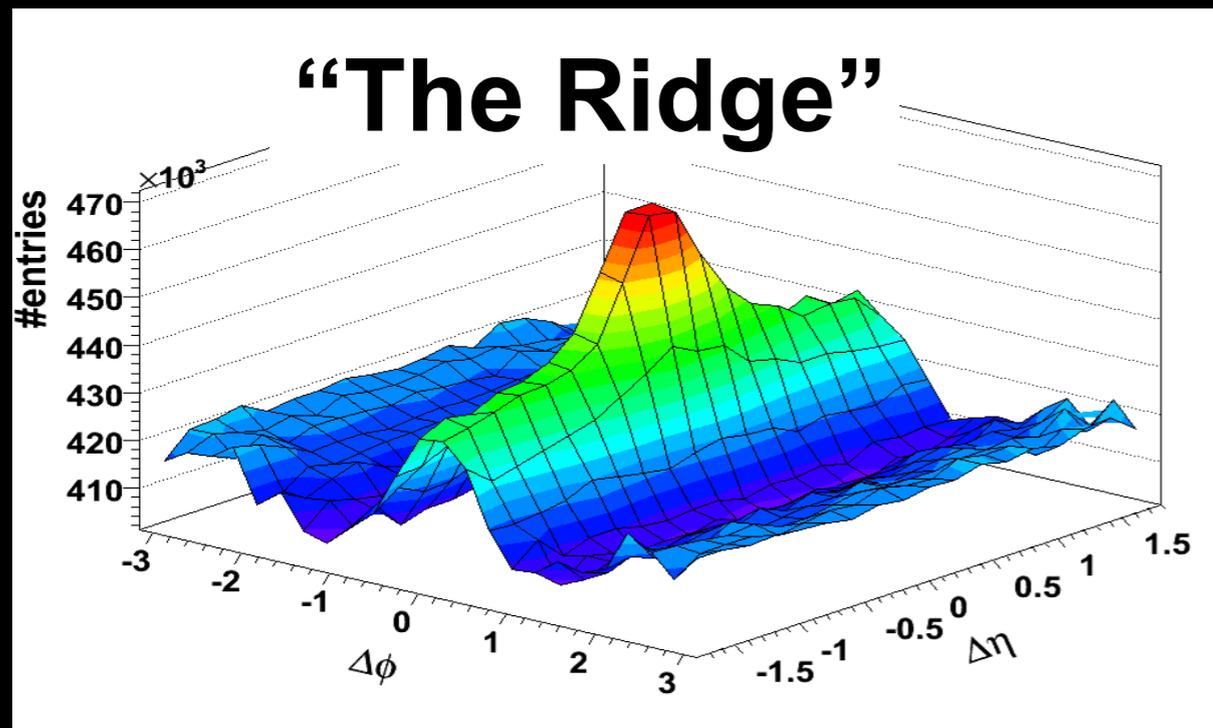
STAR arXiv:1010.0690



CMS arXiv:1009.4122

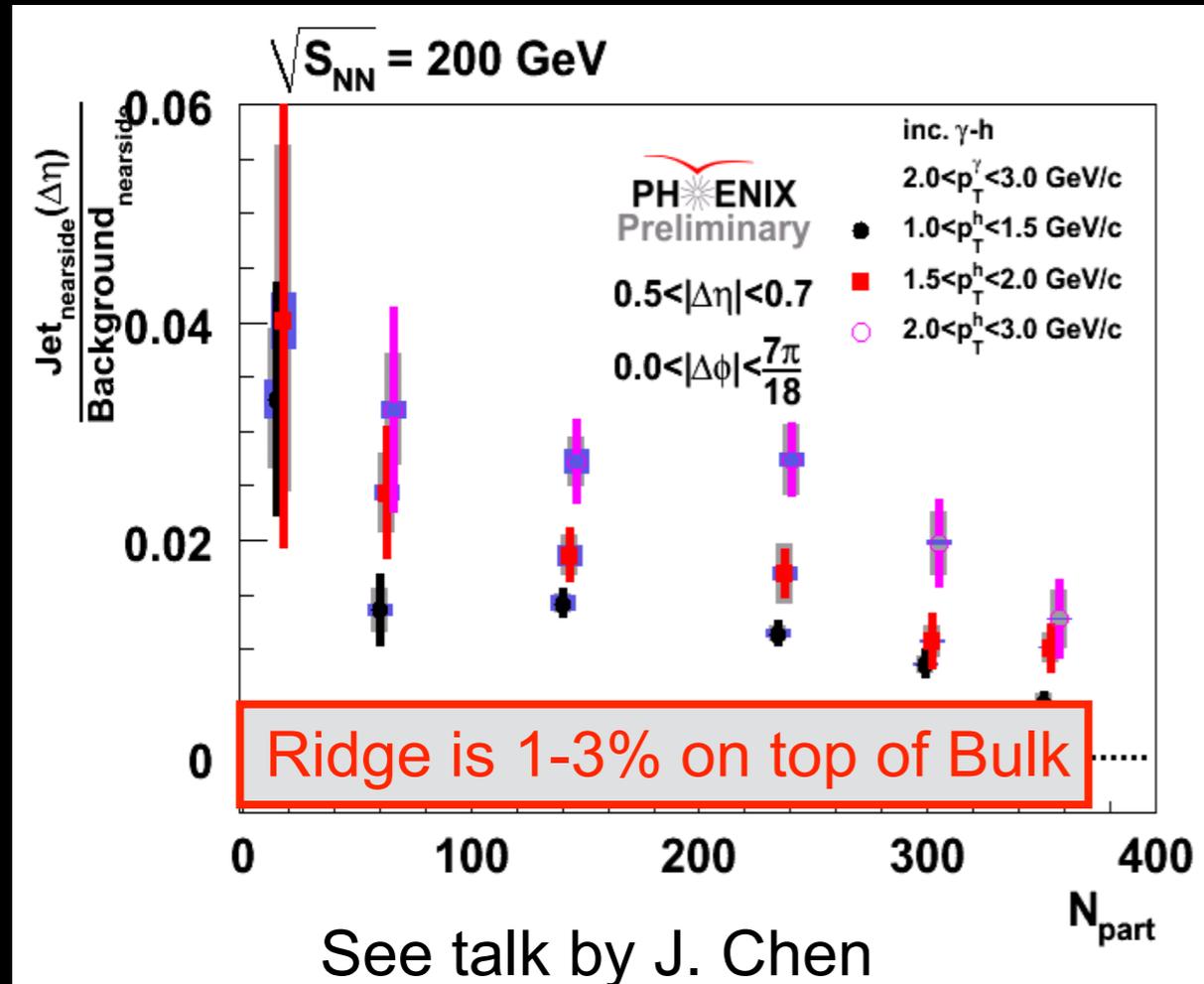


# Medium Response? QM 2009



Basic Properties:

1.  $p_T$  spectra similar to bulk (or slightly harder)
2. baryon/meson enhancement similar to bulk
3. Scales per trigger like  $N_{part}$  similar to bulk

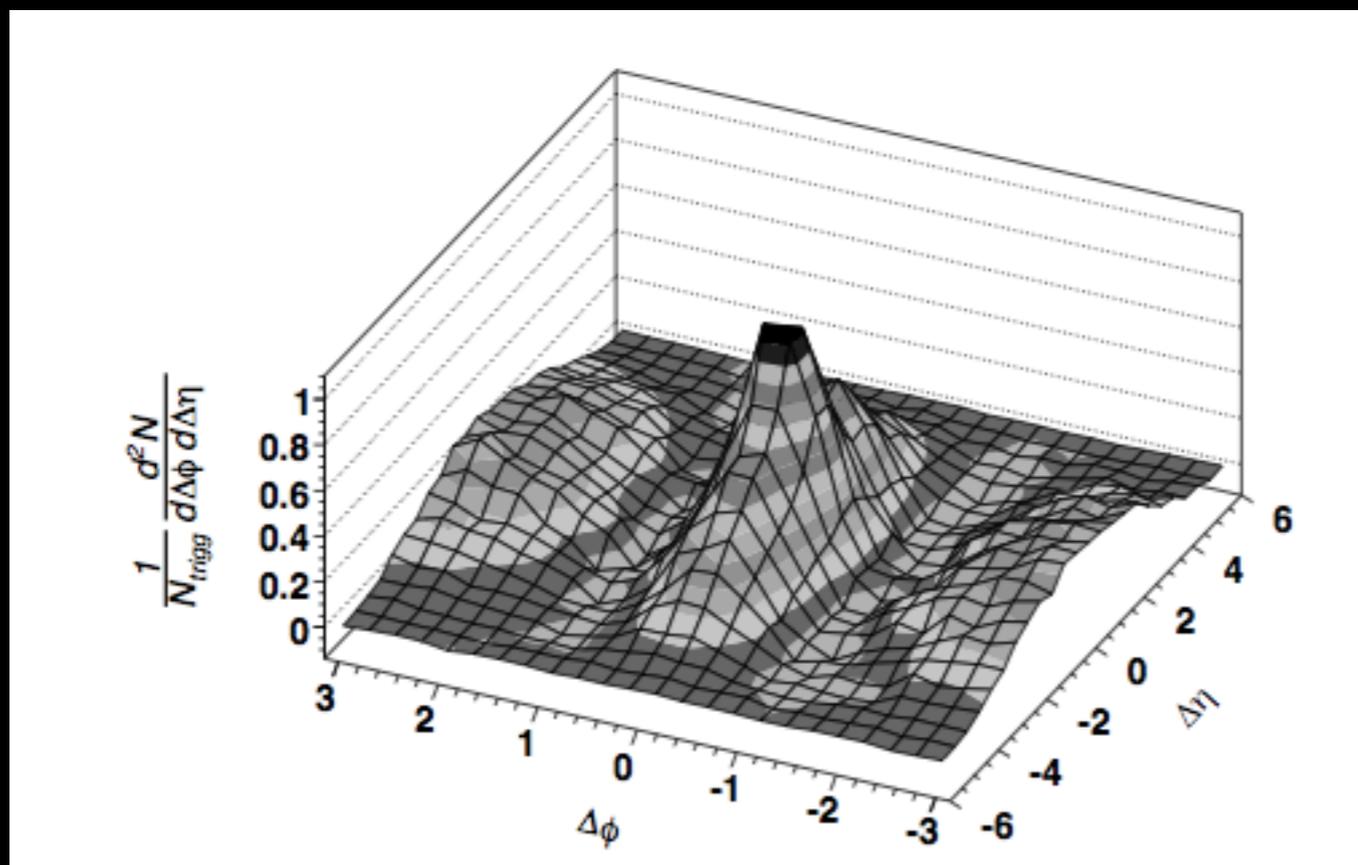


"Theoretical Free-For-All"

Paul Stankus

"Theorists, help us Kill your model."

Brian Cole QM08



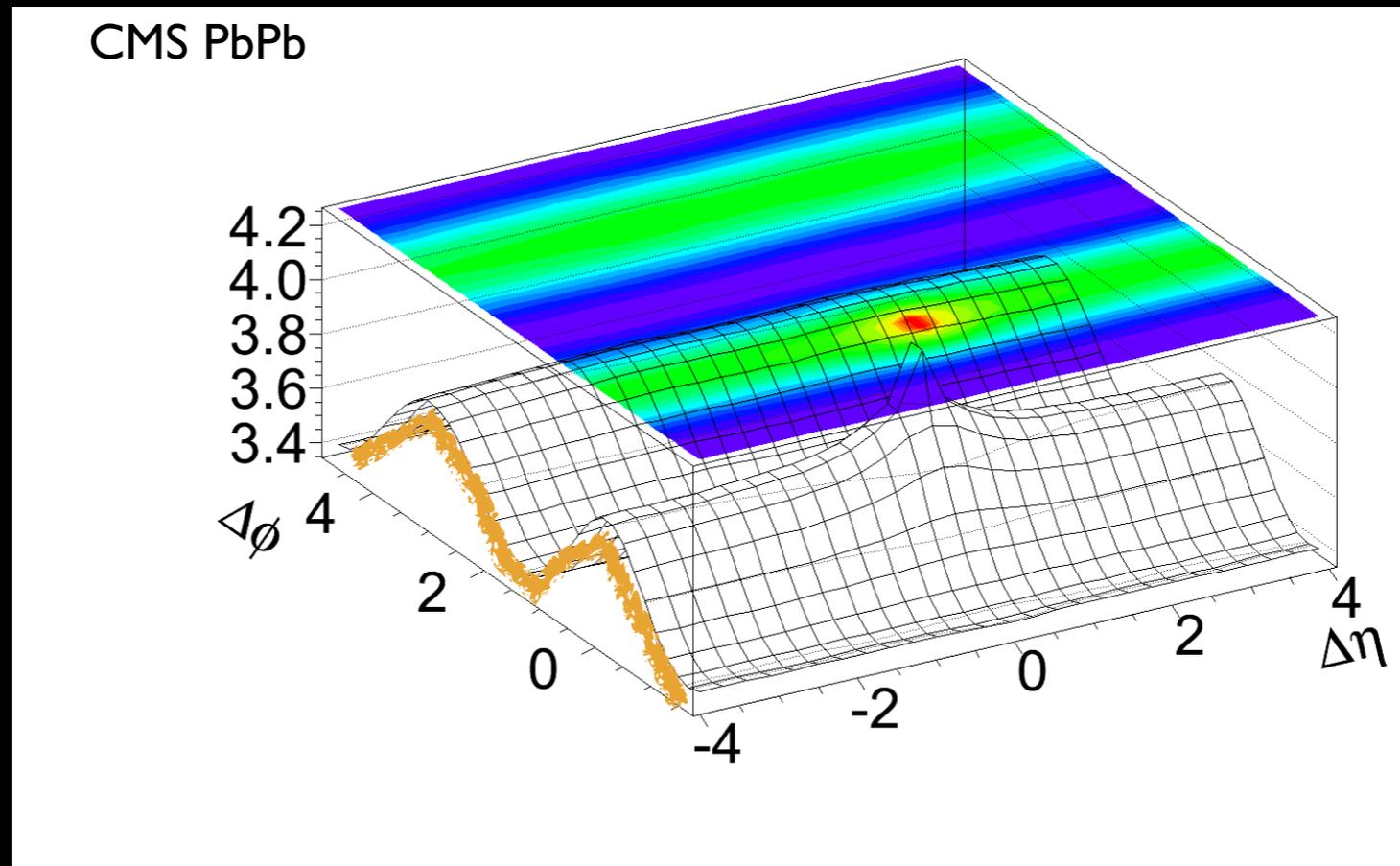
Takahashi et al, arXiv:0902.4870  
 Phys.Rev.Lett. 103 (2009) 242301

Hydro with MC Glauber initial conditions

### 3.5. Ridge effect

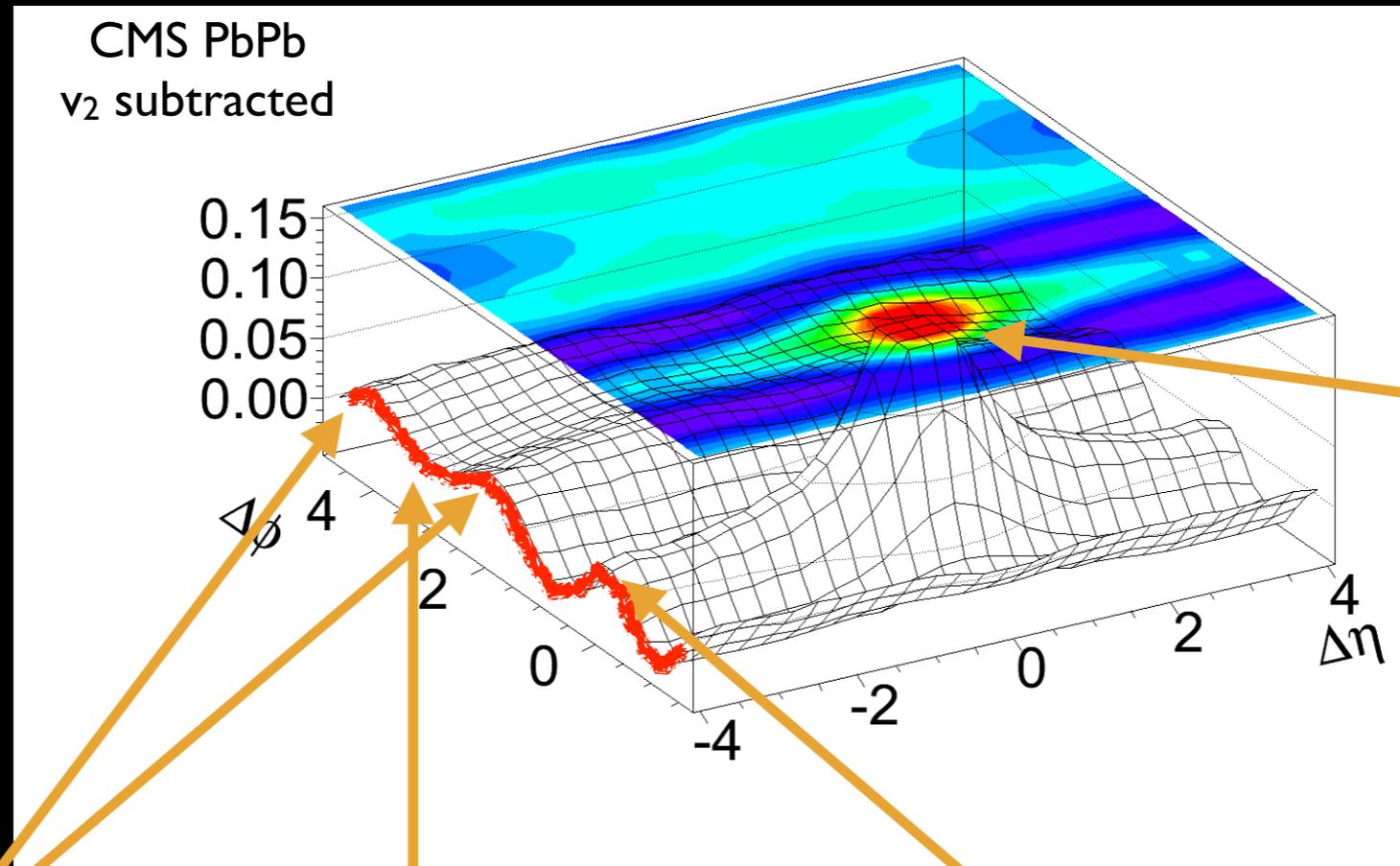
Another effect, which is produced naturally by the longitudinal baton structure of IC, as shown in Fig.1, is the so-called *ridge phenomenon* which has been experimentally seen in high- $p_T$  nearside correlations [12]. Since

Subtract or suppress elliptic flow:  
Complex remnant correlation structure



# Beyond elliptic flow

Subtract or suppress elliptic flow:  
Complex remnant correlation structure



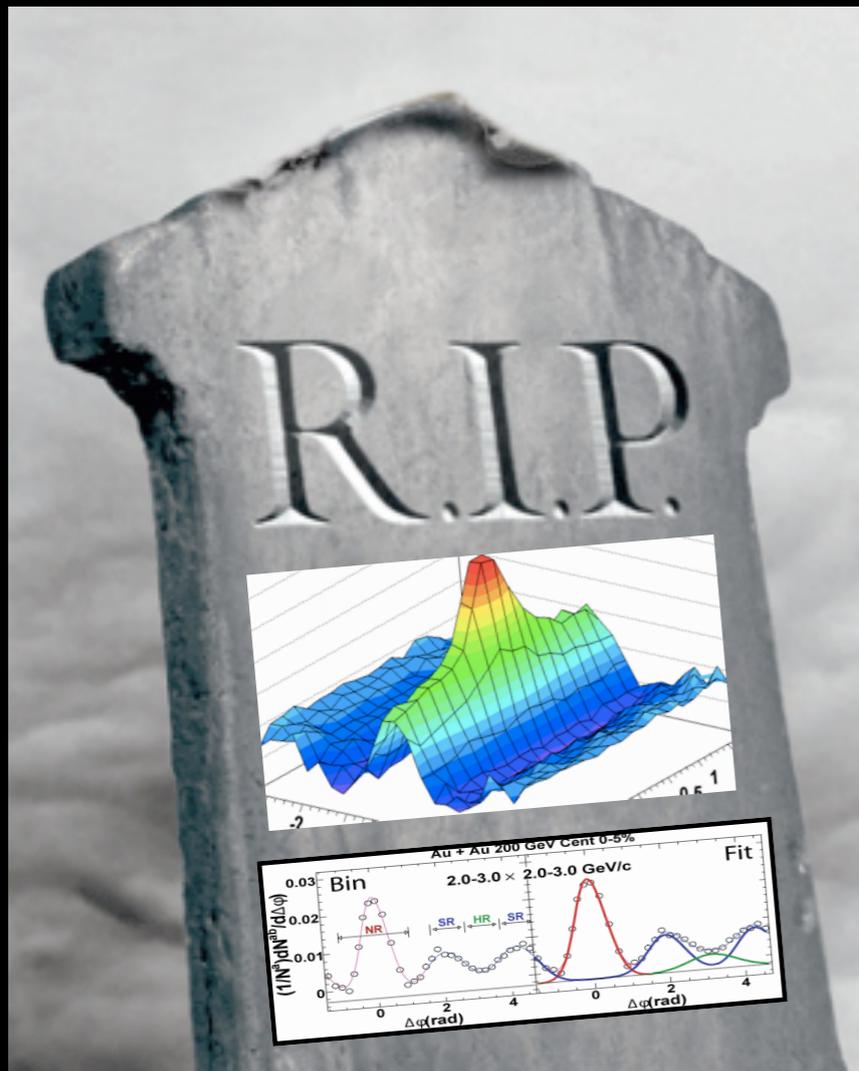
“Shoulder”  
“Mach cones”  
PHENIX (2005)  
STAR (2010)

“Head”  
“ $\pi$  region”  
PHENIX (2005)  
STAR (2010)

“Ridge”  
STAR (2006)  
PHOBOS (2008)

Jet correlations

2005-2010:  
Extensive literature (100's of papers)  
on each of these features



“Triangular flow”  
and  
“Participant triangularity”  
No ridge, no mach cone, just flow

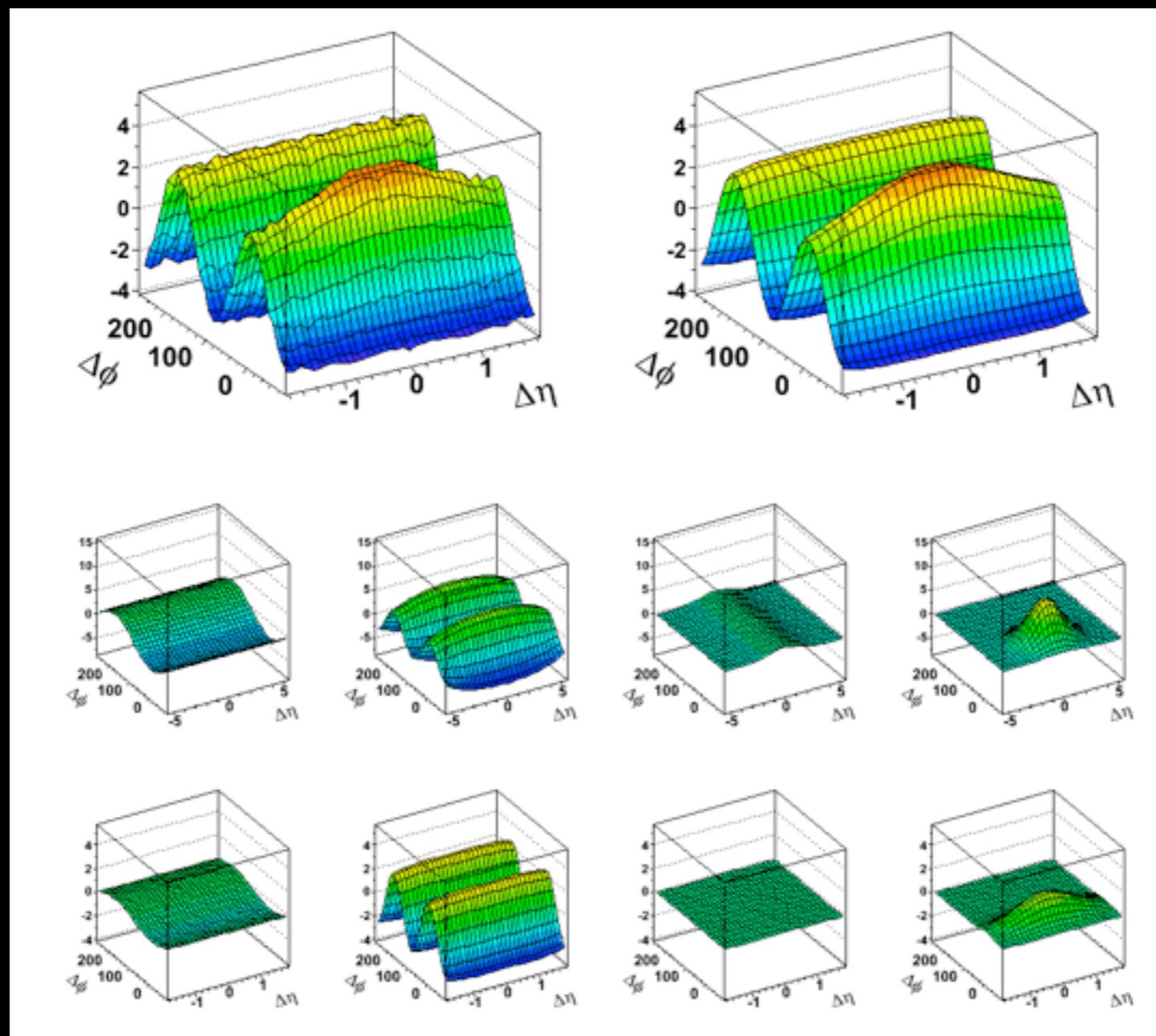
Burak Alver  
Mithig meeting  
10/13/2009

Geometry fluctuations and  
final state correlations:

Mishra et al arXiv:0711.1323

Takahashi et al, arXiv:0902.4870

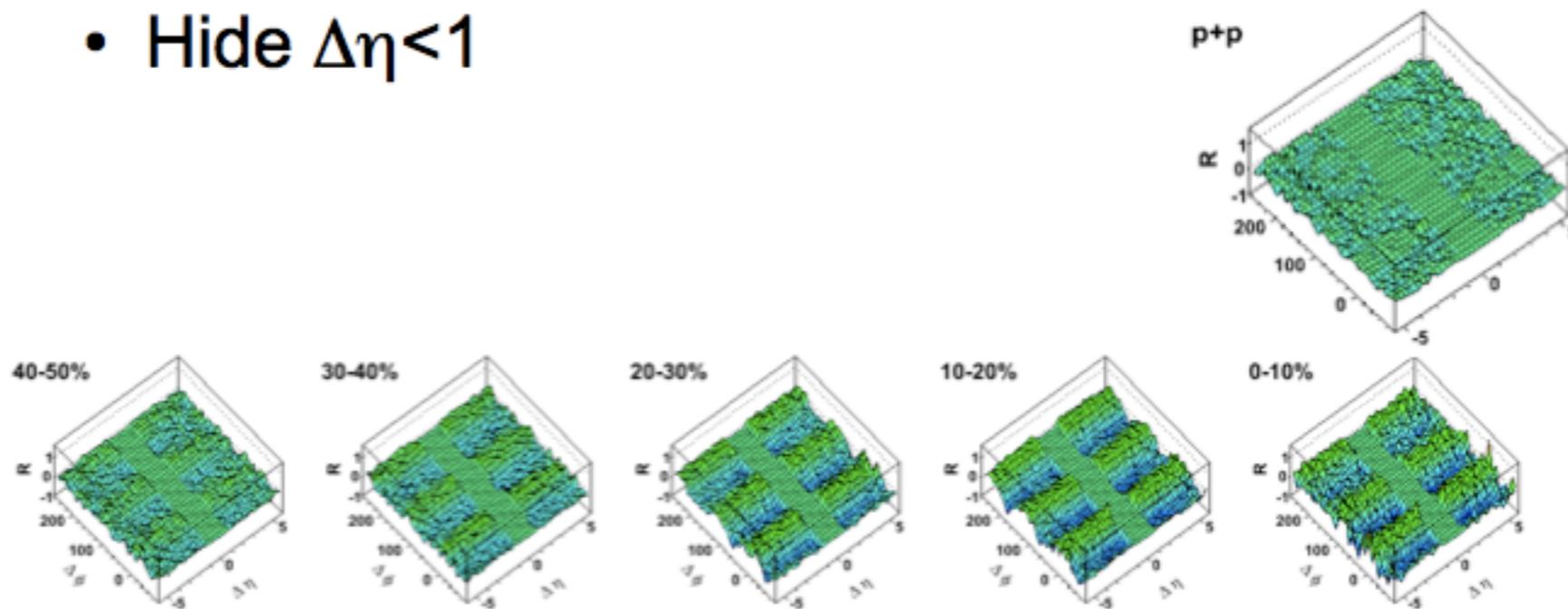
Sorensen, arXiv:1002.4878



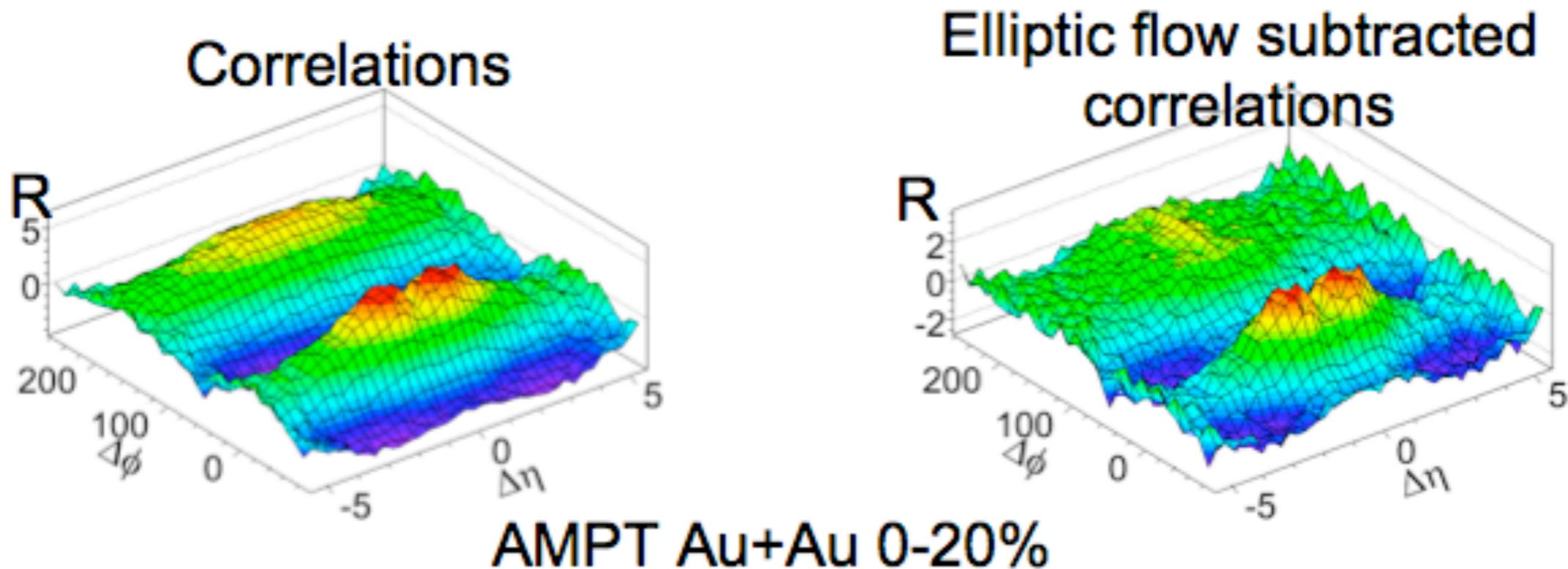
## Decomposing CFs à la STAR

## Let's look at "non-flow"

- Flow subtraction is tricky
  - Let's just take out all First and Second FC
- Normalize in bins of  $\Delta\eta$  (a la  $\Delta\eta$  of ZYAM)
- Hide  $\Delta\eta < 1$



AMPT model: Glauber initial conditions, collective flow



AMPT model also produces similar correlation structures that extend out to long range in  $\Delta\eta$ .

Lin et. al. nucl-th/0411110

“Triangular flow”

and

“Participant triangularity”

No ridge, no mach cone, just flow

Burak Alver

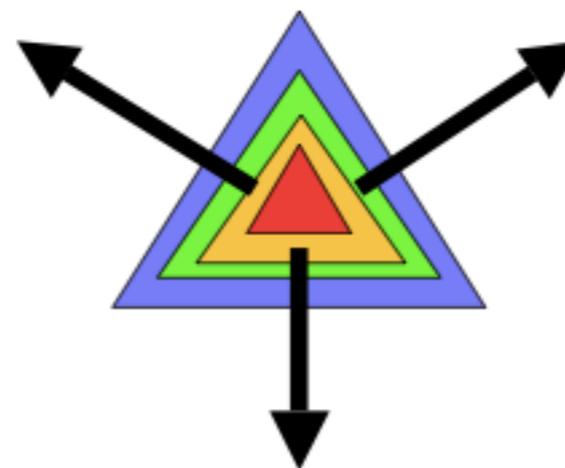
Mithig meeting

10/13/2009

## “Participant Triangularity” $\tau$

If the initial geometry were a triangle,  
then we would expect global  $v_3$

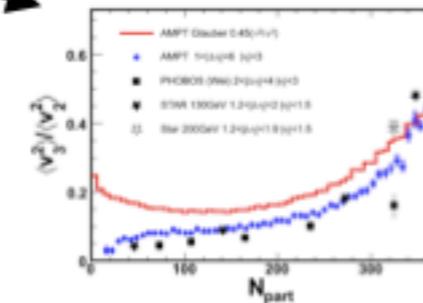
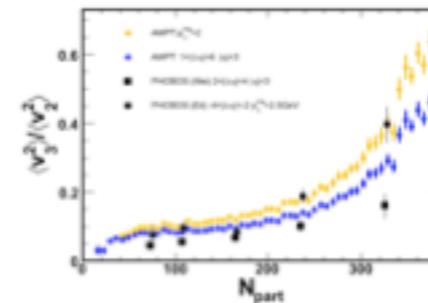
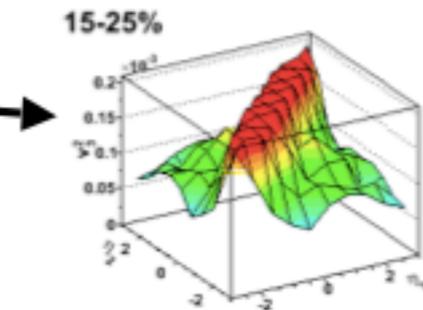
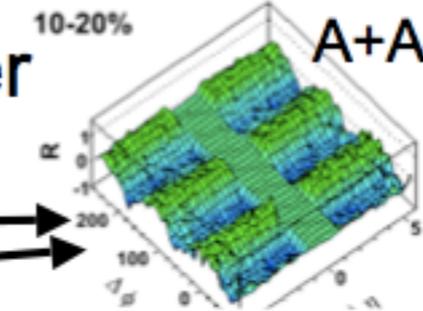
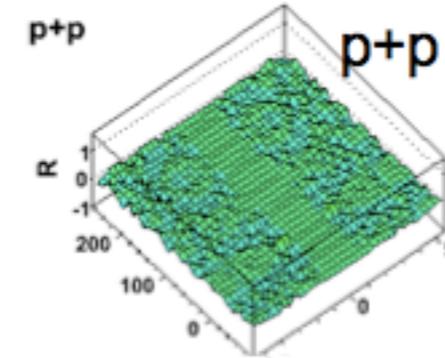
Define  
participant triangularity  $\tau$   
analogous to  
participant eccentricity  $\varepsilon$



## “Triangular Flow”

Why should we think the third Fourier coefficient is due to a global effect?

- It is large
- It is there at large  $\Delta\eta$
- It is a function of  $\eta$  rather than  $\Delta\eta$
- $v_3 / \tau$  : makes sense.
- $v_3(p_T)$  makes sense.



(This is an outline slide)

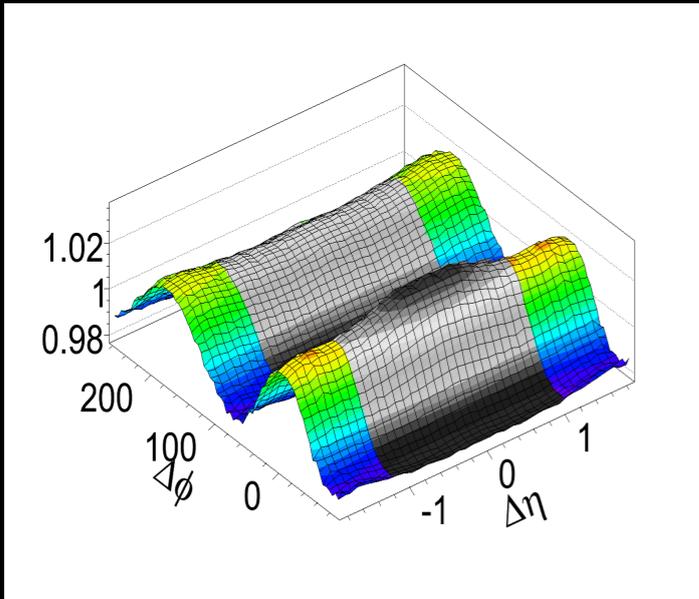


# Comparison to published data

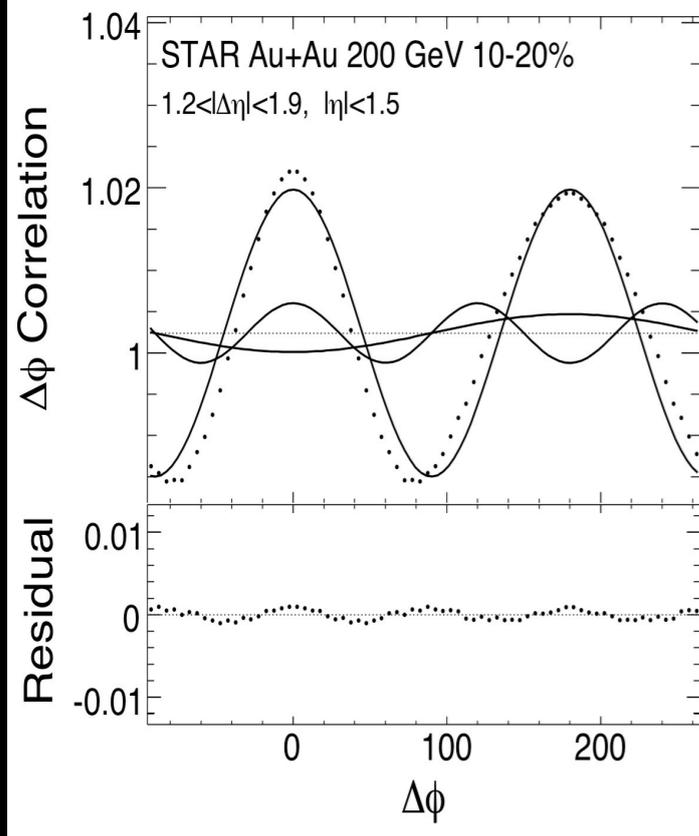


Burak Alver, GR, arXiv:1003.0194 (PRC in press)

## Inclusive correlations

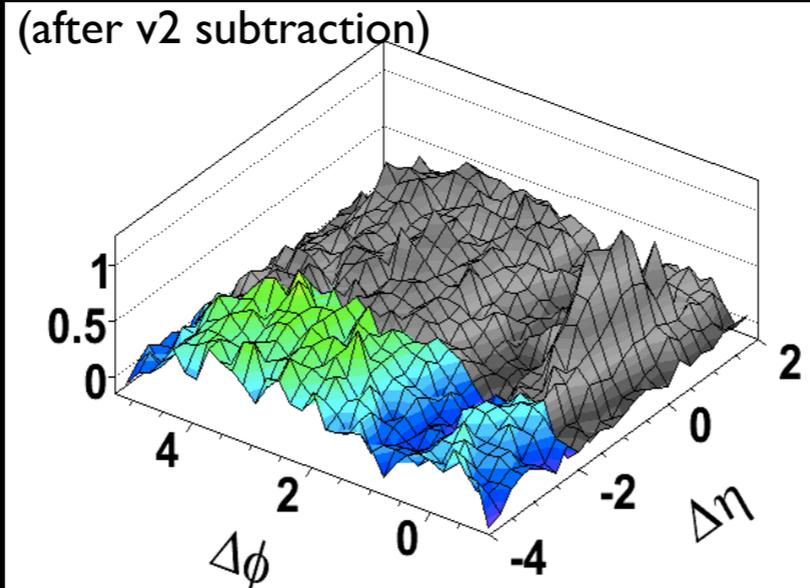


arXiv:0806.0513

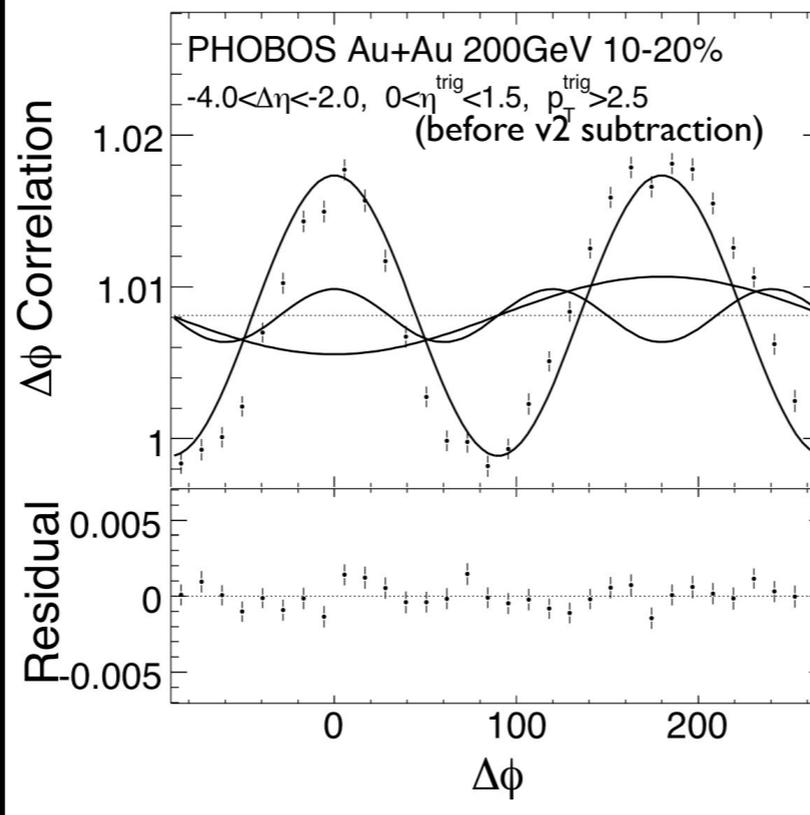


## Triggered correlations

(after  $v_2$  subtraction)



PRL 104, 06230 (2010)



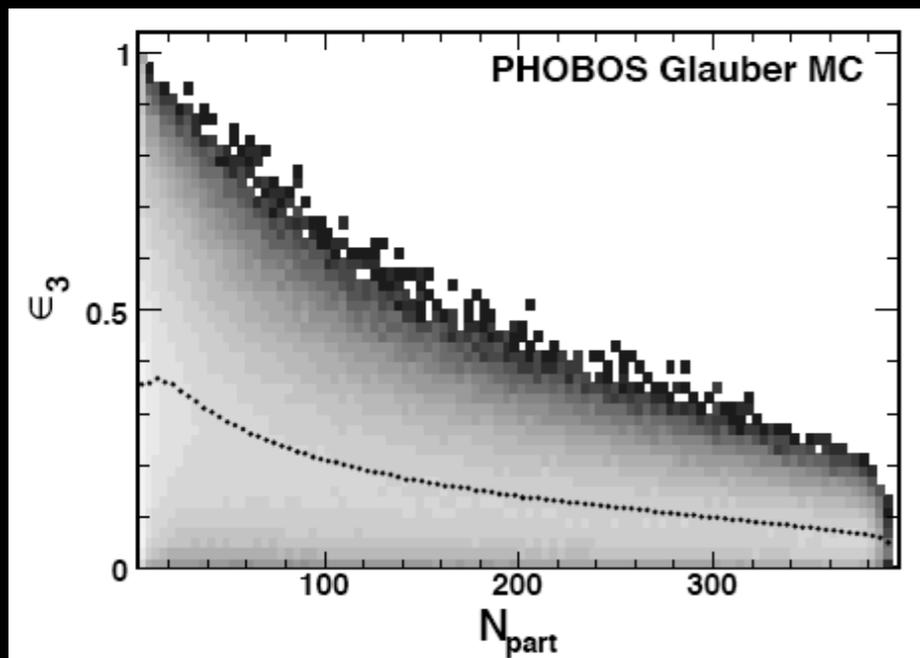
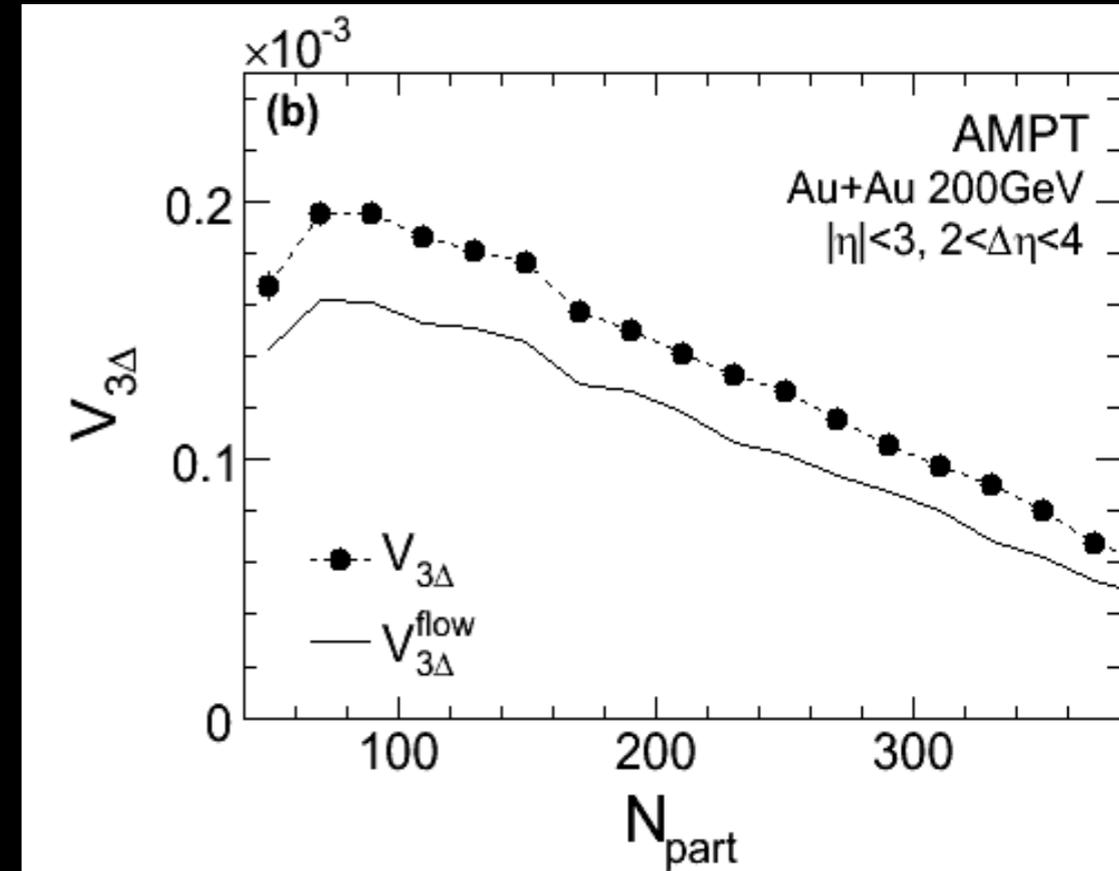
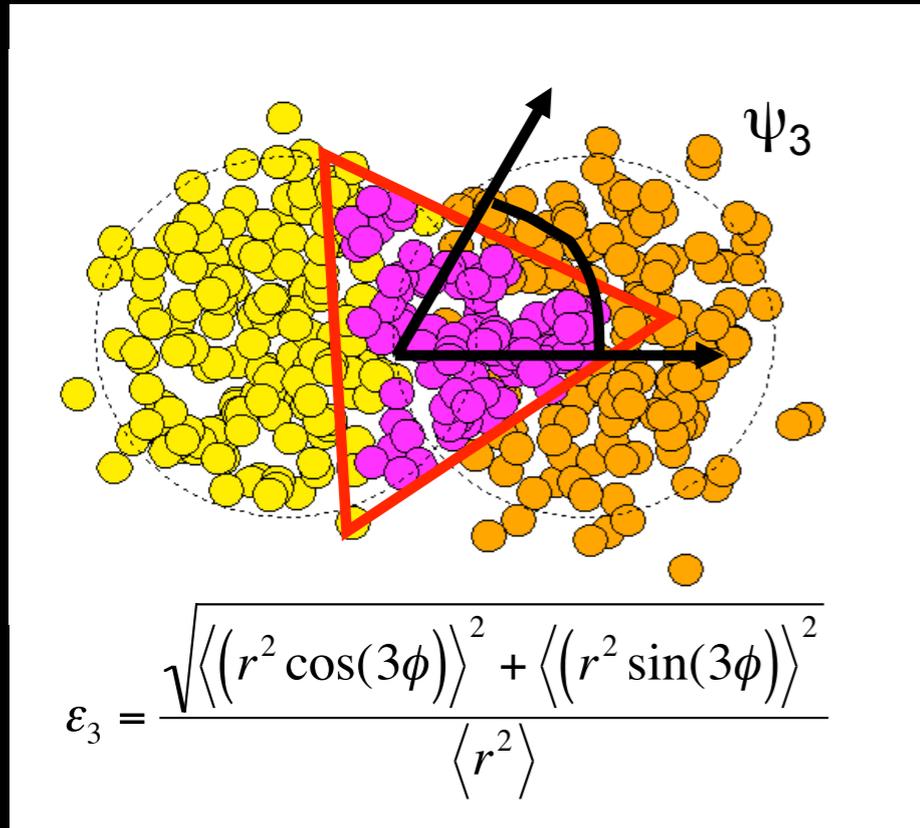
Published correlation data (STAR, PHOBOS) show  $v_3$  component!

Flow contribution to long-range “ridge” and “broad away-side”

This is purely a fluctuation effect - no fluctuations, no  $v_3$ !

n.b.  $\Psi_2$  and  $\Psi_3$  are uncorrelated - triangular flow is not visible in  $v_2$  event plane analysis

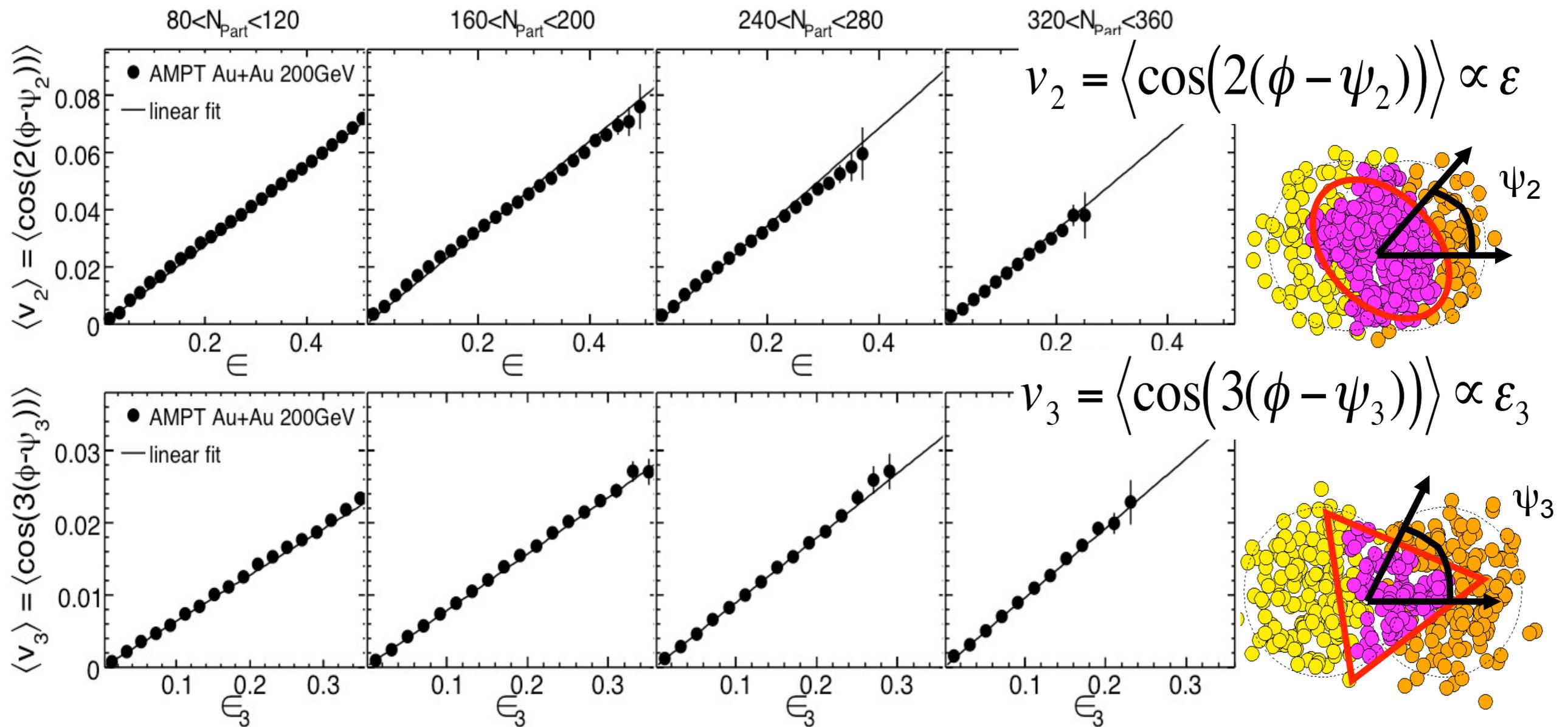
## Participant Triangularity



Just like elliptic flow reflects event-by-event eccentricity, “triangular flow” ( $v_3$ ) reflects event-by-event “triangularity” ( $\varepsilon_3$ )



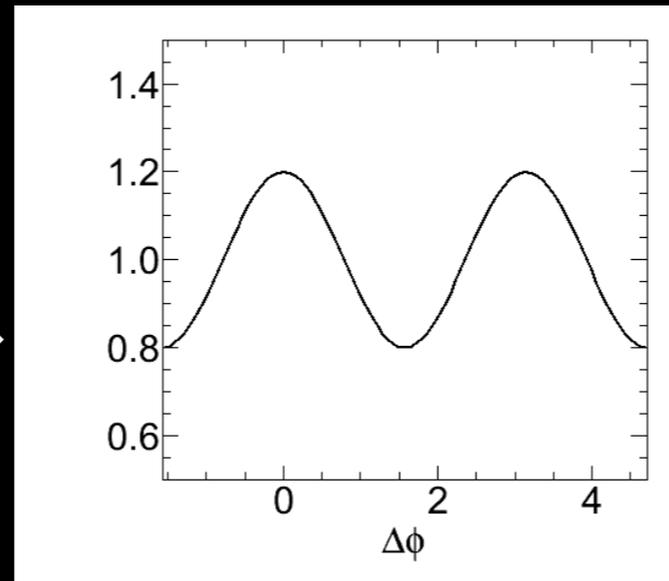
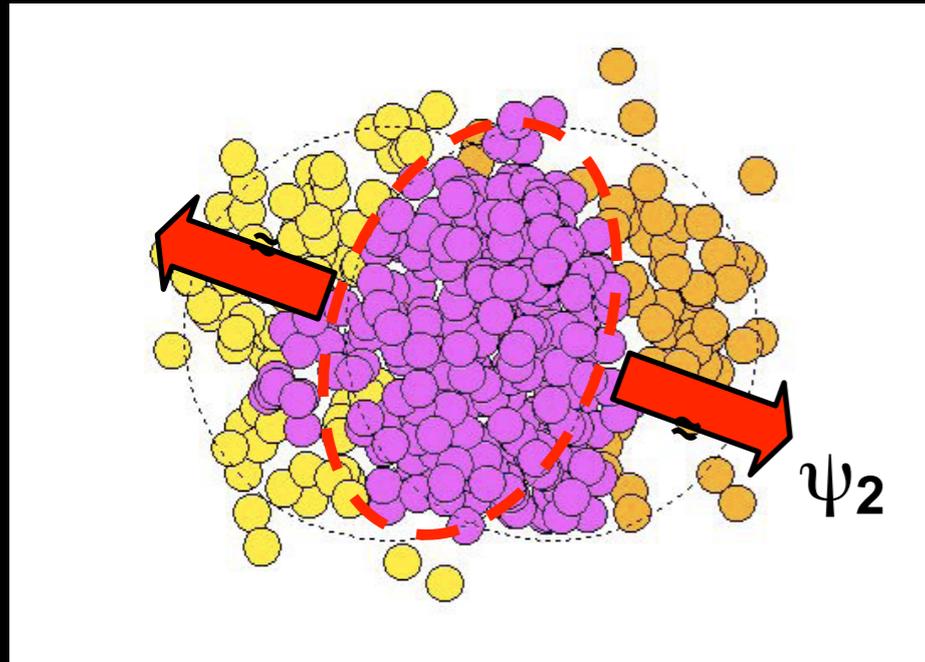
# Elliptic and triangular flow



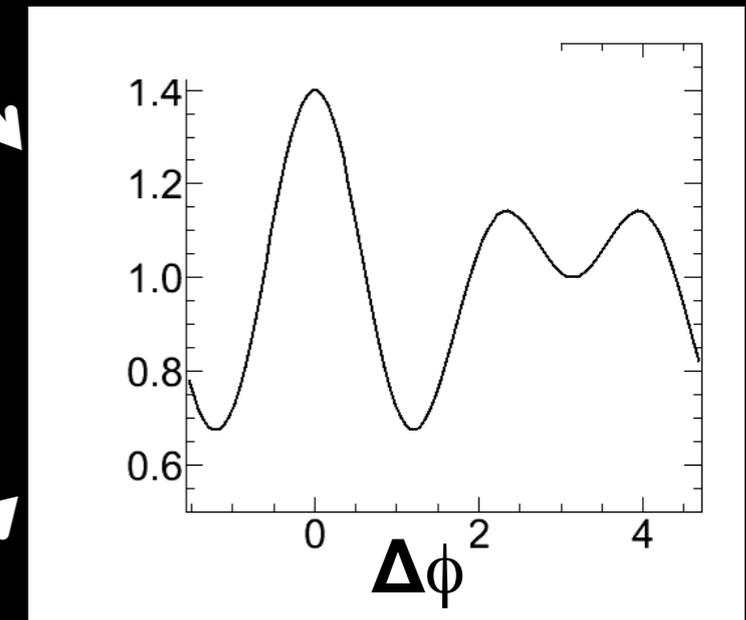
Burak Alver, GR, arXiv:1003.0194

# “No ridge, no mach cone, just flow”

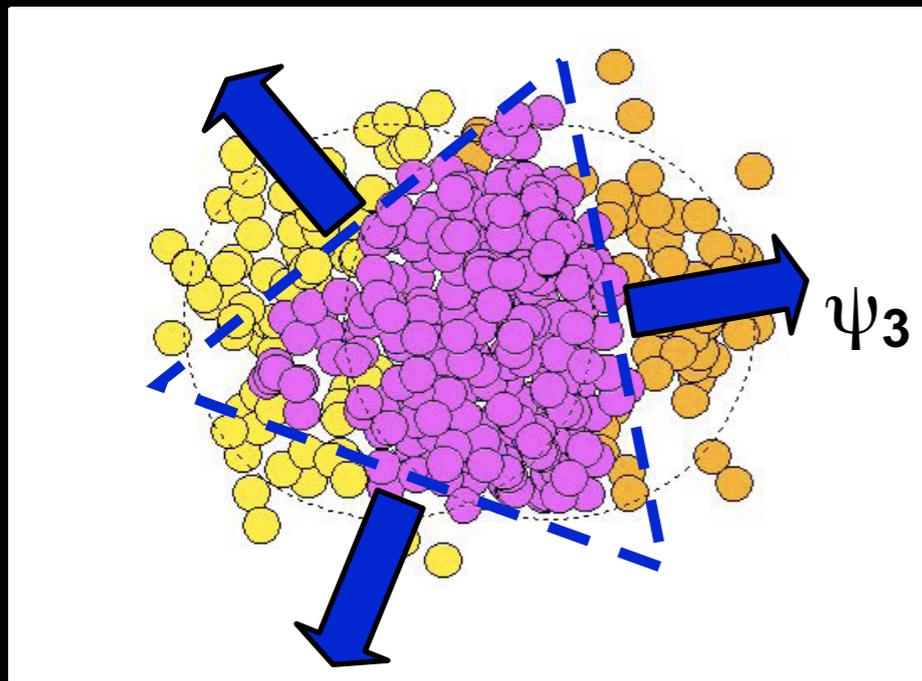
Elliptic flow ( $v_2$ )



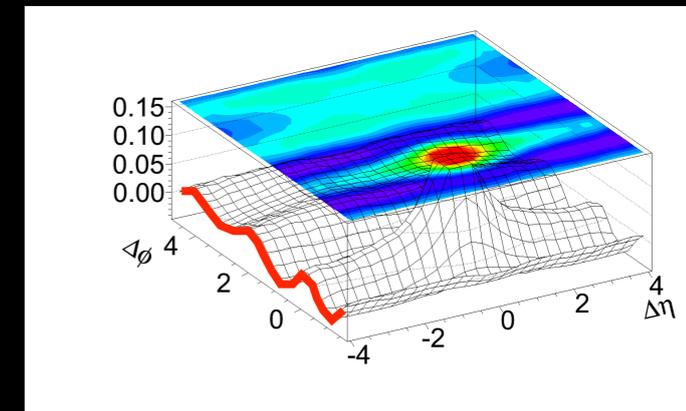
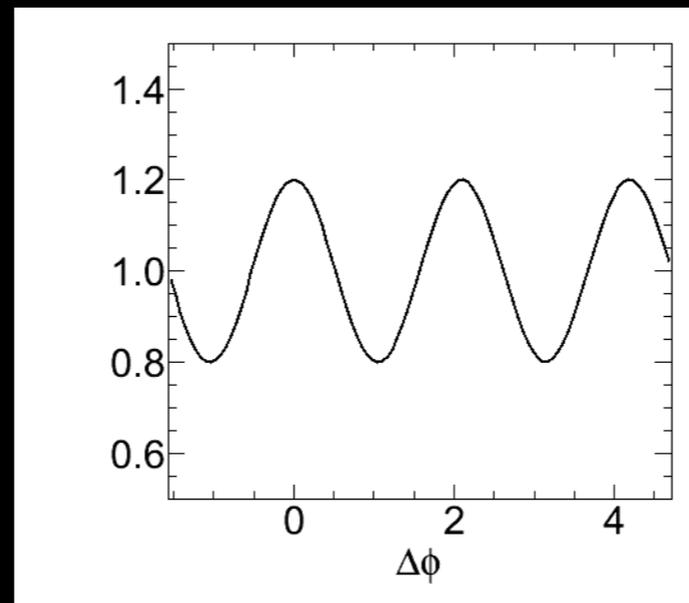
Add  $v_2^2$  and  $v_3^2$



Triangular flow ( $v_3$ ) from fluctuating initial condition



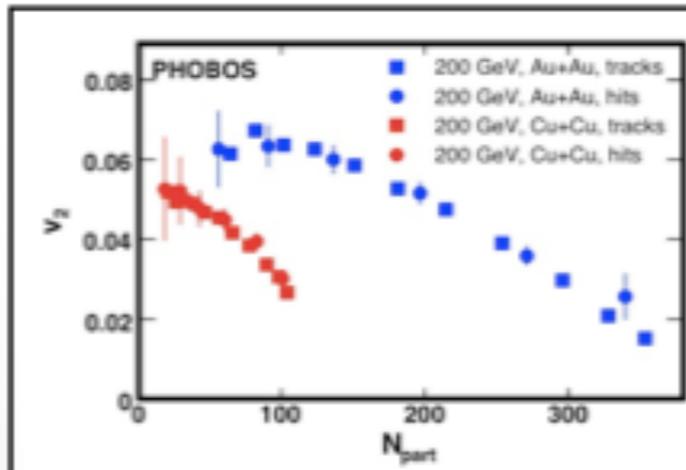
2-particle correlation functions



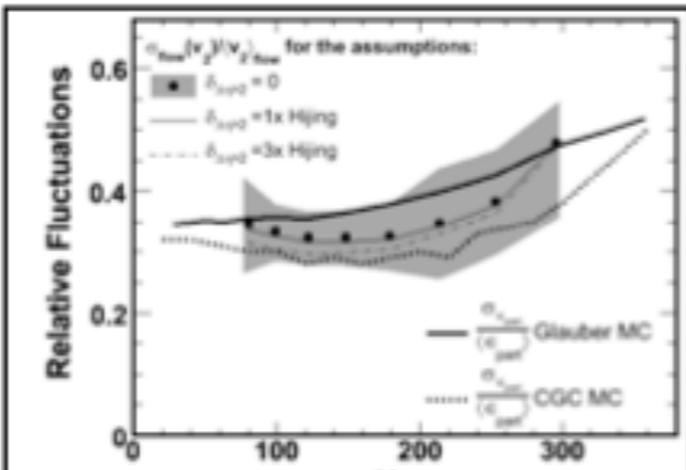


# Initial geometry fluctuations

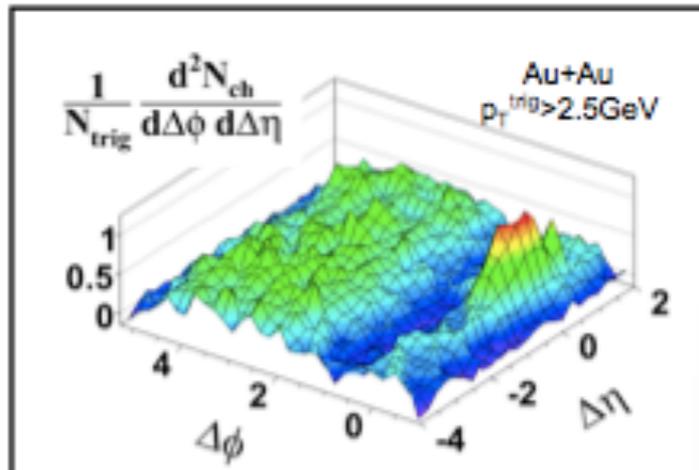
## A consistent picture



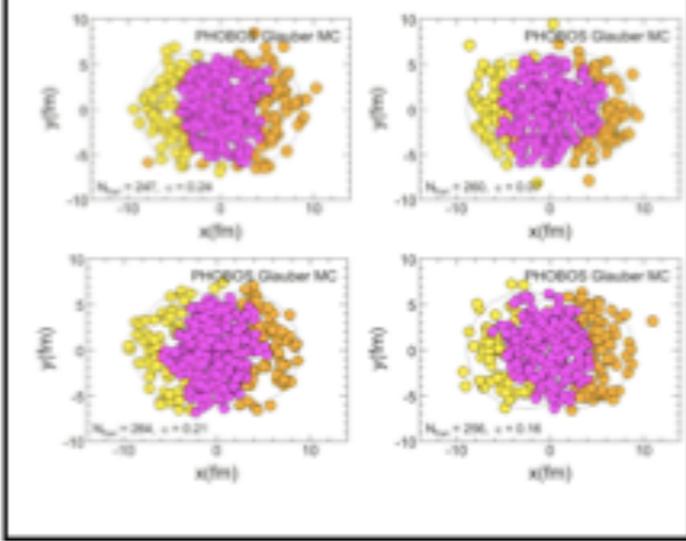
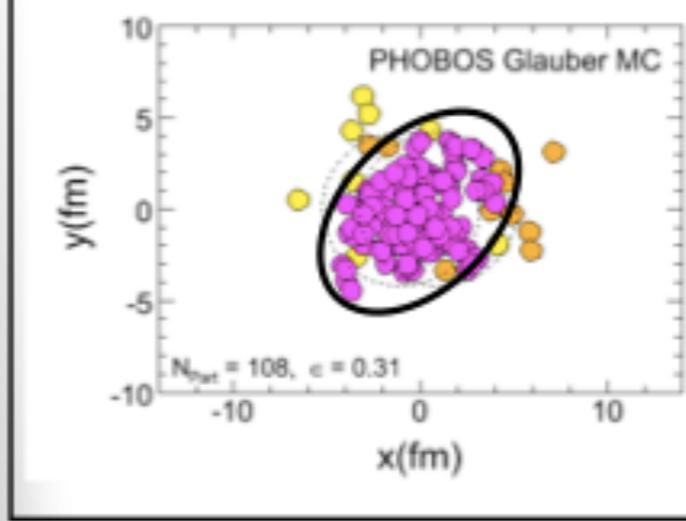
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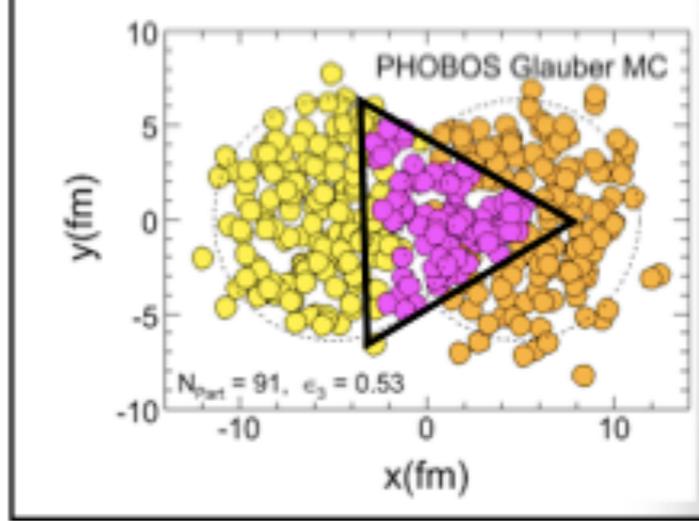
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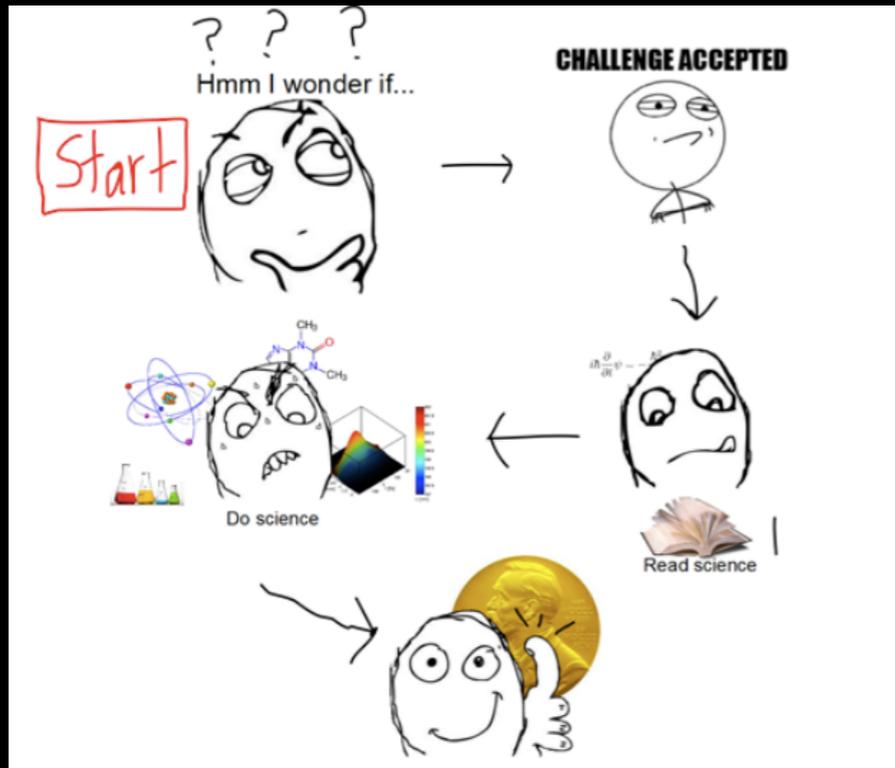
19



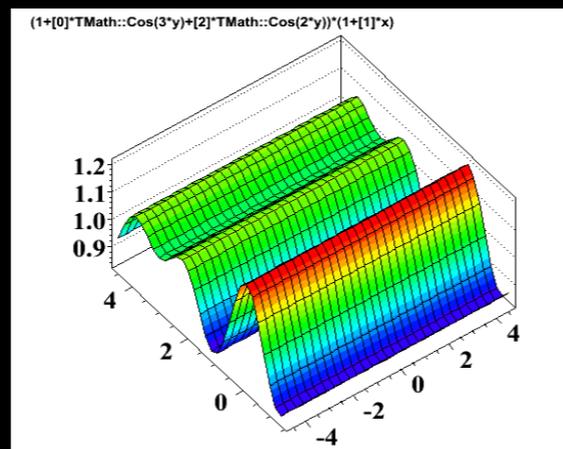
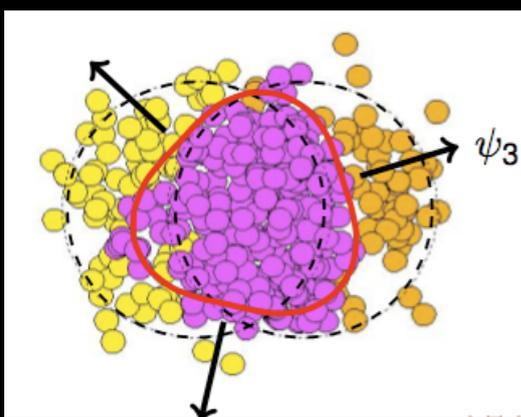
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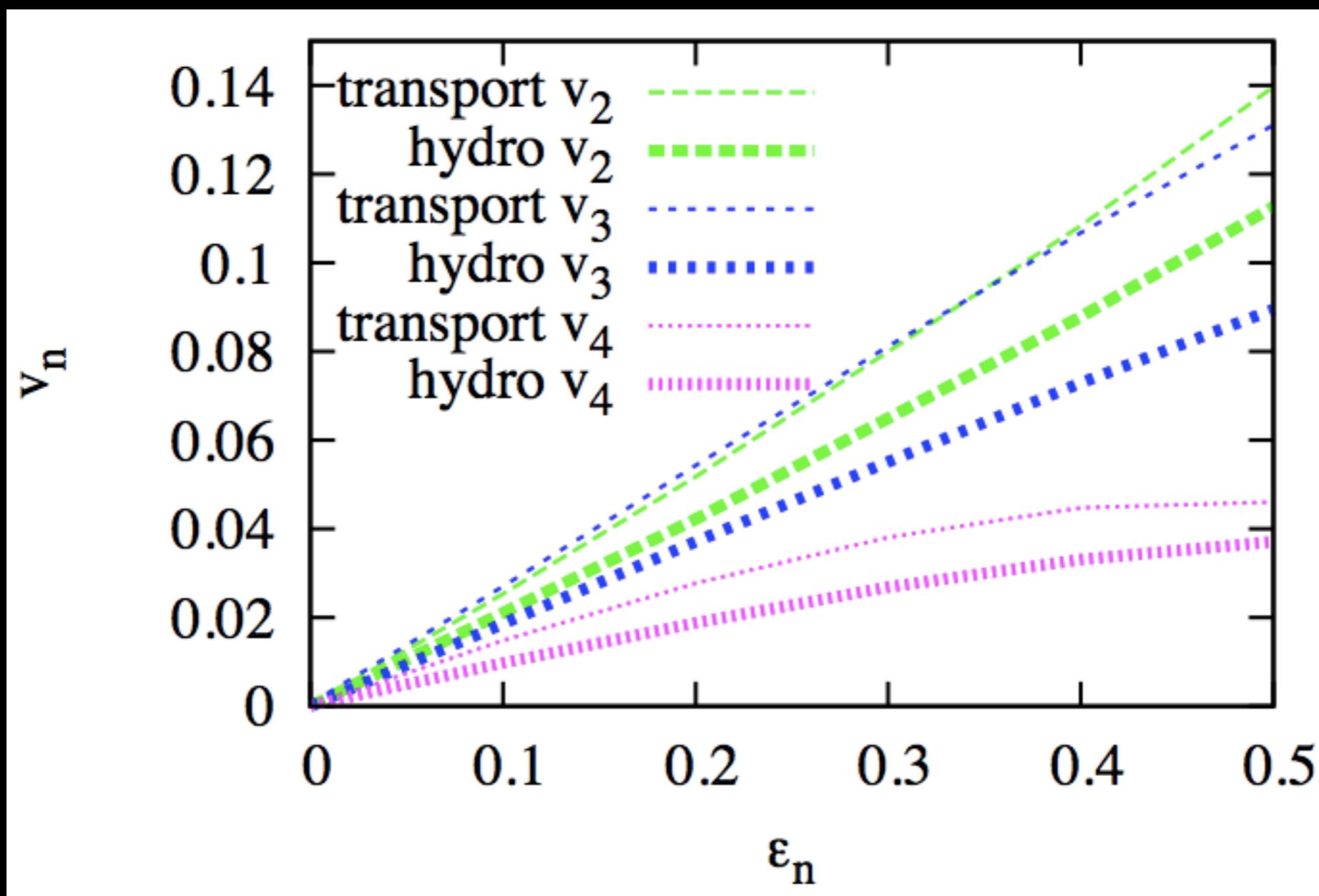


# Romantic notion of science

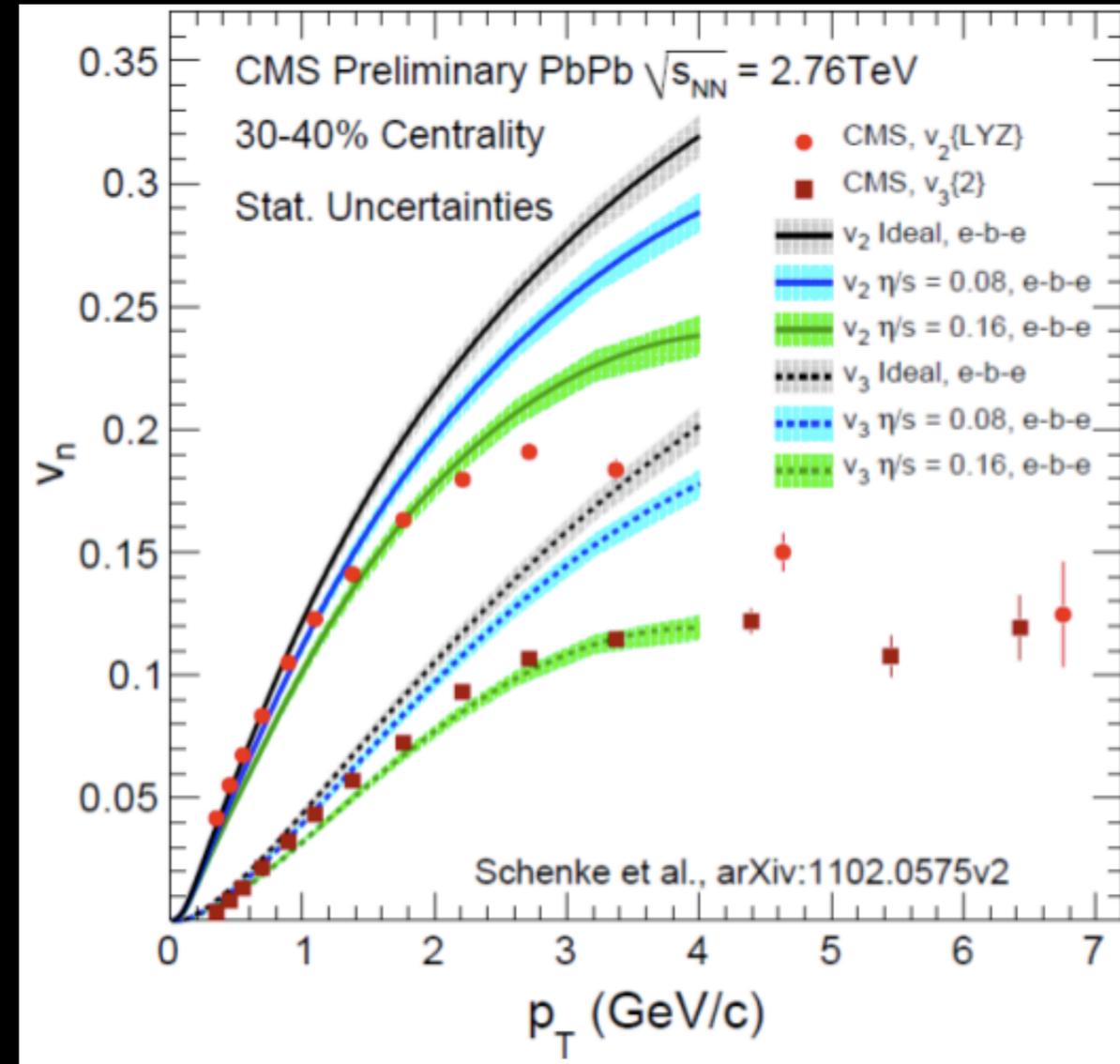
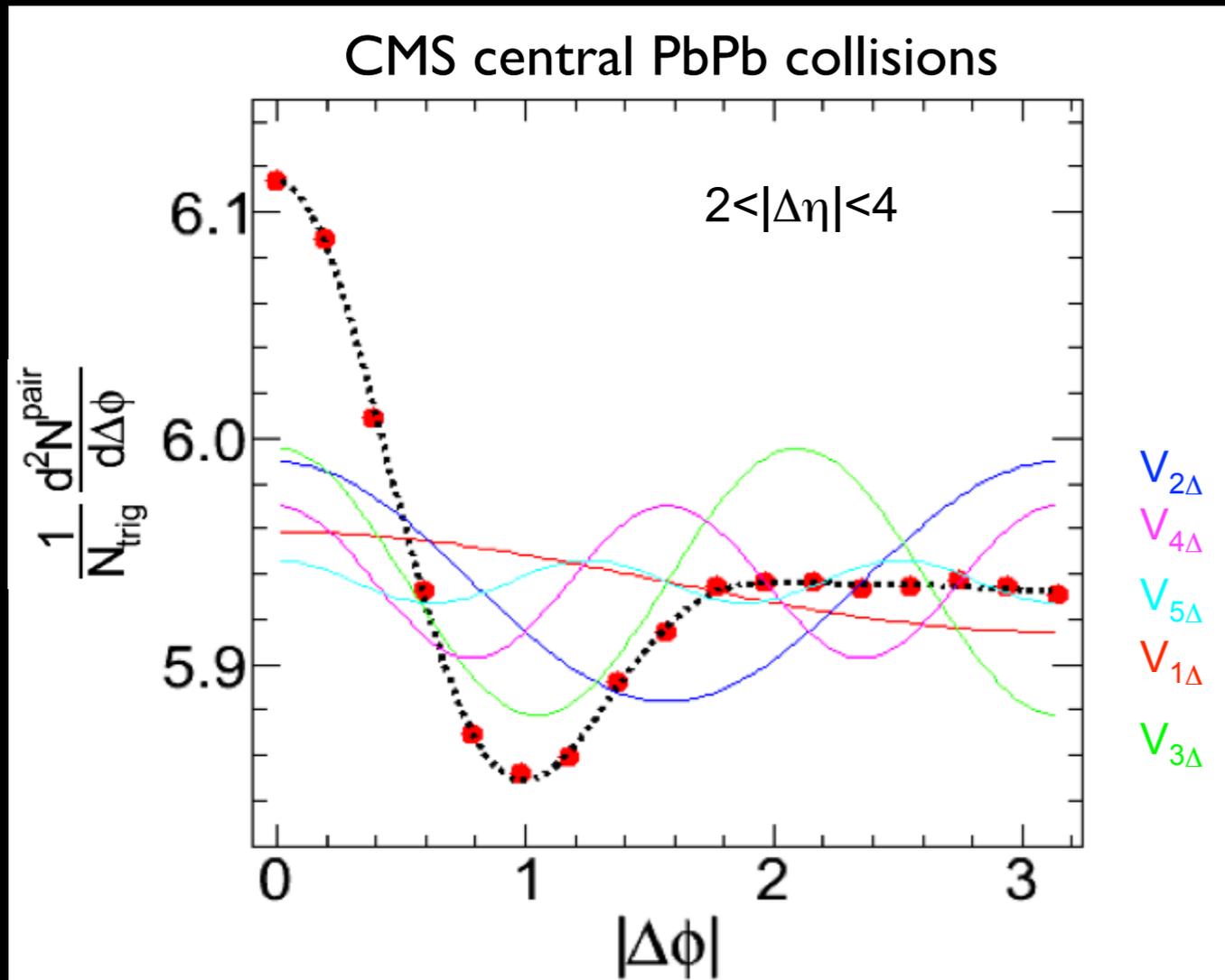


# Science in real life





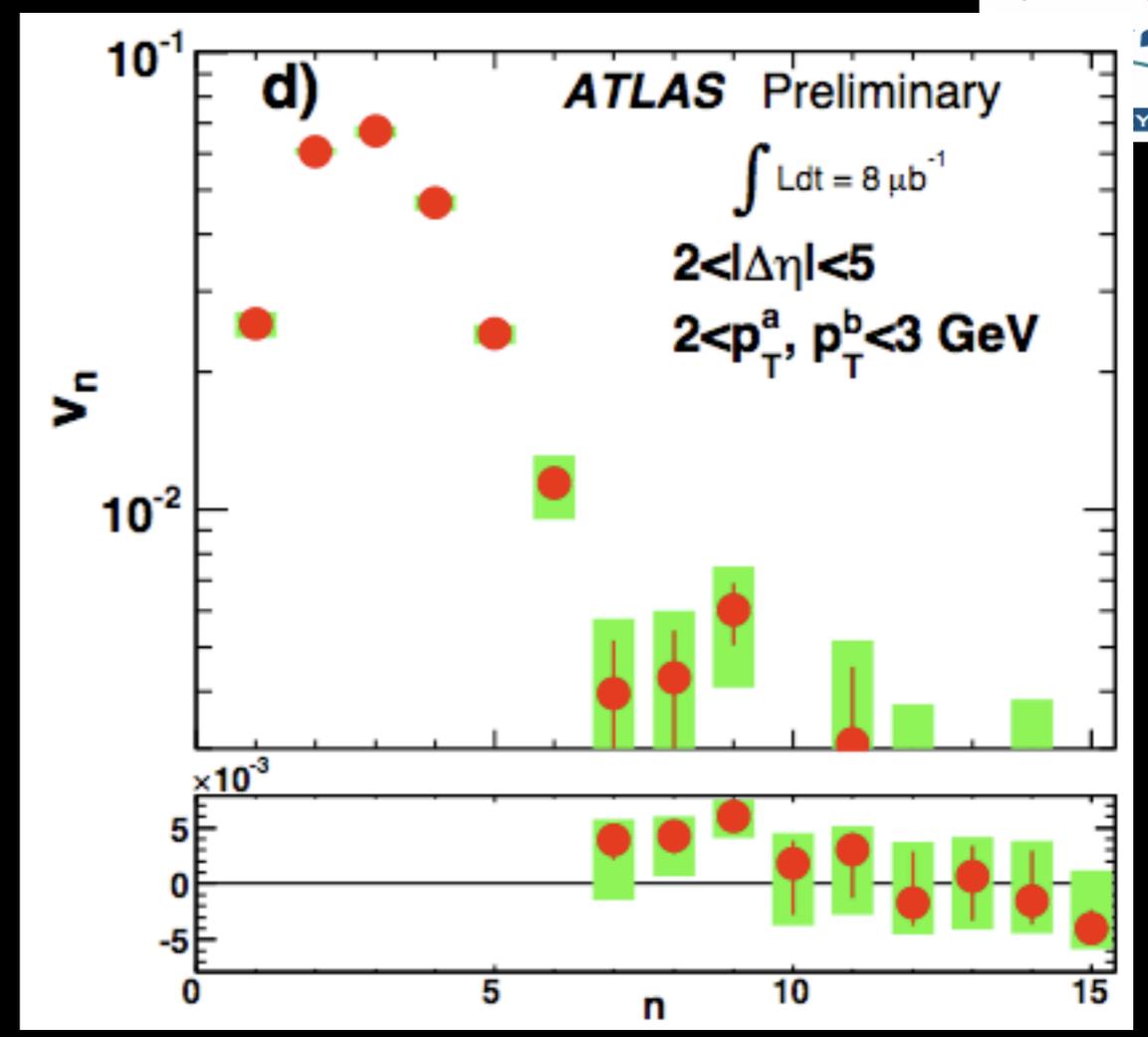
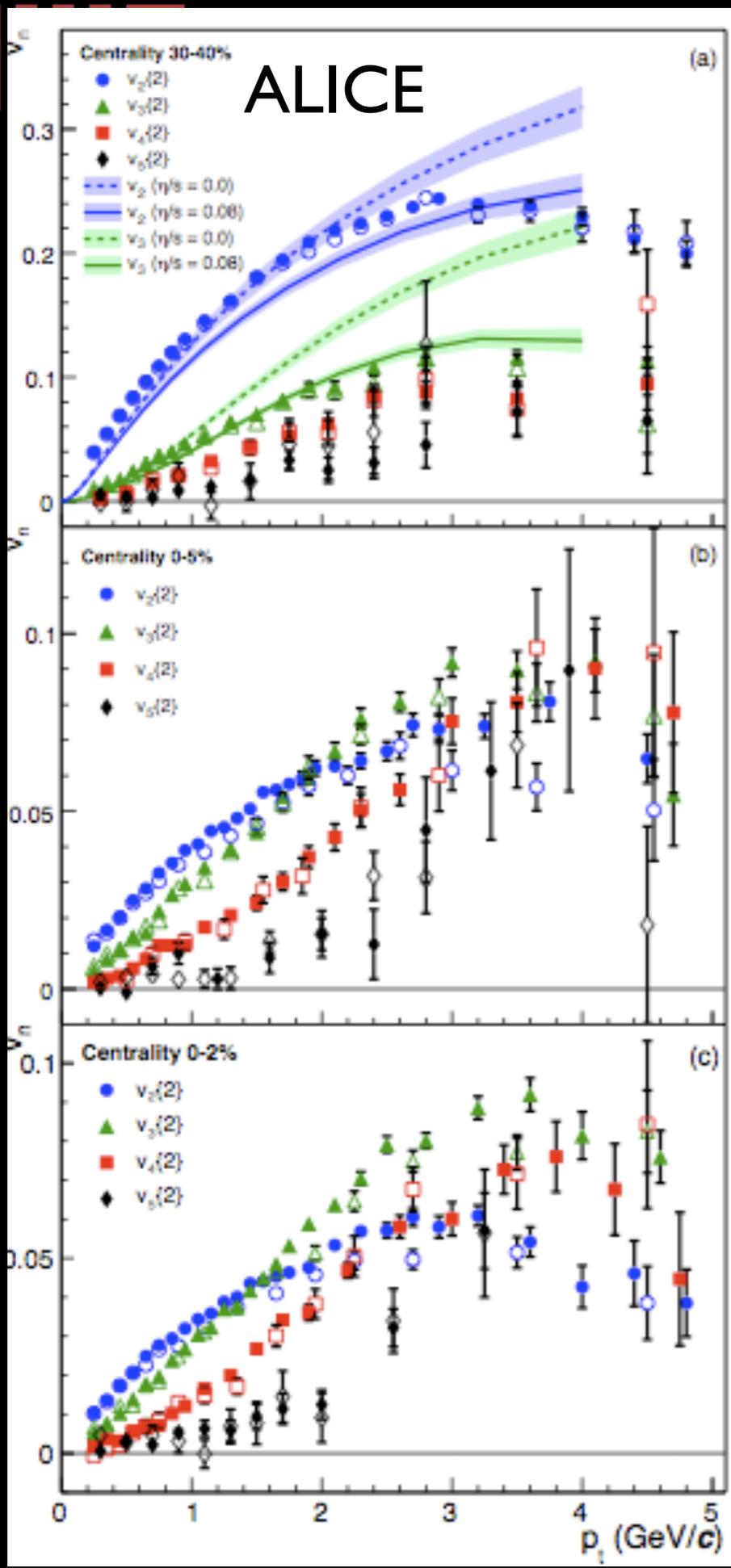
Alver, Gombeaud, Luzum, Ollitrault 1007.5496



Azimuthal correlations for central collisions (driven by shape fluctuations) show higher order Fourier components

(but many functions can be Fourier-decomposed...)

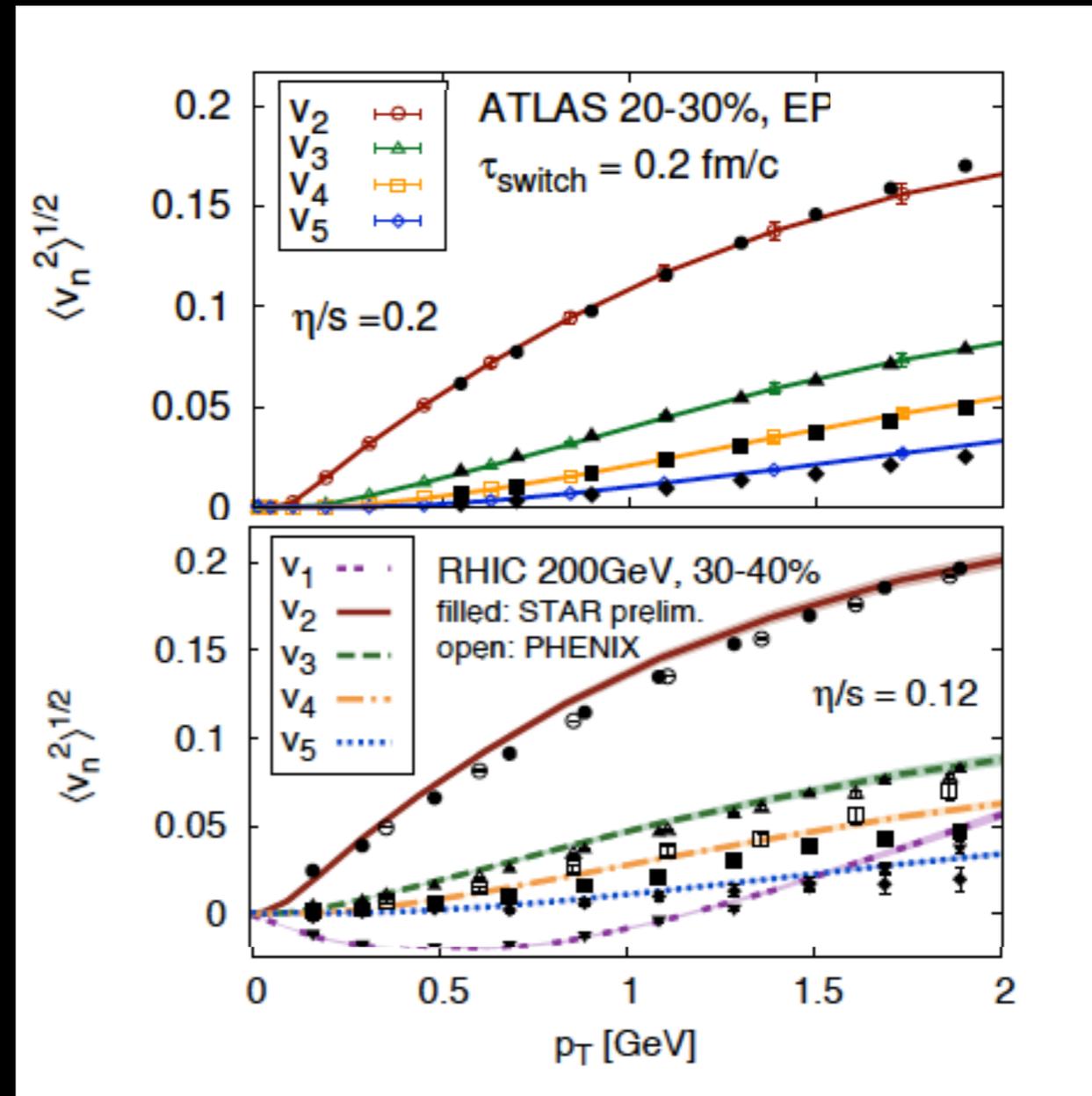
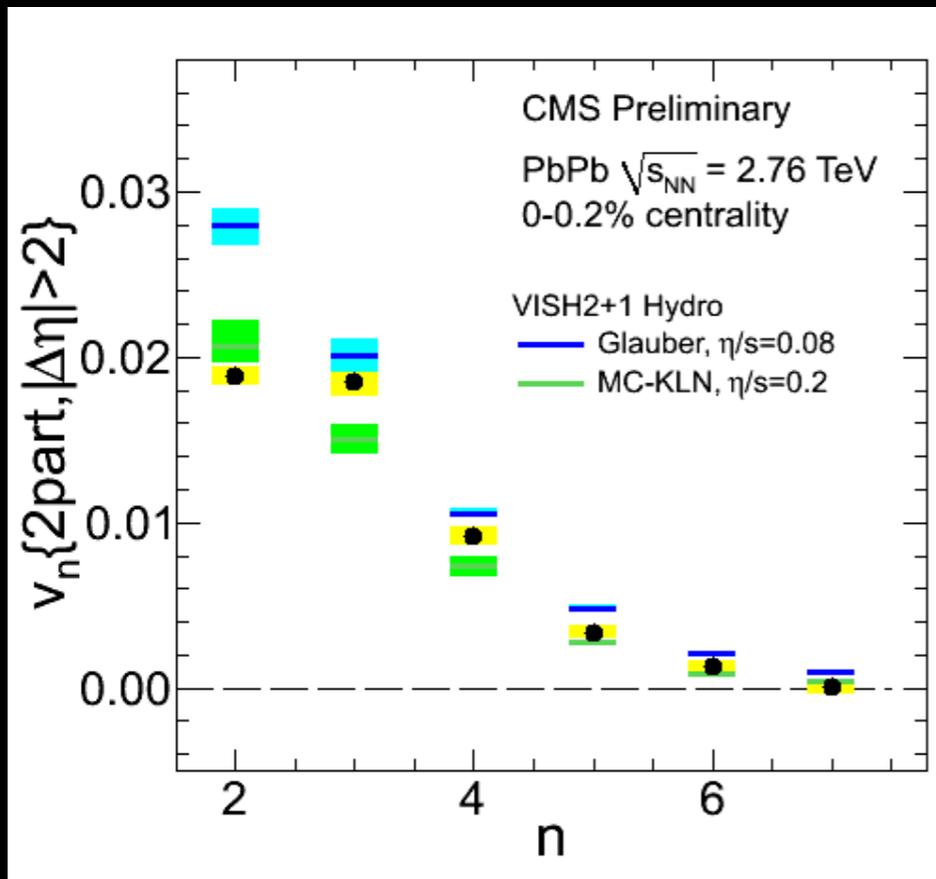
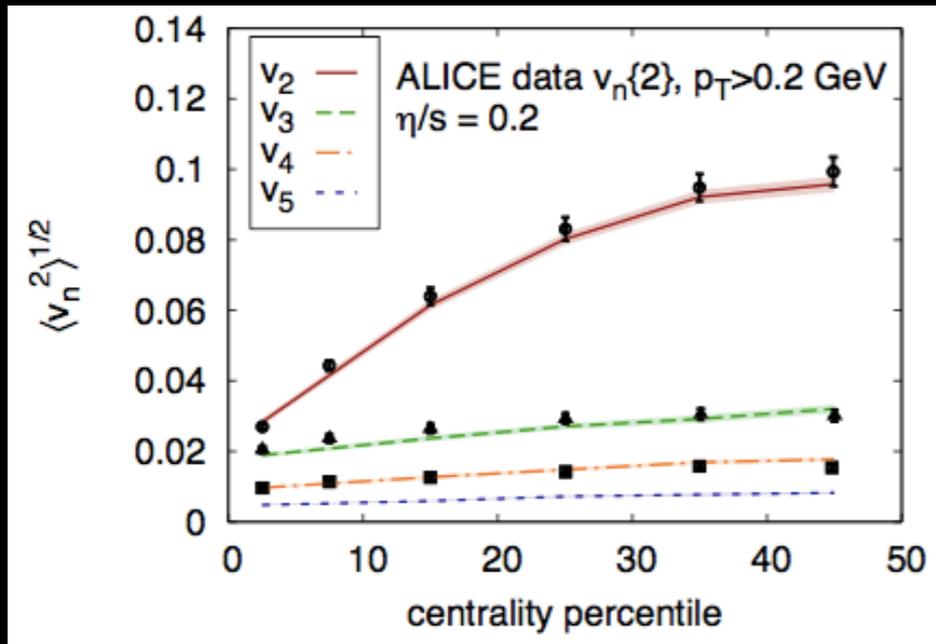
Proof is in the pudding:  
Full e-by-e viscous hydro calculations can describe  $v_3$



Fourier coefficients for central PbPb



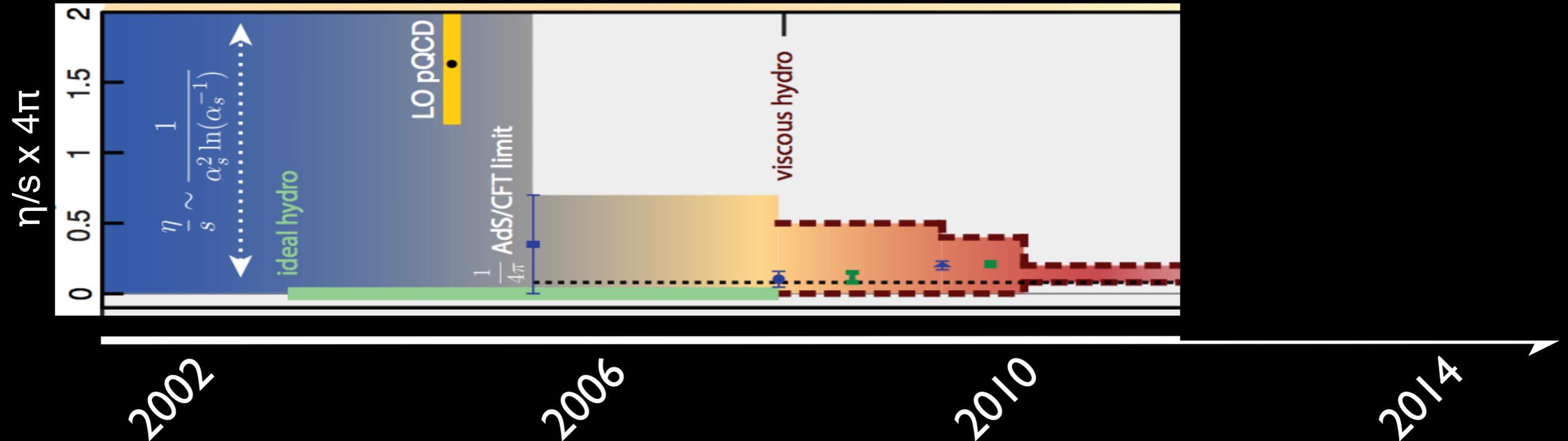
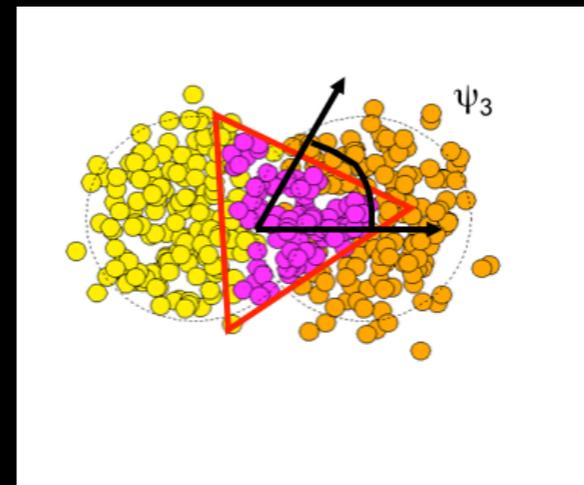
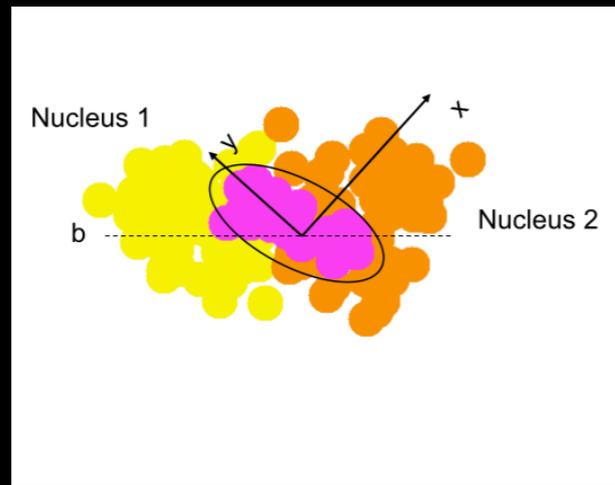
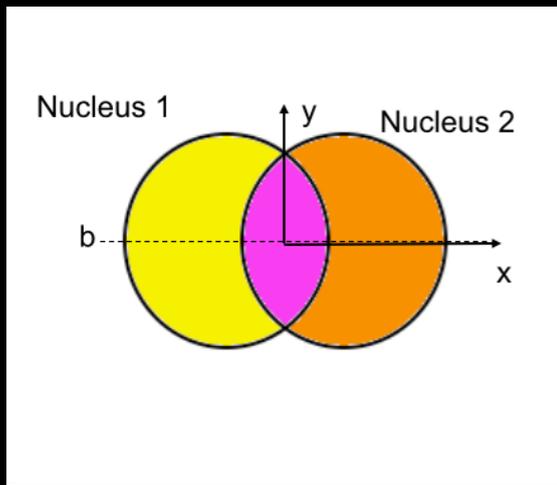
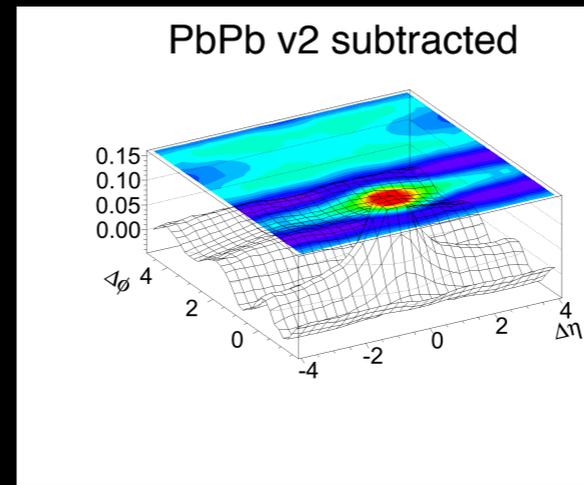
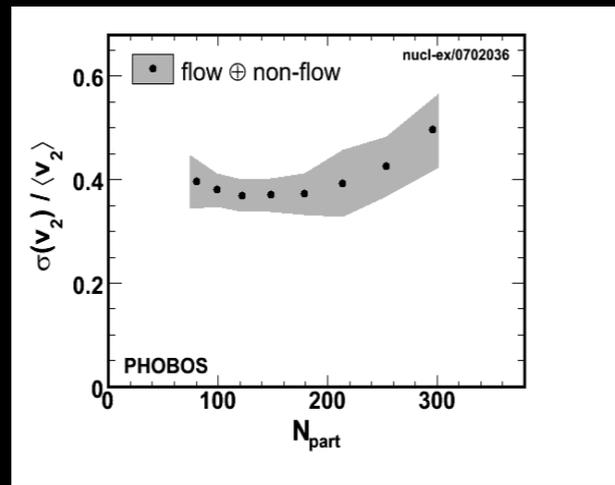
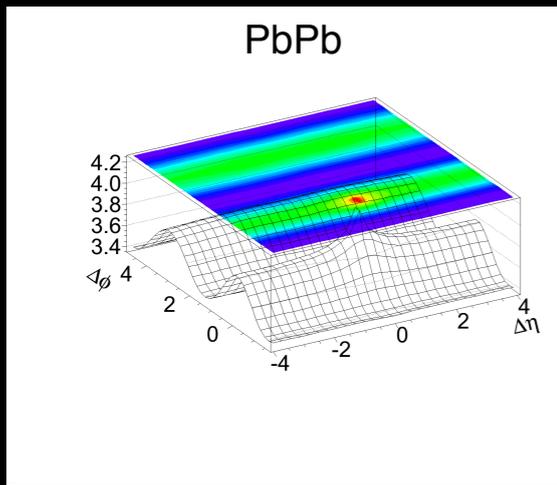
# Viscous hydro calculations vs data



Existing hydro calculations were able to describe the new flow components

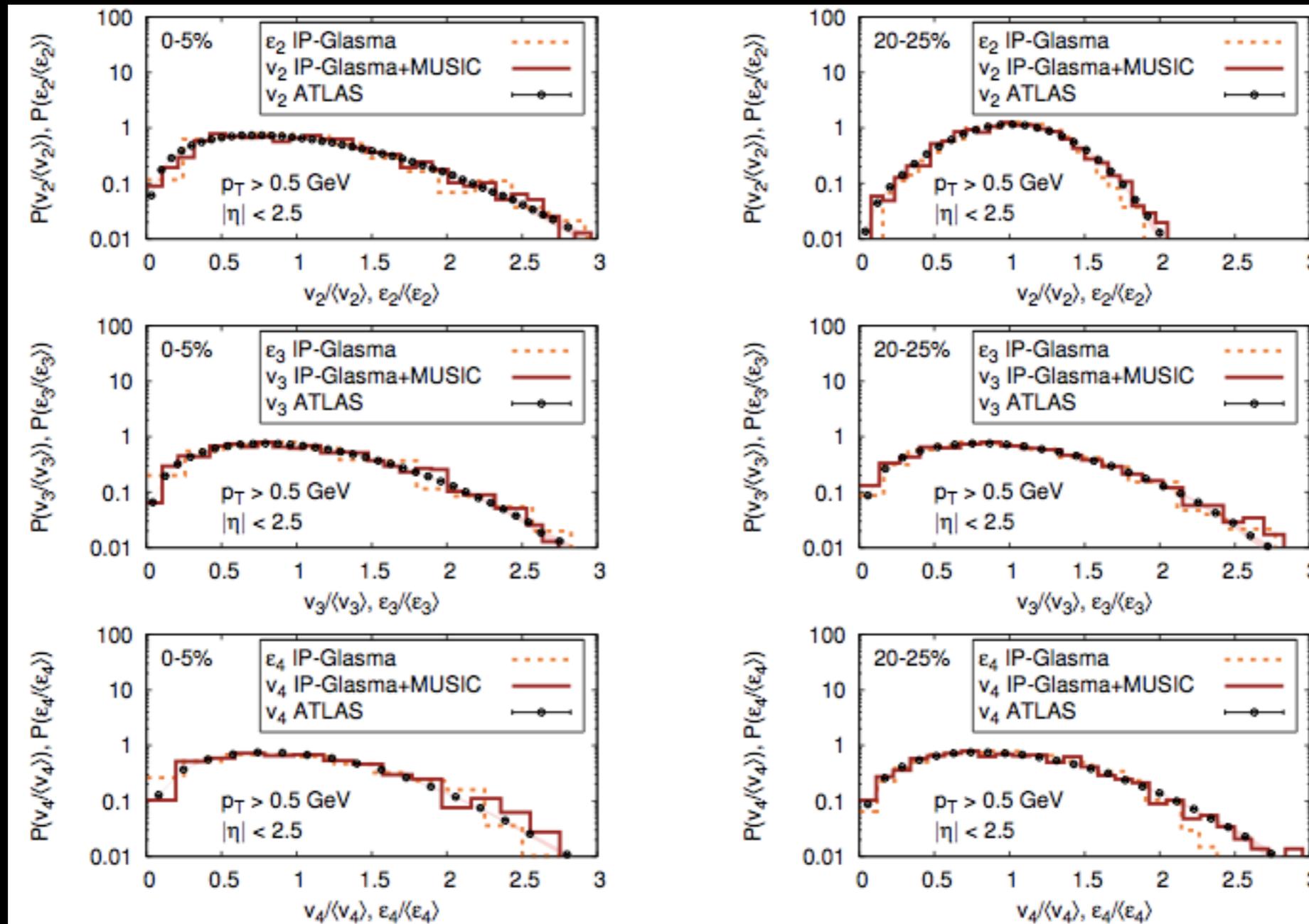


# A brief history of correlations in HIC



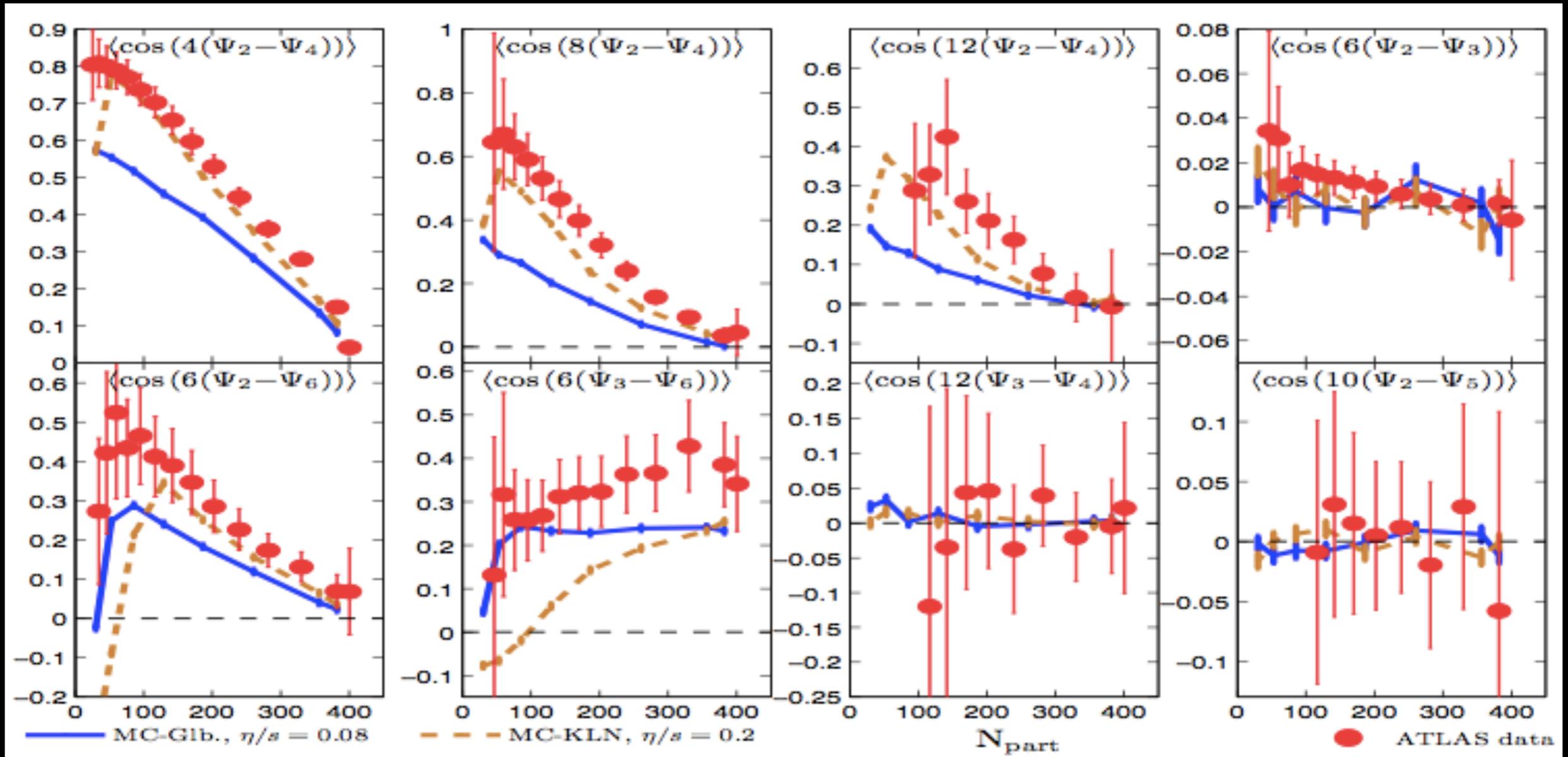


Has hydrodynamics ever predicted anything?



Information content per HI collision is obviously much lower than for CMB  
 But we have billions of events: Study event-by-event fluctuations

Calculations: Heinz et al.,



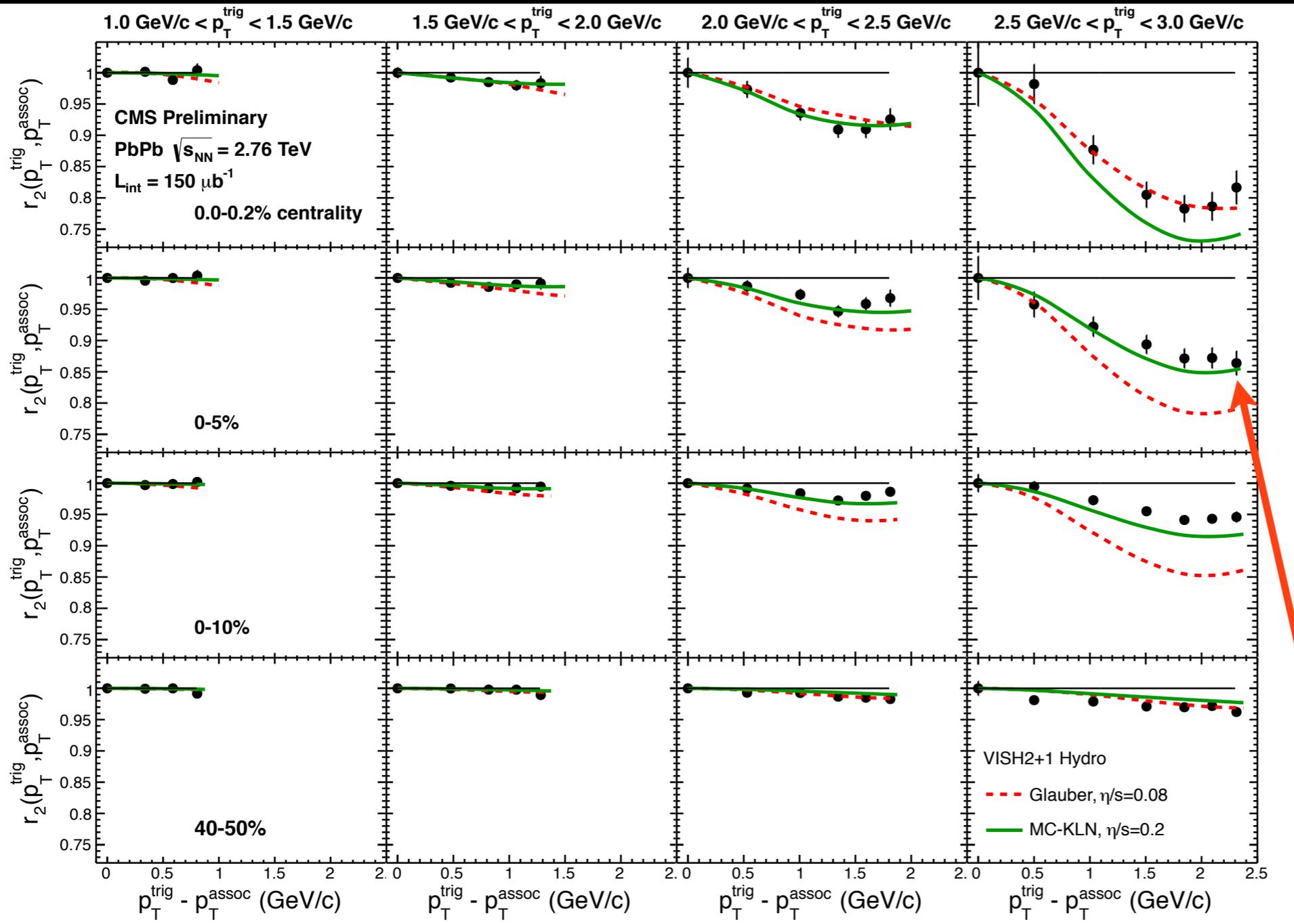
Fluctuations + hydro evolution lead to specific correlations of different order event plane angles



# Factorization breakdown



Calculations: Heinz et al., PRC 87, 034913 (2013)



$$r_n \equiv \frac{V_{n\Delta}(p_T^{trig}, p_T^{assoc})}{\sqrt{V_{n\Delta}(p_T^{trig}, p_T^{trig})} \sqrt{V_{n\Delta}(p_T^{assoc}, p_T^{assoc})}}$$

Factorization is broken as fluctuations lead to  $p_T$  dependent event-plane angle

# QCD-inspired models

- **MC-KLN**:  $k_T$  factorization, positions of the nucleons fluctuate

*Drescher and Nara, Phys. Rev. C 76, 041903 (2007), Albacete and Dumitru, arXiv:1011.5161*

- **MCrcBK**: + fluctuations added to match multiplicity distribution in pp collisions

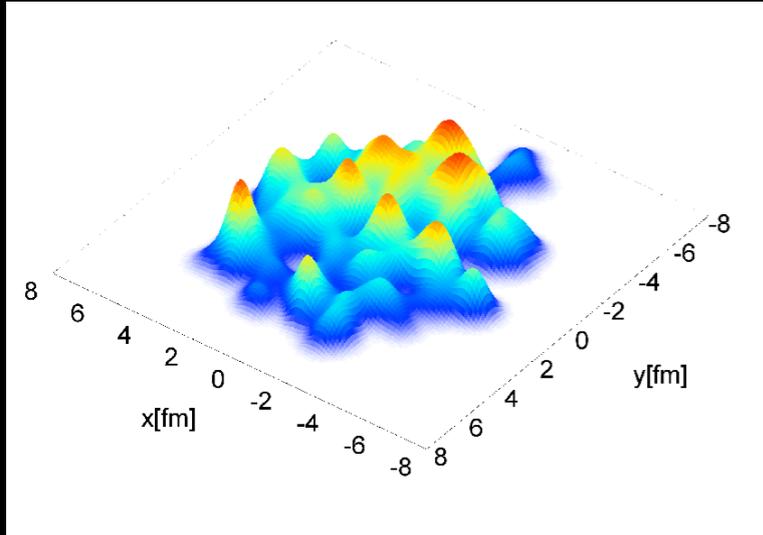
*Dumitru and Nara, Phys. Rev. C 85, 034907 (2012)*

- **DIPSY**: +multiple gluon cascade

*Flensburg, arXiv:1108.4862*

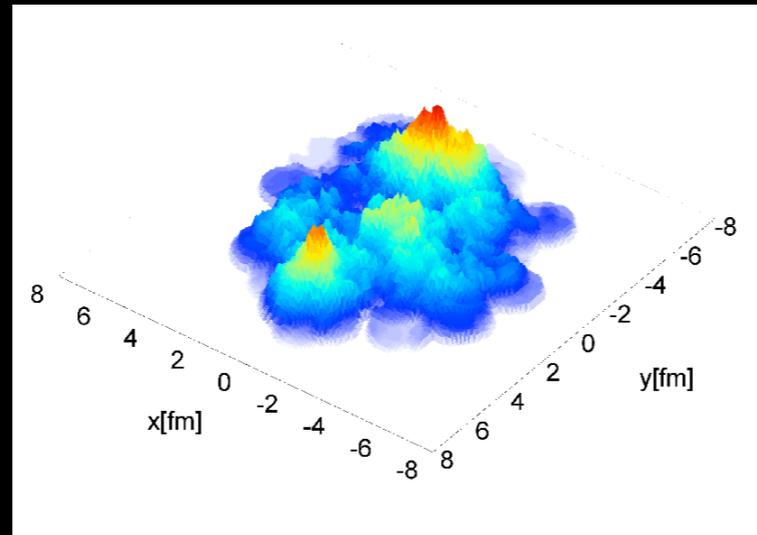
- **IP Glasma**: no  $k_T$  factorization, non-linearities, fluctuations of color charges within a nucleon

*Schenke, Tribedy, Venugopalan: PRL108 (2012), 252301*



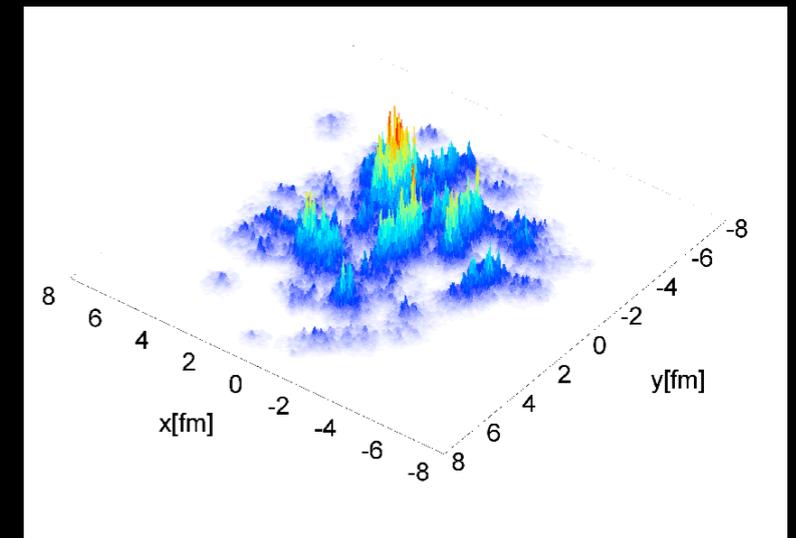
## MC Glauber

(Billiard-ball nucleons)



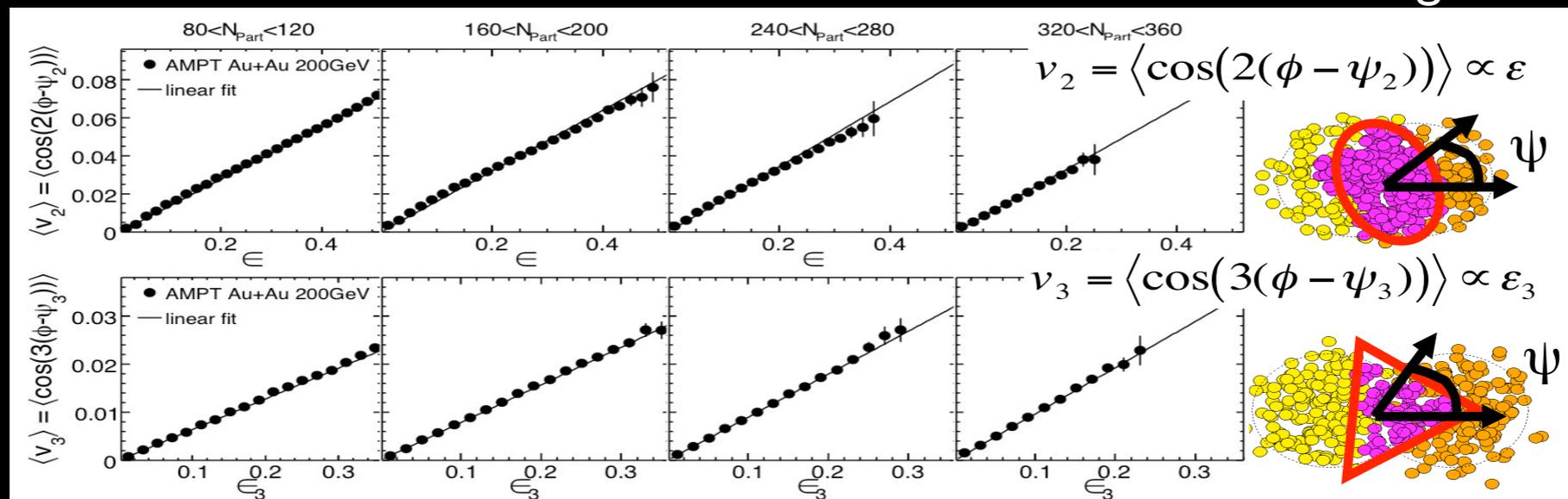
## MC KLN

Nucleonic fluctuations  
+ saturation effects

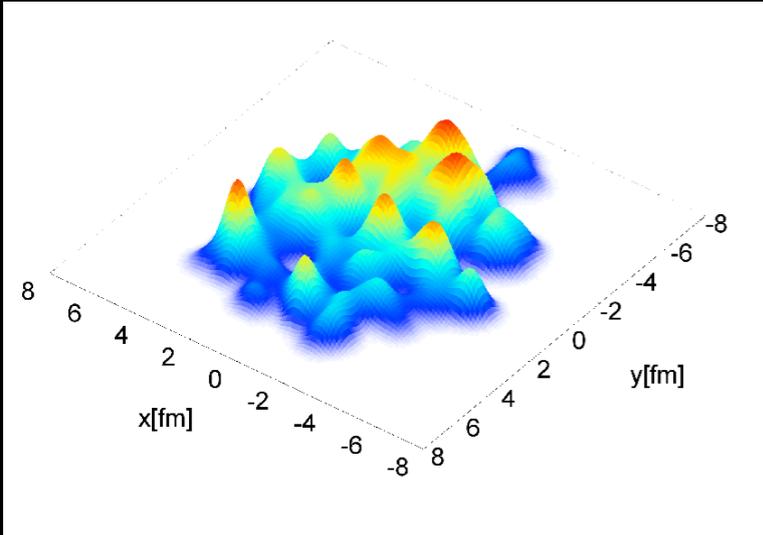


## IPGlasma

Nucleonic fluctuations  
+ quantum fluctuations of  
gluon fields

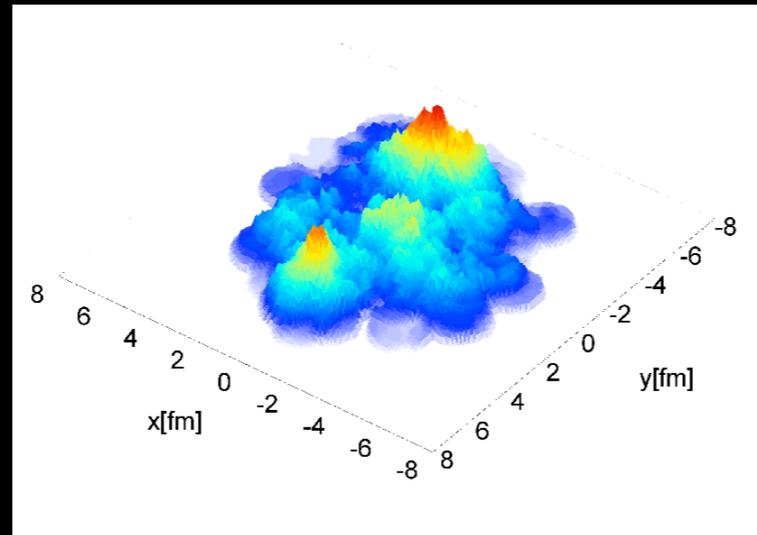


Recall: Linear response of elliptic ( $v_2$ ) and triangular flow ( $v_3$ ) to initial eccentricities  $\epsilon_2$  and  $\epsilon_3$

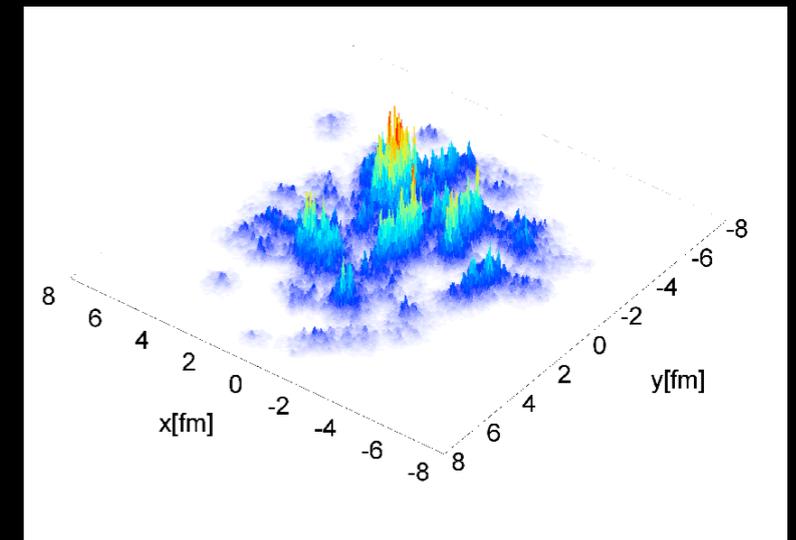


## MC Glauber

(Billiard-ball nucleons)



## MC KLN

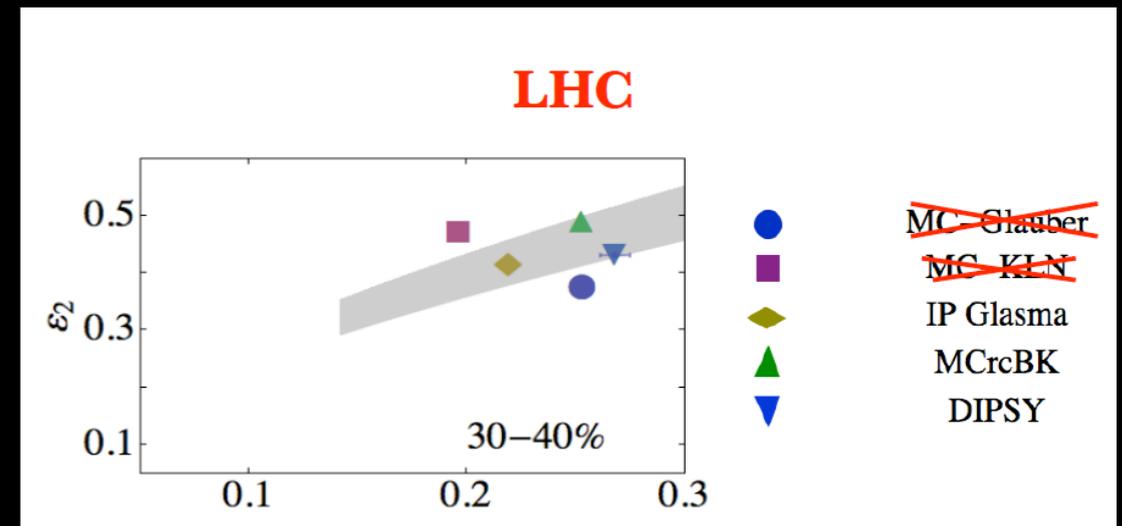
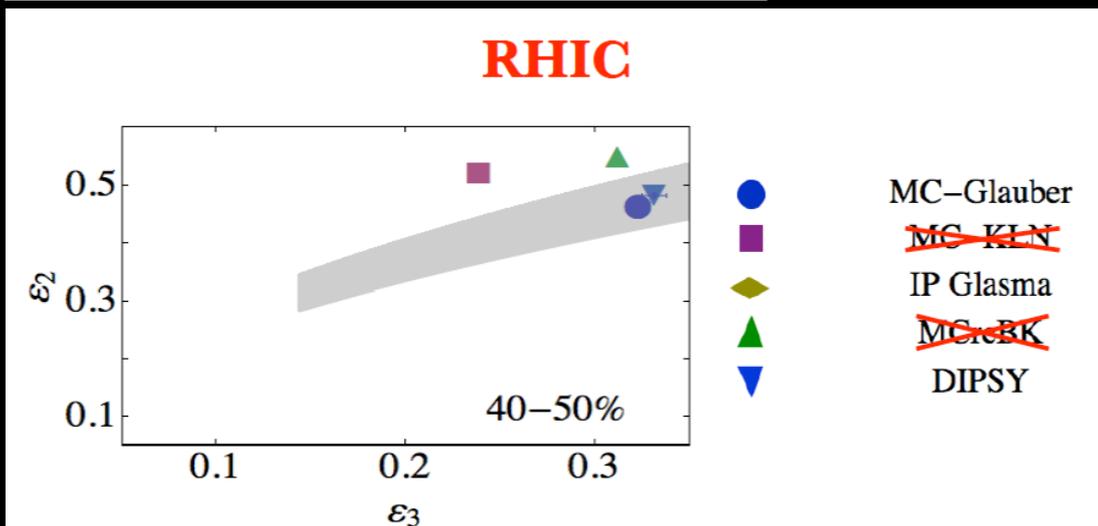


## IPGlasma

Nucleonic fluctuations  
+ quantum fluctuations of  
gluon fields

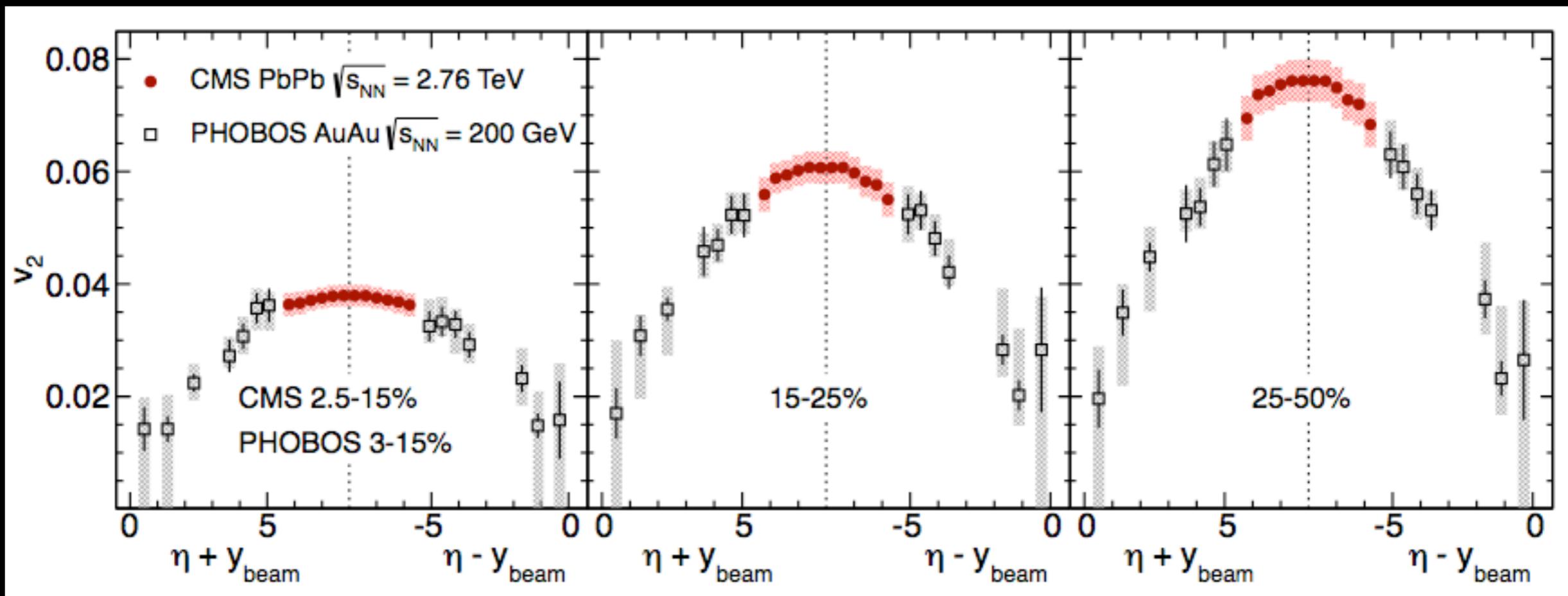
$$\sqrt{\langle \varepsilon_2^2 \rangle} = C \left( \sqrt{\langle \varepsilon_3^2 \rangle} \right)^k$$

Parametrized  
hydro response





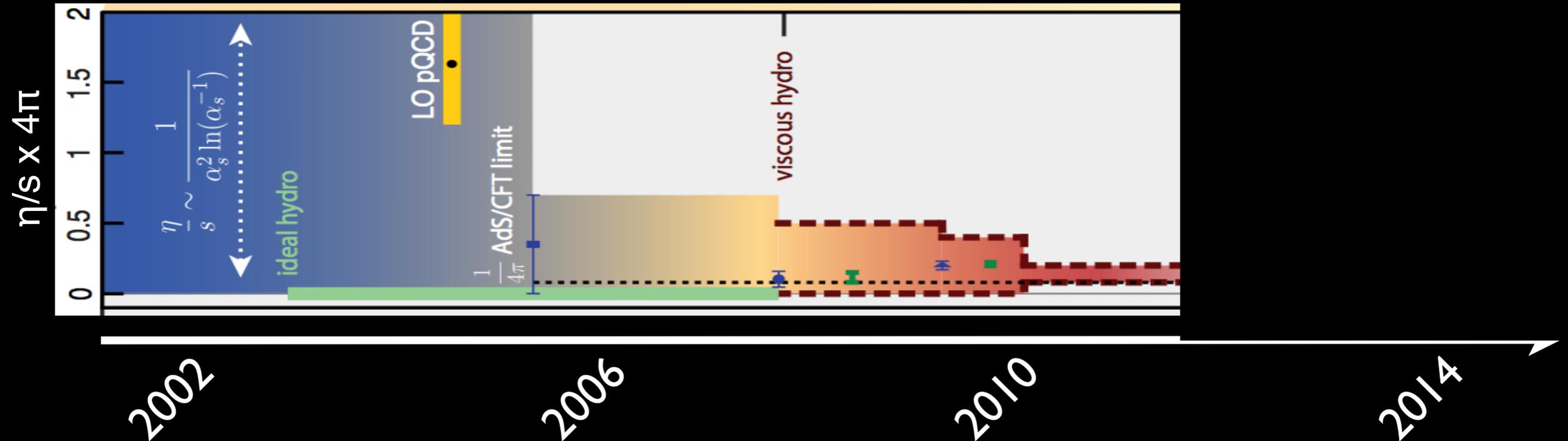
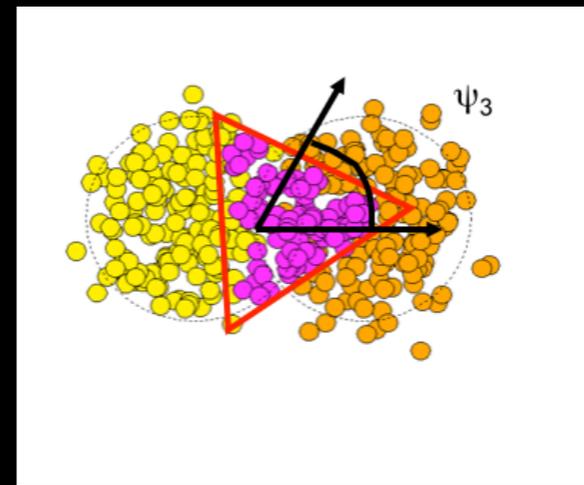
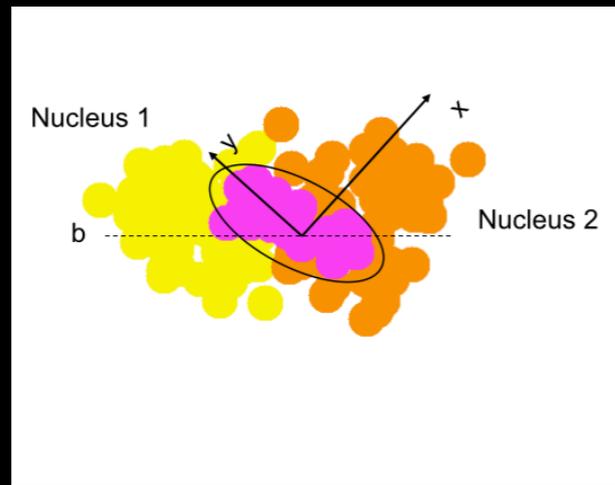
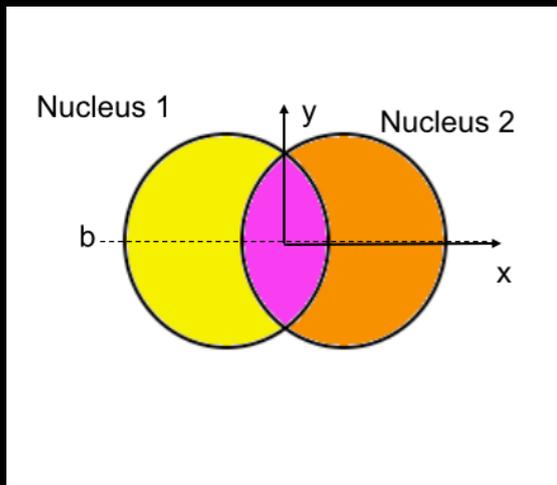
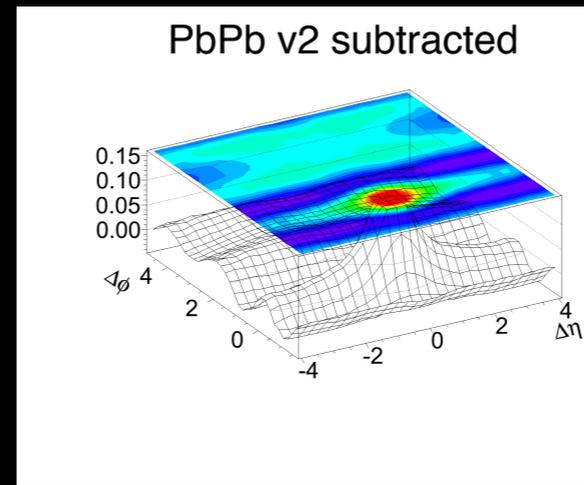
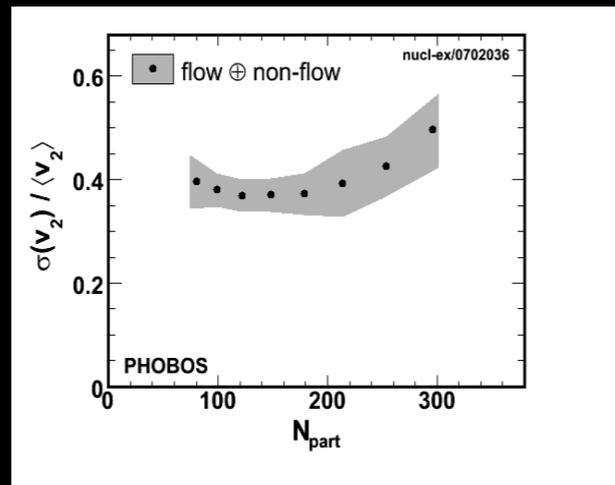
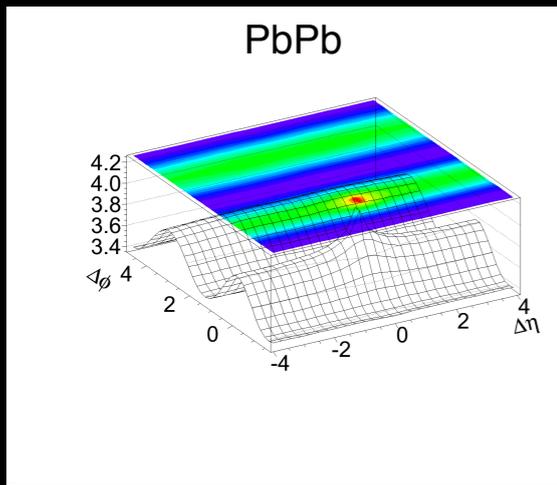
# What about the longitudinal direction?



Profound questions remain, e.g. what is the longitudinal structure of the source?



# A brief history of correlations in HIC



# Beyond elliptic flow

PHENIX/PHOBOS/STAR  
Au+Au 200 GeV  
2006-2010

ALICE/ATLAS/CMS  
PbPb 2.76 TeV  
Nov 2010-Mar 2011

